



**Solution of Large Underestimation Problem in the
Monte Carlo Calculation with Hard Biasing
-In Case with Geometry Input Data Created
by CAD/MCNP Automatic Converter-**

Hiromasa IIDA, Nobuo KAWASAKI*, Chikara KONNO
Satoshi SATO and Akiyuki SEKI

Fusion Neutronics Group
Fusion Research and Development Directorate

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〒319-1195 茨城県那珂郡東海村白方白根 2 番地 4
電話 029-282-6387, Fax 029-282-5920, E-mail: ird-support@jaea.go.jp

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Tel +81-29-282-6387, Fax +81-29-282-5920, E-mail: ird-support@jaea.go.jp

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–In Case with Geometry Input Data Created by CAD/MCNP Automatic Converter–**

Hiromasa IIDA, Nobuo KAWASAKI*, Chikara KONNO, Satoshi SATO and Akiyuki SEKI⁺

Division of Fusion Energy Technology
Fusion Research and Development Directorate
Japan Atomic Energy Agency
Tokai-mura, Naka-gun, Ibaraki-ken

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An inconvenient experience was encountered, in which we have different answers depending on applied weight window values, in the nuclear analysis of the benchmark problem for CAD/MCNP interface programs, being developed under the ITER R&D task. Biasing can enhance calculation speed, but should not give different answers. Mechanism of this large underestimation is clarified. It is caused by the combination of the following two facts;

- * When one of particles in a history has got lost, MCNP cancels all tallies calculated during the history and all banked particles are thrown away (never tracked).
- * When we have distributed micro geometry errors in input data, important histories, which give significant contribution to tallies, will have many splitting and have “lost particle” with higher probability in the case of hard biasing.

These two facts lead to selective canceling of important histories. An attempt to eliminate this inconvenience has been made, by modifying the subroutine “hstory” of MCNP. The modification has been done very successfully and eliminated the large underestimation, giving the same answer independently from applied weight window values.

Keywords: Monte Carlo Code, MCNP, Hard Biasing, Weight Window, Large Underestimation, Lost Particles

⁺ Center for Computational Science & e-Systems

* SGI Japan, Ltd.

強バイアス法を用いたモンテカルロ法計算に於ける過小評価問題の解決 —CAD/MCNP 自動変換プログラムで作成した入力の場合—

日本原子力研究開発機構 核融合研究開発部門

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飯田 浩正、川崎 信夫*、今野 力、佐藤 聡、関 暁之+

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ITER の R&D タスクとして行なっている「CAD/MCNP 自動変換コードの開発」において、ベンチマーク問題の解析中、適用「weight window」の違いによって MCNP が異なる答えを出すという不都合な事例に遭遇した。「weight window」法を含む“biasing”は計算速度を上げることがあっても、異なる答えを出す様なことがあってはならない。本研究では、この「大きな過小評価」が起こるメカニズムを明らかにしプログラムの修正を行ったので報告する。「大きな過小評価」は、以下の 2 つの事実の組合せでおこる：

- * MCNP はあるヒストリーの演算中に“lost particle”を検出すると当該ヒストリー中に計算された全てのタリーをキャンセルしてしまう。また、その時点で splitting 等の結果バンクに蓄積されていた粒子はその後追跡されることはない。

- * 微小形状エラーが入力に存在するとき、強バイアスの場合、“lost particle”を生じる確率はヒストリーの重要度に大きく左右される。

この結果、MCNP は選択的に重要度の高いヒストリーをキャンセルする事になる。

上記問題の解決を図るため、MCNP のサブルーチンのひとつである“hstory”の修正を行なった。テスト計算の結果、プログラムの修正は適切に行なわれ、MCNP は適用「weight window」に左右されず同じ答えを出すようになったことが確認された。

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1. Introduction

Monte Carlo codes, such as MCNP [1,2], are powerful tool for nuclear analyses of devices with complicated geometry. When device size is very large and significant attenuation of radiation fluxes is expected, it is not easy to obtain results with good statistics even with using large supercomputers. In such cases we very often use biasing techniques, such as “weight windows”. Those biasing can enhance calculation speed (can give smaller fractional standard deviation with the same number of histories or CPU time), but should not give different answers (mean values).

Because of severely limited space available for radiation shield in and around the torus, very high accuracy is required for nuclear analysis of ITER machine. In order to minimize estimation uncertainty in ITER device design calculation, most detailed and accurate geometry modeling is essential and development of CAD/MCNP interface program has been conducted under the ITER R&D task. Four parties (China, EU, Japan and USA) participate in this task and outcomes of their activities have been discussed periodically at the ITER task meetings.

In the course of analyzing the ITER benchmark problem defined in this task, we had an inconvenient experience in which different answers were obtained depending on applied weight window values. In this report we describe the problem we encountered (Chapter 2), causes of the problem (Chapter 3) and method to avoid such inconvenience (Chapter 4). Resolving the problem requires modification of one of the subroutines of MCNP program. The algorithm modification was made for MCNP5[1], although MCNP4C[2] was used in initial analysis of ITER CAD/MCNP benchmark problem. Those two versions showed the same behavior concerning with the present problem as shown in Appendix A.

2. Large underestimation encountered in the ITER CAD/MCNP benchmark neutron flux calculation

Figure 2.1 shows a 3-D CAD drawing of the ITER benchmark problem indicating four calculation tasks with MCNP. An inconvenience was observed in the calculation of the last one ④, which requires flux calculation in seven spheres located behind the port plug. Flux level behind the port plug is five orders of magnitude lower than that at the plasma region and most of neutrons which appear in the spaces behind the plug have reached there passing through the gaps between the plug and the port wall.

Comparison between Japanese and Chinese results [3] is shown in Fig. 2.2. Neutron fluxes in the plug region by the both parties agree very well, but those behind the plug have systematic difference between the both. Japanese calculation shows lower fluxes than Chinese one by a factor of about two. Neutron fluxes in the plug region and those in the gaps between the plug and port wall are shown in Fig. 2.3. This figure indicates that calculated fluxes in the plug by the both are similar but those in the gaps give quite different flux distributions.

By using MCNP plotting function, the gap geometries and surrounding material definitions were carefully compared between automatically converted geometry input data of the both parties. The both parties used their own CAD/MCNP conversion programs, namely MCAM [4] by China and GEOMIT [5] by Japan. However, no significant difference was found between the both geometry data.

In the following discussion, we mainly discuss about neutron fluxes in the first (closest to plasma) sphere (“#1 sphere” here after) out of the seven spheres behind the port plug in the equatorial port, since the difference between the both is systematic as shown in Fig. 2.2.

Figure 2.4 shows behaviors of the calculated fluxes in the #1 sphere by using the both input data changing width of weight windows (WUPN). There are two Japanese input data with different tables of weight windows (lower bound values); one is made for the target of photon flux in the sphere (ww1) and the other is for neutron flux in the sphere (ww2). In MCNP, lower bound of weight window of each cell is given as a table (WWN) and upper bound of the window is defined as $WWN \times WUPN$, where WUPN value is separately given in a WWP card (default is $WUPN=5$). When we increase the width (WUPN), particles experience less splitting, becoming closer to non-biasing (analog) run. By observing this figure, one can imagine that Japanese calculations have some problem since their results change depending on the value of the weight window width and they agree with Chinese ones when the window width becomes sufficiently large. Especially from the fact, that Japanese calculation gives different neutron flux by just changing “mode” (neutron and photon calculation or that of neutron only), we are sure that something very inconvenient has happened in Japanese calculation.

Prominent difference in the input data by Japanese conversion code (GEOMIT) and by Chinese (MCAM) is that no lost particle is observed in the run with the latter input data while some lost particles with the former input data. Although the frequency is not high (about one lost particle per a ten thousand of histories), this is suspected to be a reason of the large underestimation in Japanese calculation with hard biasing. The explanation of the mechanism of this underestimation is given in the next chapter.

3. Mechanism of the large underestimation with hard biasing technique

3.1 Micro geometry errors in MCNP input data automatically converted from CAD files

Depending on maturity of interface programs developed in the four parties (China, EU, Japan and U.S.A), there exist “micro geometry errors” in input data automatically converted with them. An example found in the input data converted by the Japanese conversion program GEOMIT is shown in Fig. 3.1. The size of this example of error is very small (see Fig. 3.1 (e)). It is not easy to find this kind of error visually by using MCNP built-in plotting function but their locations can be found in the MCNP output list as a part of data concerning “lost particle”. Because of the small error size, they usually do not affect calculation results, as far as hard biasing is not employed. However, they can give significant effect when hard biasing is used by the unique procedure of handling “lost particles” in MCNP. The next section explains this procedure.

3.2 The procedure of handling “lost particles” in MCNP

When a particle gets into “geometry error” like “undefined space”, MCNP has a trouble in continuing tracing particle trajectory and clears all tallies calculated during random walks of all particles in this history (like Fig. 3.2 (b)), not only giving up its tracing and clearing tallies due to the specific “lost particle” (see Fig. 3.2 (a)). There still should be many particles in the bank in hard biasing cases when lost particle is detected, since splitting could have happened many times especially when this history provided particles in “very important locations (or cells)” for the calculation purpose of the present problem. Those particles banked during this history will be thrown away (never be traced, more exactly). Then, MCNP starts tracking the next history, without recording any tallies in this troubled history. This algorism can be confirmed by examining the subroutine “hstory” or conducting a simple test calculation which is shown in Appendix B.

This handling procedure of lost particles is not convenient for us when we have many scattered micro geometry errors in our input data. The probability with which the particle get into one of the micro errors can be proportional to the particle population. In hard biasing cases we make weight of particle very low and increase number of particles (or population) leading to high probability of “got lost”, when it comes to important places (cells).

The combination of the above two facts leads to an inconvenient result of procedure, in which MCNP selectively cancels tallies of important histories, while it keeps all other histories as they are. This mechanism in hard biasing cases can lead to a large underestimation of calculation results.

4. Solution of the problem and modification of MCNP program

In order to solve the above problem, it is necessary to keep all histories in the calculation result without canceling their tallies even if lost particles are detected during random walks of those histories. Then MCNP may accept existence of micro geometry errors even in hard biasing cases. It can be a reasonable solution for this problem since it is not easy work to fully improve our CAD/MCNP automatic conversion program to completely eliminate micro geometry errors from its converted data.

In attempt of this solution, we slightly modified algorithm of “hstory” in MCNP5 program. Actually, modification was made in the following two steps.

(1) First step modification : **mod1**

We inserted a few lines into the subroutine hstory program in order to

- * avoid clearing calculated tallies during the present history, but
- * go back to “bankit(100)” to take next particle stored in the bank

Continuation of tracking the “lost particle” itself is given up.

(2) Second steps : **mod2-1 and mod2-2**

In addition to the above (1), continuation of tracking the “lost particle” was made in two ways, slightly moving particle location forward in the flight direction of the present particle.

- * move back to the proper location in the program to continue tracking of the particle.(mod2-1)
- * stop tracking the present particle trajectory but bank a new particle which has the same weight, flight direction, energy, etc. as the lost particle, but slightly moved location forward. (mod2-2)

A list of lines modified in the subroutine “hstory.F90” is attached in Appendix C.

5. Test calculations

With using the modified subroutine “hstory”, re-run of the ITER benchmark calculation has been conducted. Input data are the same as used in Fig. 2.2, which were automatically converted with GEOMIT from CAD data file. Test calculation results are shown in Fig.5.1, which compares neutron fluxes in #1 sphere by the MCNP with/without modified subroutine “hstory”. All results with the modified subroutine, agree with Chinese result not depending on weight window and/or calculation mode (include photon calculation or not), while MCNP with original “hstory” gives very different answers depending on weight window and calculation mode. Since all three modifications give practically the same answer, tracking the “lost particle” itself looks not important, at least for the present problem. Concerning with Chinese input data, we have no “lost particle” and the same results were obtained by the both of original and modified MCNP versions.

Figure 5.2 shows dependence of weight window width on the neutron flux tested with using mod2-2 “hstory.F90” subroutine. Now we have the same answer independently from WUPN value as shown in this Fig. 5.2. Comparison with Chinese results is again shown in Figs. 5.3 and 5.4, replacing Japanese results with the new ones, which are obtained with modified MCNP5 (mod2-2). Fluxes in the gaps between port wall and the plug calculate by the both parties agreed very well, as well as those behind the port plug.

Figures 5.5 and 5.6 show similar comparison but with other participating countries; EU and USA[3]. From theses figures we can say that agreement of the results from all participating parties’ CAD/MCNP interface programs on ④ task in Fig.2.1 has become rather well, leaving only 20 % difference among them.

6. Concluding remarks

The following conclusions have been obtained through the present study;

(1) Mechanism of large underestimation in cases of hard biasing was clarified. It is caused by the combination of the following two facts;

- * MCNP cancels all tallies calculated during the history, when one of particles in the history gets lost.

- * When we have distributed micro geometry errors, for example ITER 3-D model converted with our CAD/MCNP interface program (GEOMIT) from CAD drawing, important histories will have “lost particle” with high probability in the case of hard biasing.

These facts lead to selective canceling of important histories.

(2) The attempt to eliminate the large underestimation in case of hard biasing has been made successfully, by modifying subroutine “hstory.F90” of MCNP5. Three modified versions of the “hstory.F90” have been produced and they give practically the same answer, which is judged to be reasonable.

(3) Comparison of the fluxes behind the port plug with other parties (CN, EU and USA) showed rather good agreement, indicating that the CAD/MCNP interface programs developed by all parties are approaching a matured level for using in real design analysis of the machines.

(4) All calculations of the ITER CAD/MCNP benchmark problem, conducted by MCNP4C were re-conducted by using MCNP5, and it is confirmed that the both version give same answers.

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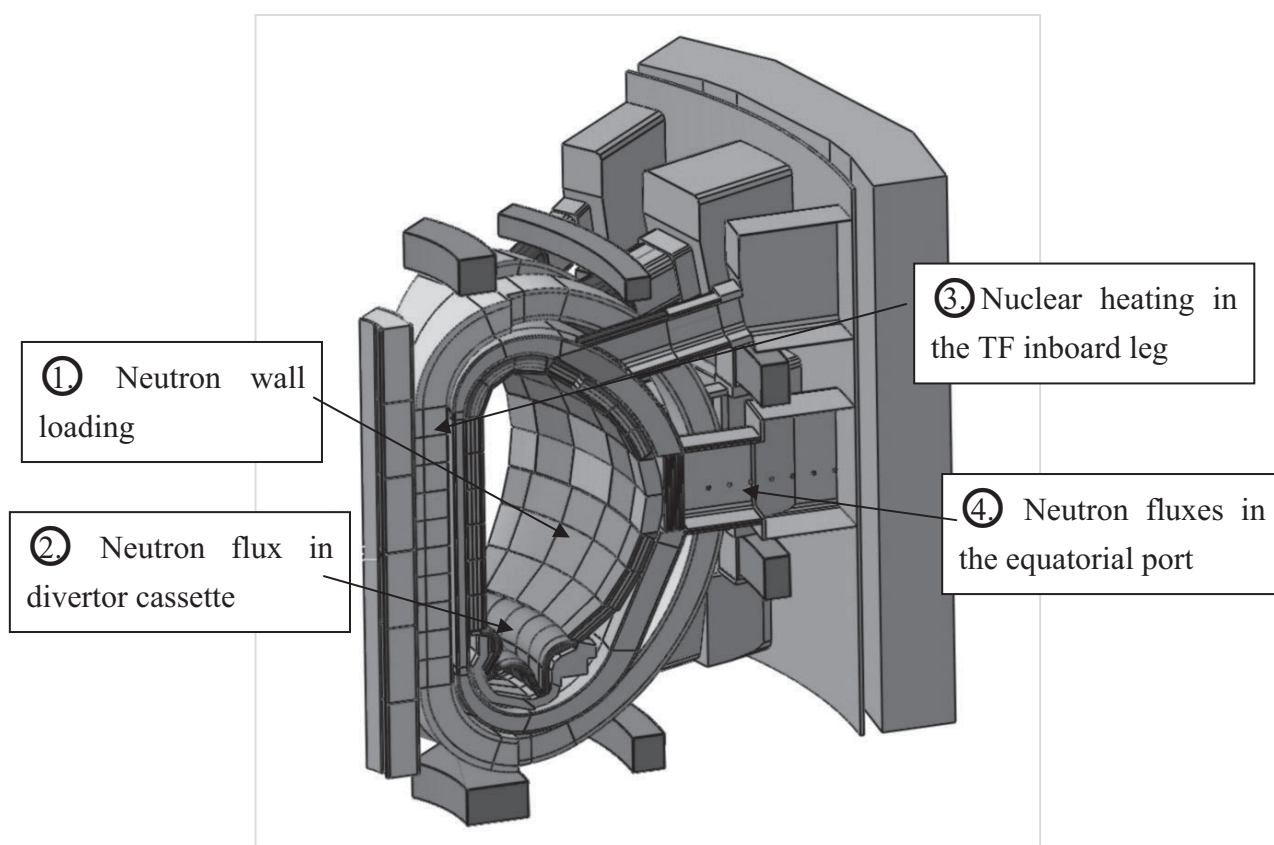


Fig. 2.1 CAD model of the ITER benchmark geometry and MCNP calculation tasks

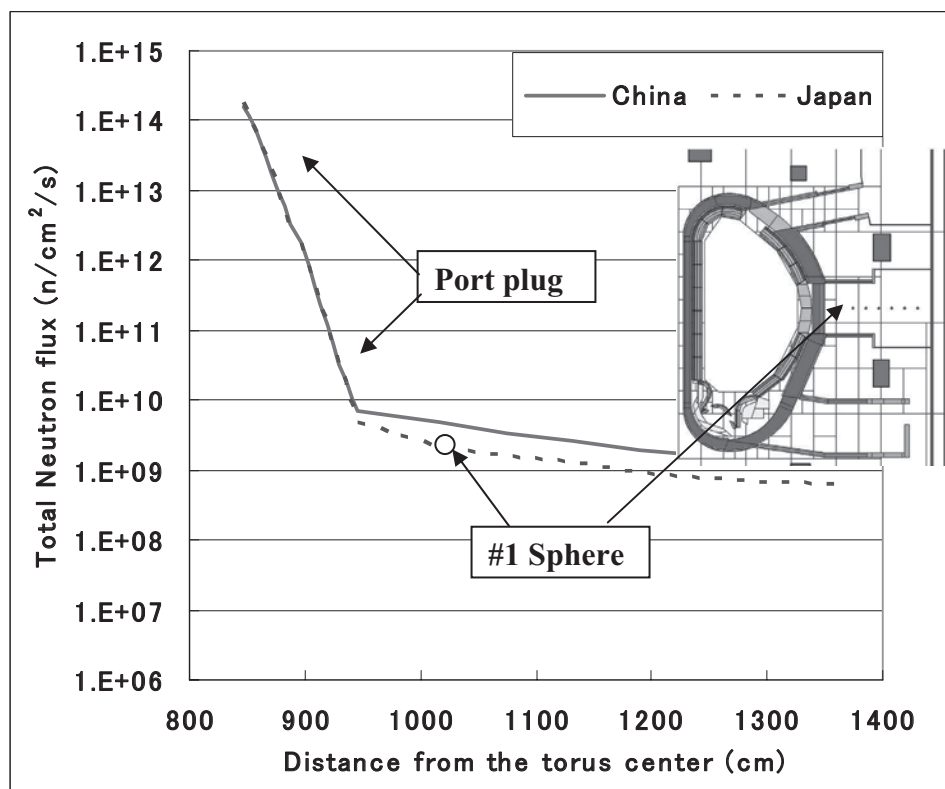


Fig. 2.2 Comparison of neutron fluxes in and behind the port plug between Japanese and Chinese calculations (with MCNP4C): ④ task shown in Fig. 2.1

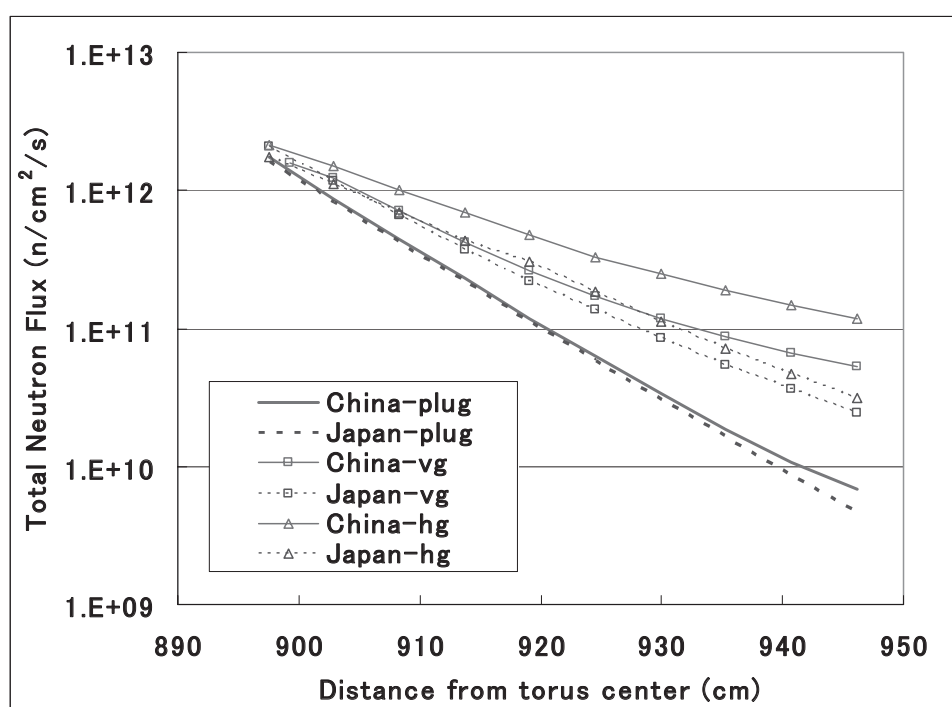


Fig. 2.3 Comparison of neutron fluxes in the port plug (plug), those in the vertical gaps (vg) and horizontal gaps (hg) between Japanese and Chinese calculations (with MCNP4C)

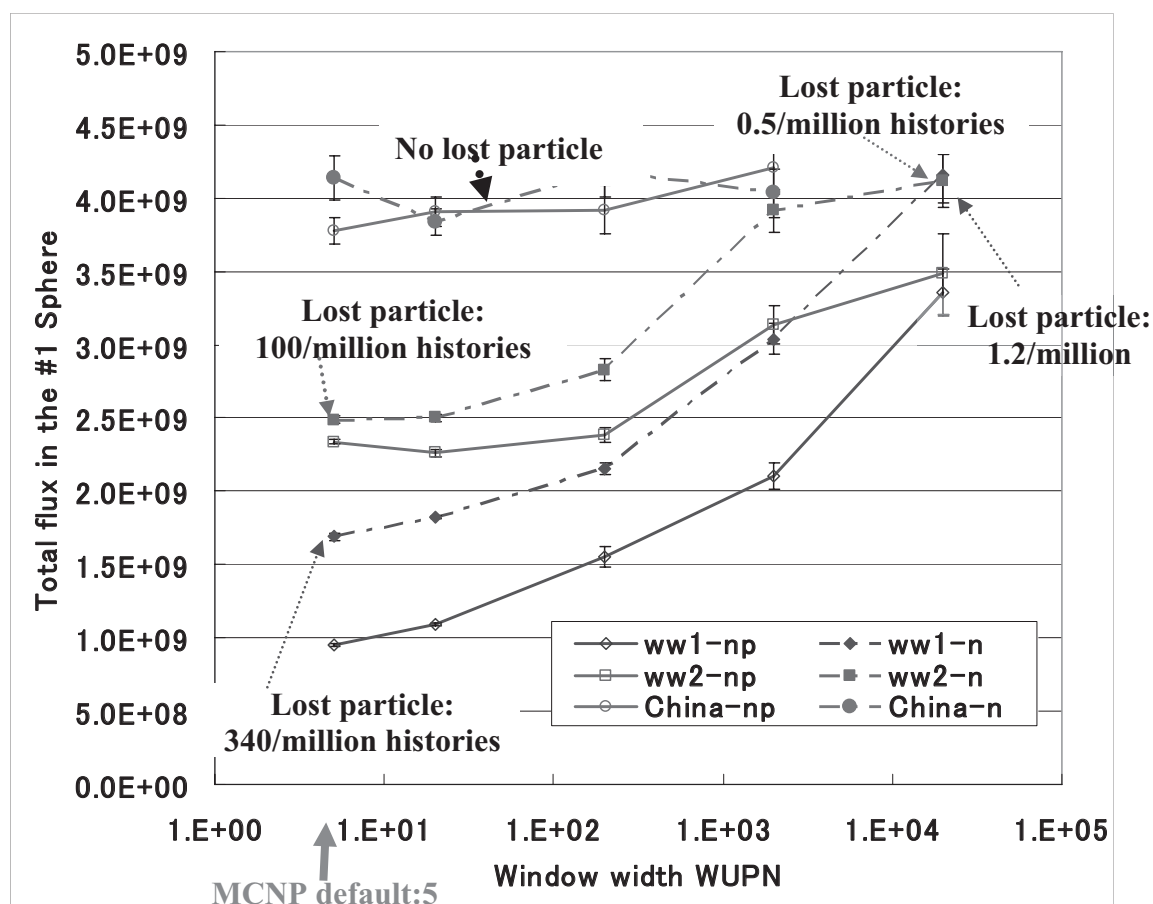


Fig. 2.4 Effect of weight window width (WUPN) on the neutron flux at the #1 sphere behind the port plug (with MCNP4C)

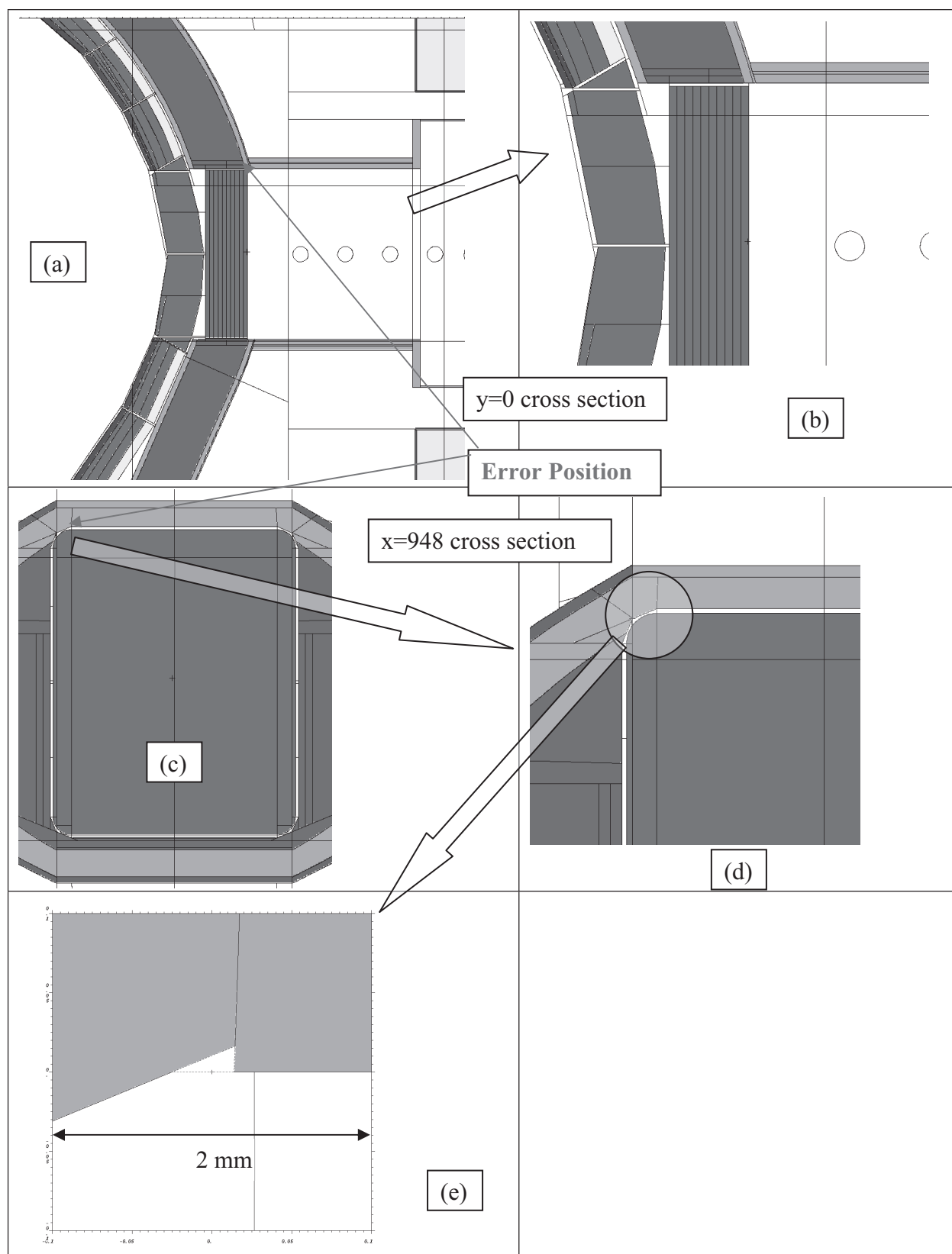


Fig. 3.1 An example of micro geometry errors in the input data automatically converted with GEOMIT from CAD data file

One History

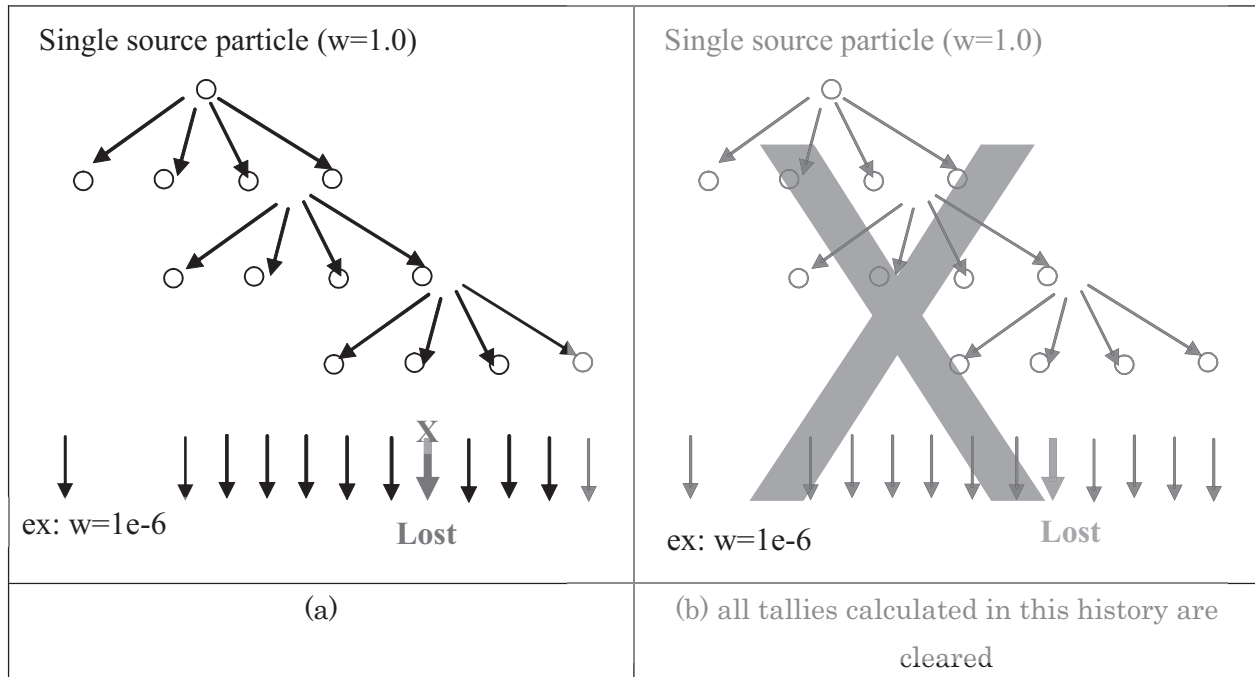


Fig. 3.2 Schematic explanation of MCNP lost particle handling procedure

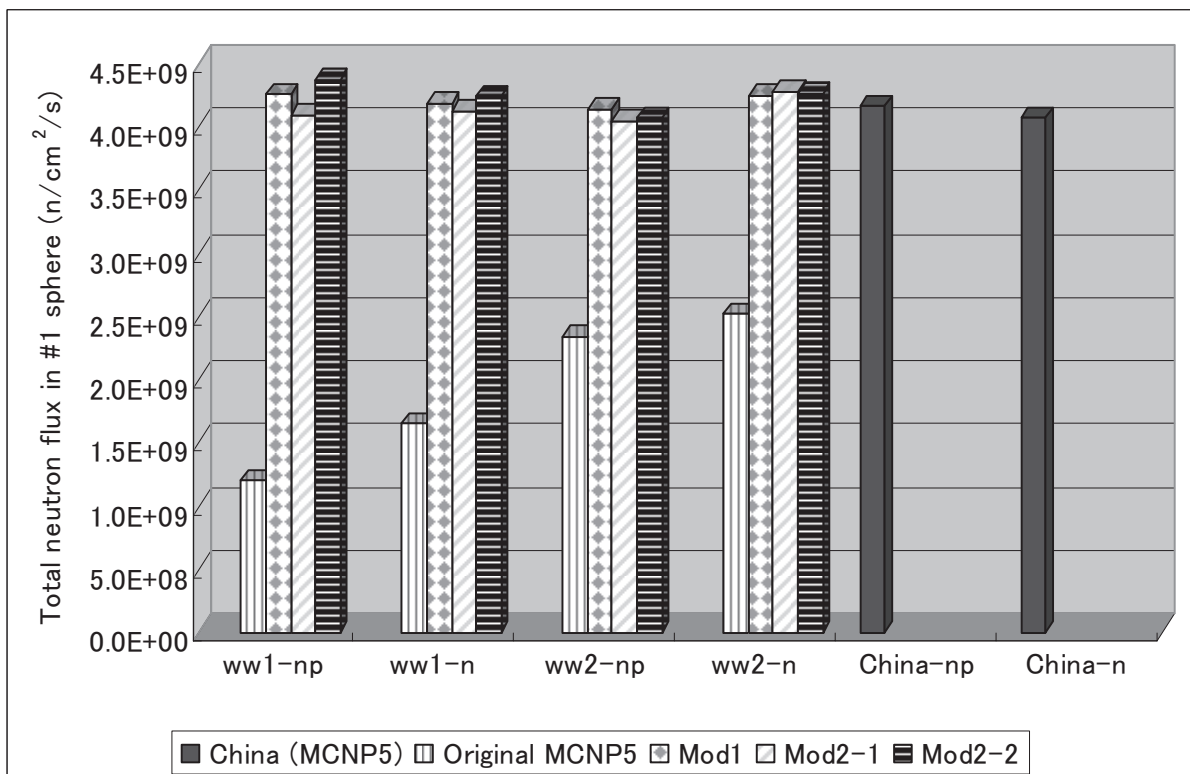


Fig. 5.1 Neutron fluxes in the #1 sphere behind the port plug of the ITER benchmark problem with modified subroutines "hstory.F90" (WUPN value: default of 5)

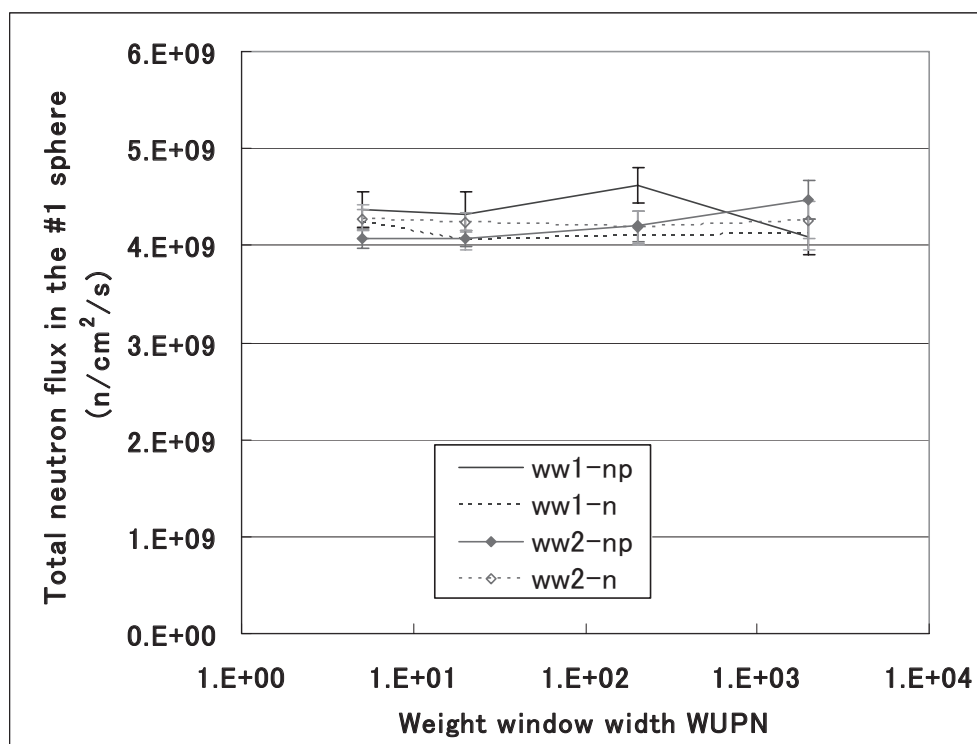


Fig. 5.2 Effect of weight window width (WUPN) on the neutron flux at the #1 sphere behind the port plug with modified MCNP5: (mod2-2)

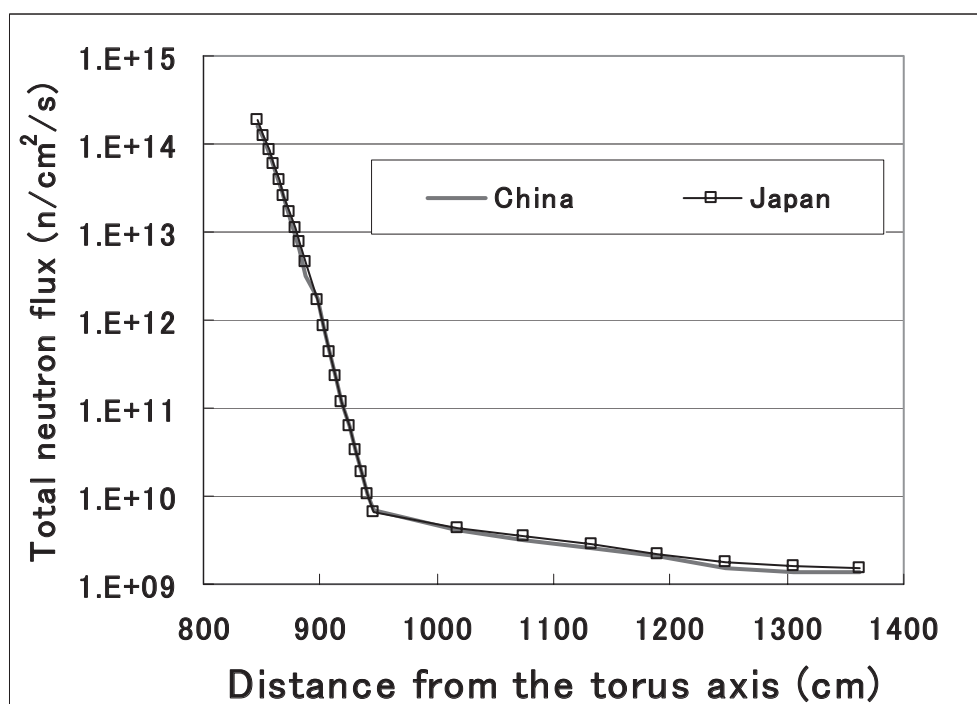


Fig. 5.3 Comparison of neutron fluxes in and behind the port plug between Japanese (with modified MCNPv5: mod2-2) and Chinese calculations (④ task shown in Fig.2.1)

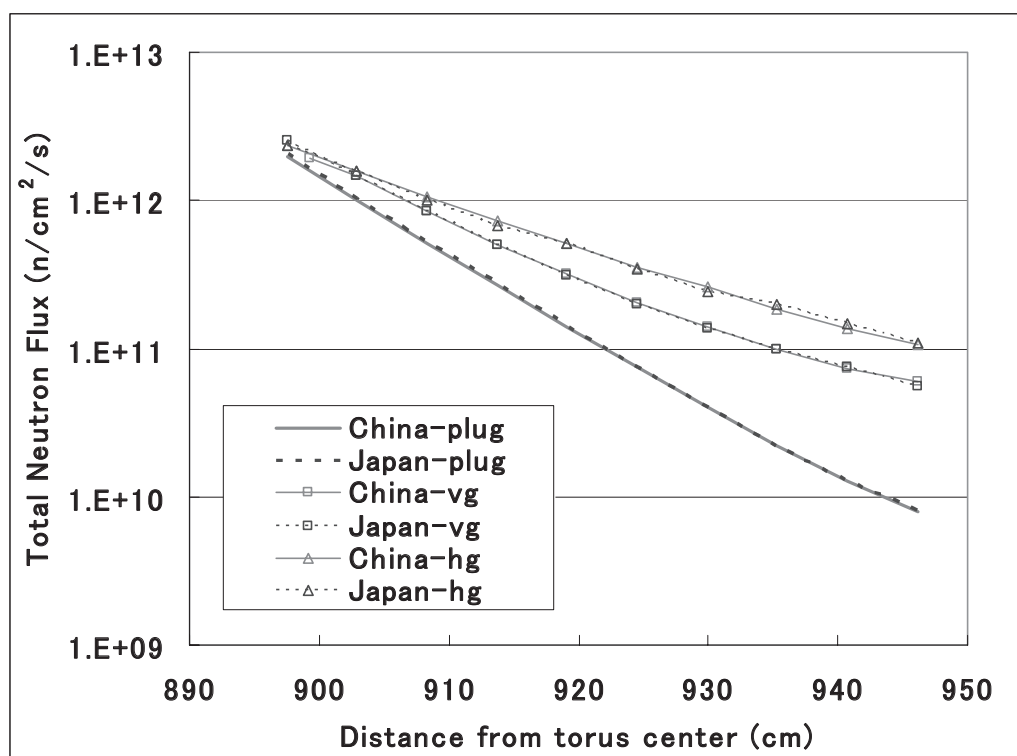


Fig. 5.4 Comparison of neutron fluxes in the plug(plug), those in the vertical gaps (vg) and horizontal gaps (hg) between Japanese (with the modified MCNP5:mod2-2) and Chinese calculations

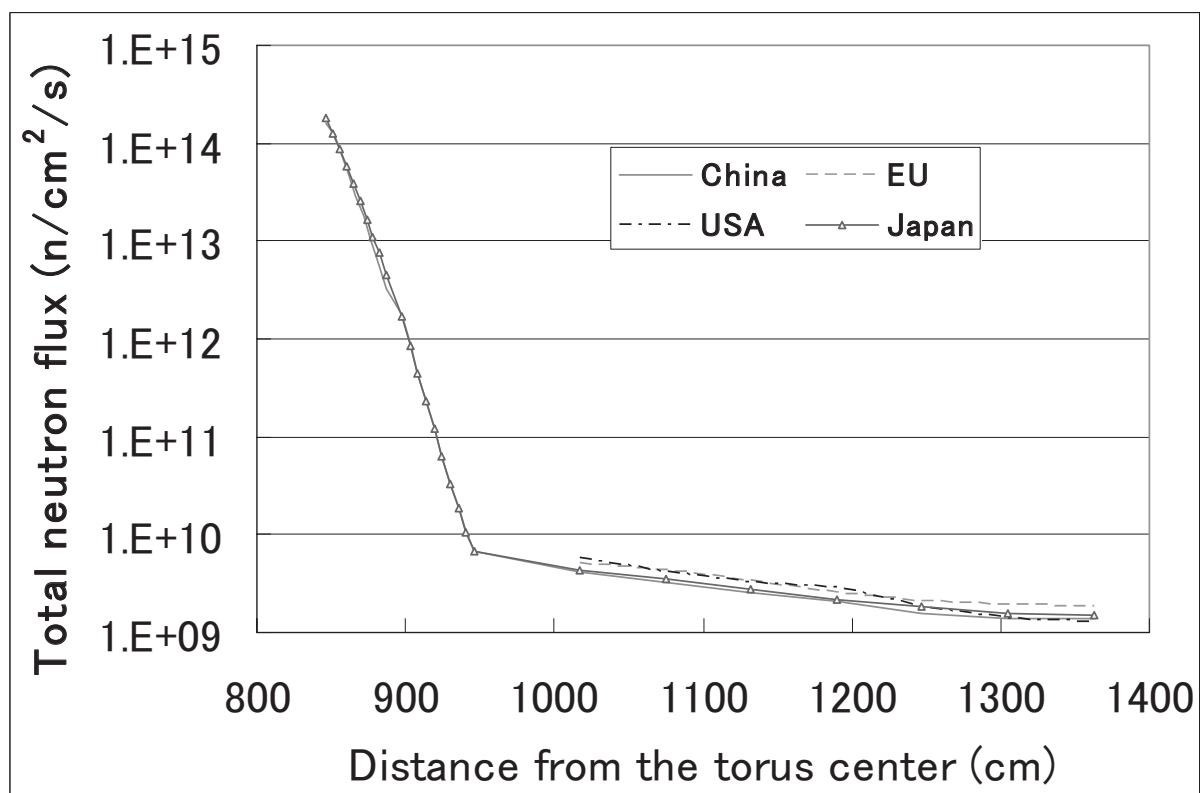


Fig. 5.5 Comparison of neutron fluxes in and behind the port plug among all participating parties in the CAD/MCNP benchmark tasks (④ task shown in Fig.2.1)

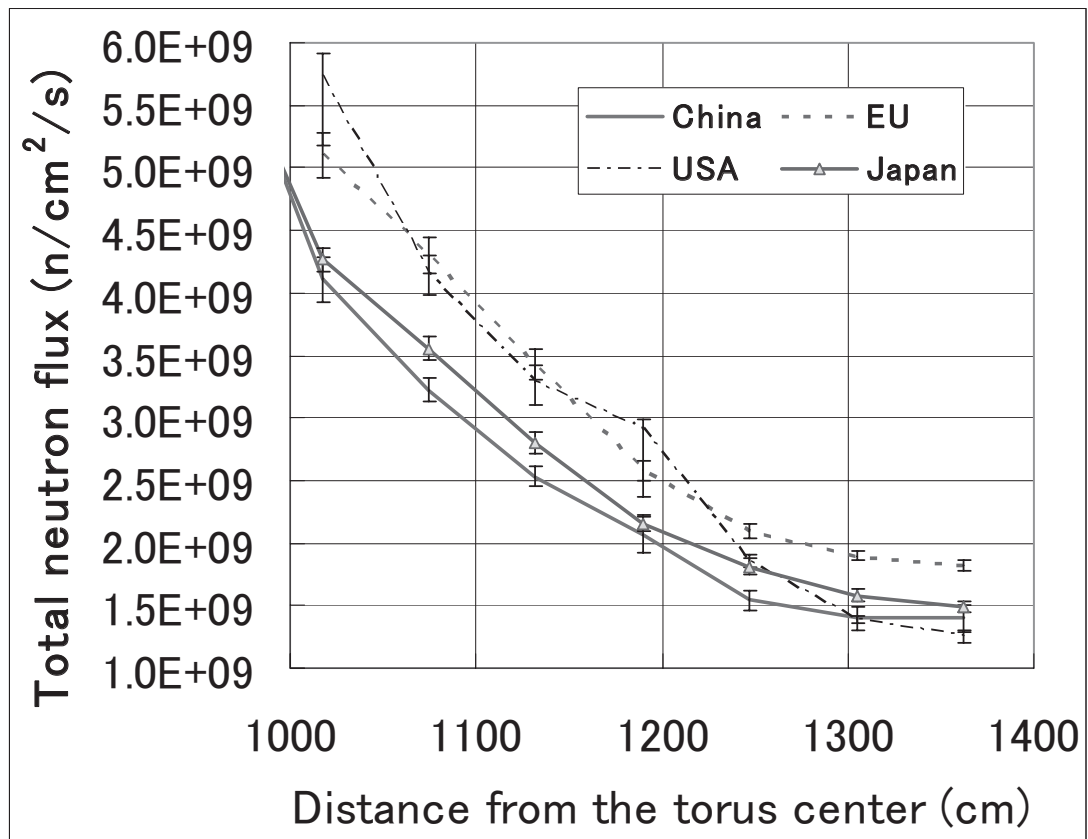


Fig. 5.6 Comparison of neutron fluxes behind the plug among all participating parties in the CAD/MCNP benchmark tasks (④ task shown in Fig.2.1)

Appendix A Comparison of neutron fluxes behind the plug between MCNP4C and MCNP5

The underestimation problem was detected in calculations with MCNP4C and we modified MCNP5 program. Then, in order to assure that the both version of MCNP gives same results, comparative study has been done. The results are shown in Fig. A-1. It can be said that we have almost complete agreement between results of the both versions of MCNP.

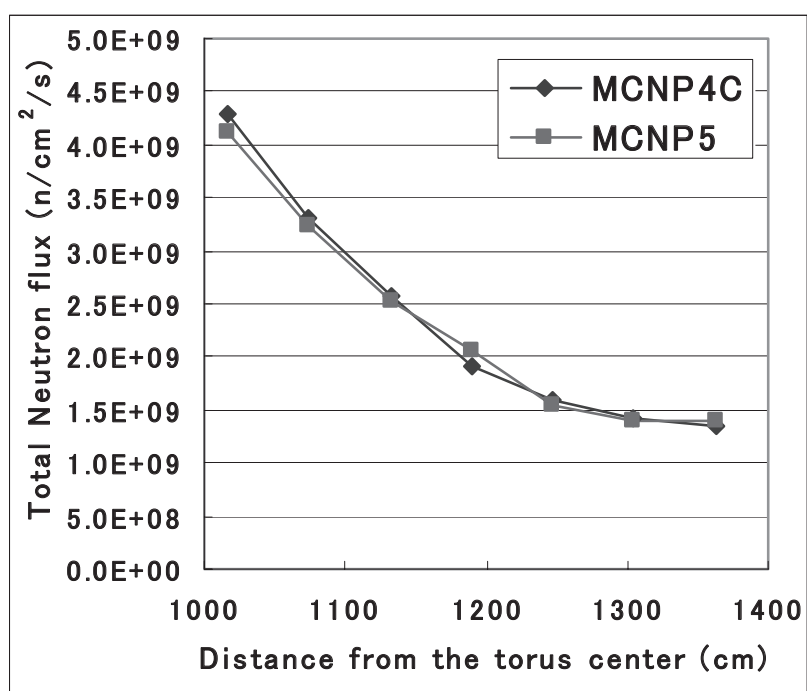


Fig. A-1 Comparison of neutron fluxes behind the plug between those calculated with MCNP4C and MCNP5 (Chinese input data is employed)

Appendix B Test calculation for confirming “lost particle handling” in MCNP

In order to visually understand how MCNP behaves when it gets “lost particle”, a test run has been made as shown in Fig.B-1. The calculation model is an empty box of 2m x 2m x 2m, with a 2 cm thick thin region (from $x=-1$ to 1) of source cell. Source particles born homogenously in this region and travel in parallel with X axis (40% in - direction and other 60% in + direction).

At $x=-50$ a very thin region (1 mm) is located where calculation condition is changed as follows;

- (a) Case 1 : normal void region with importance=1
- (b) Case 2 : normal void region with importance=0, then particles are killed when they have reached here.
- (c) Case 3 : this region is “undefined region” in MCNP input data; namely “geometry error”.

When particles reach this region, they are got lost.

Test calculation results are shown in Fig.B-2. When the geometry is correctly produced (case 1) the result looks very reasonable and ratio of flux levels on $x>0$ side and $x<0$ side is 6 : 4 corresponding to emission rates of source particle in both directions. When importance in the thin region at $x=-50$ cm is set to be 0 (case 2), the flux at $x < -50$ cm become zero, but no change in other location. If we have geometry error (undefined region) at $x= -50$ cm, all fluxes at $x < 0$ become zero although particles should have been traced normally until reaching the thin region at $x=-50$ cm. This shows clearly that MCNP cancels all tallies it calculated during the present history.

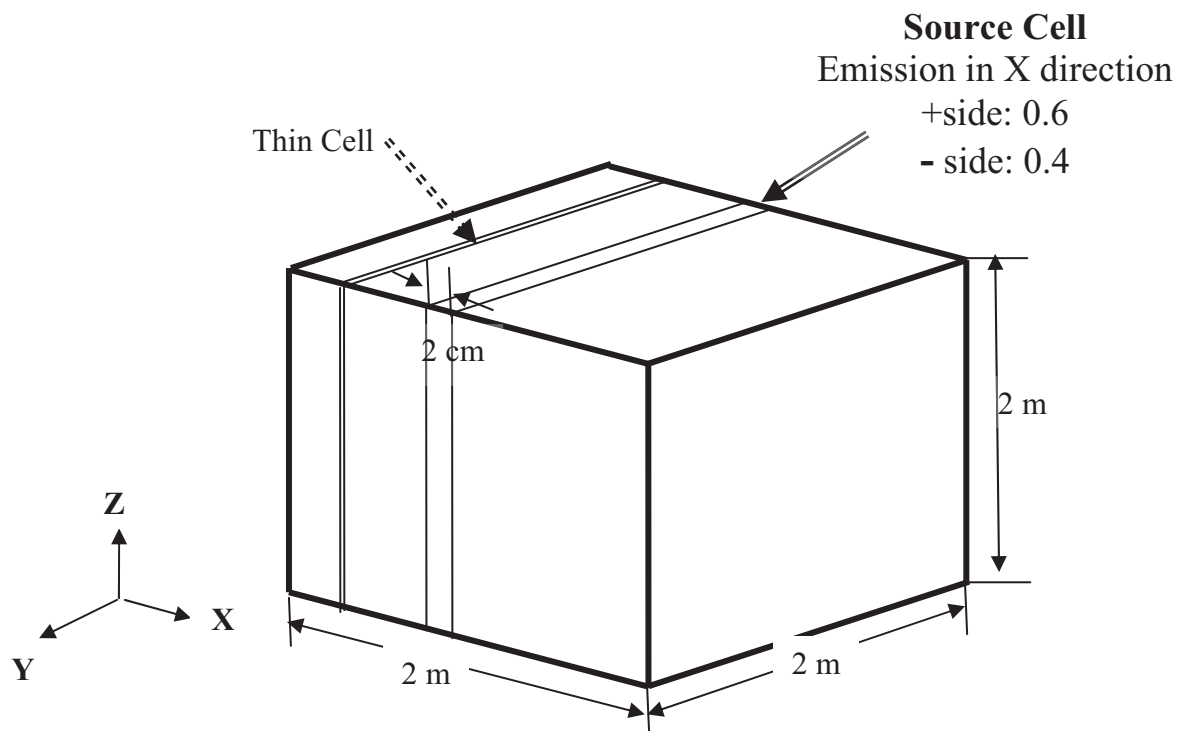


Fig. B-1 Test calculation model for examining how “lost particles” are treated in MCNP

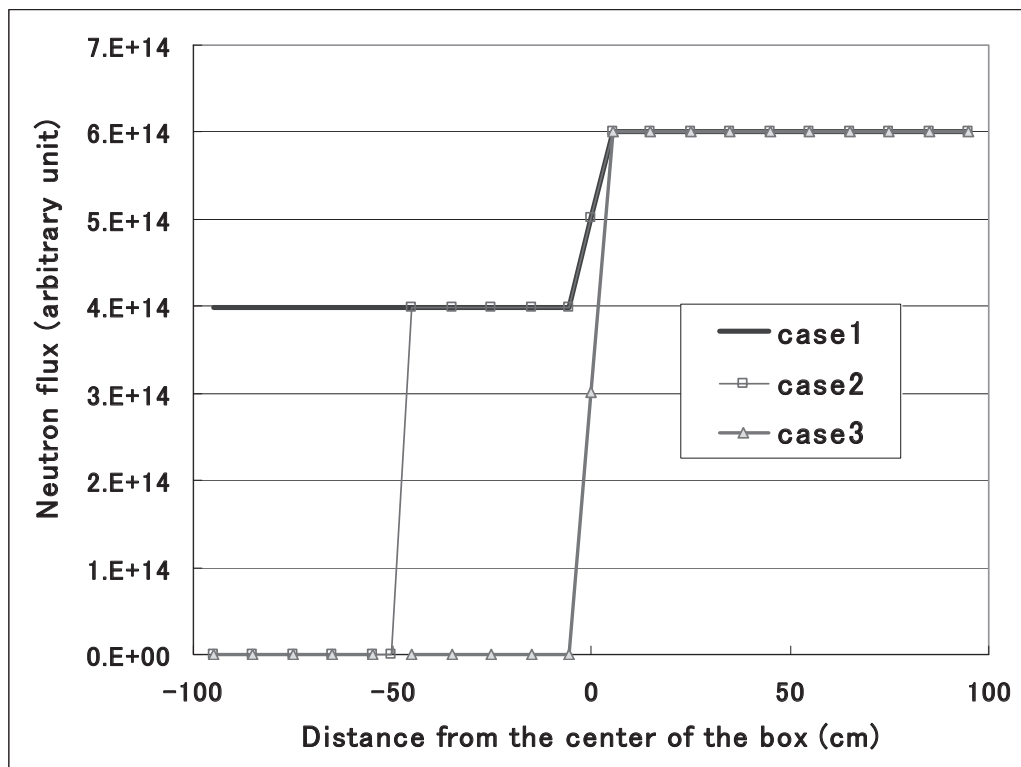


Fig. B-2 Results of the test calculation for examining how “lost particles” are treated in MCNP

Appendix C Modified part of subroutine “hstory.F90” of MCNP5

The modified parts of the subroutine “hstory.F90” of MCNP5.1.40 are listed below for the cases of (1)mod1, (2)mod2-1 and (3)mod2-2.

(1) Modification of “hstory.F90” for mod1

541a542,543

```
>      kdb = 0
>      go to 260
```

(2) Modification of “hstory.F90” for mod2-1

1,3d0

```
< !+ $Id: hstory.F90,v 1.8.2.1 2005/10/04 22:46:16 jsweezy Exp $
```

```
< ! Copyright LANL/UC/DOE - see file COPYRIGHT_INFO
```

```
<
```

6a4

```
>      !* hstory.F90.Low_ok *
```

15,17d12

```
< #if defined(VISED) && defined(DEC)
```

```
<      use c_interfaces
```

```
< #endif
```

23a19,28

```
> !*Lower*B
```

```
>      integer :: lgcwk(mlgc+1)
```

```
>      integer :: ncel_pl(10)
```

```
>      integer :: loop_ctr
```

```
> !*Lower*B
```

```
>      integer      jerr                      !*Debug*
```

```
>      jerr=0                                !*Debug*
```

```
> !*Lower*B
```

```
>      loop_ctr=0
```

```
> !*Lower*B
```

42a48,50

```
> !  if(nps.eq.lost_num) then                      !*Debug*
```

```
> !  write(6,*) 'hstory-R: Label_020',nps,icl !*Debug*
```

```
> !  endif                                          !*Debug*
```

54c62,65


```

<  if( kdb/=0      )  go to 390
---
>  if( kdb/=0      )  then
>    jerr=1          !*Debug*
>    go to 390
>  endif
76c87,90
<    if( kdb/=0 )  go to 390
---
>    if( kdb/=0 )  then
>      jerr=2          !*Debug*
>      go to 390
>    endif
94c108,111
<  if( kdb/=0 )  go to 390
---
>  if( kdb/=0 )  then
>    jerr=3          !*Debug*
>    go to 390
>  endif
152c169,172
<    if( kdb/=0  )  go to 390
---
>    if( kdb/=0  )  then
>      jerr=4          !*Debug*
>      go to 390
>    endif
231c251,254
<  if( kdb/=0 )  go to 390
---
>  if( kdb/=0 )  then
>    jerr=5          !*Debug*
>    go to 390
>  endif
237c260,263
<    if( kdb/=0 )  go to 390
---
>    if( kdb/=0 )  then

```

```

>      jerr=6          !*Debug*
>      go to 390
>      endif
335c361,364
<      if( d==dls ) then
---
>      if( d==dls ) then
>      !*Lower*B
>      jrtry=0
>      !*Lower*E
340c369,469
<      if( kdb/=0      ) go to 390
---
>      if( kdb/=0      ) then
>      !*Lower*B
>      loop_ctr=loop_ctr+1
>      !*Debug*B
>      ! write(6,'(60hHISTORY-R: Aft_surfac,kdb/=0. nps,kdb,icl,jsu,loop_ctr,mynum=, &
>      !& i8,i4,2i8,2i4)') nps,kdb,icl,jsu,loop_ctr,mynum
>      !*Debug*E
>      if(loop_ctr.gt.100) then
>      jerr=71          !*Debug*
>      go to 390
>      endif
>      if(jrtry.eq.0) then
>      x_ww =xxx
>      y_ww =yyy
>      z_ww =zzz
>      u_ww =uuu
>      v_ww =vvv
>      w_ww =www
>      d_ww =d
>      iclwk=icl
>      jsuwk=jsu
>      kdbw =kdb
>      endif
>
>      !* shift (x,y,z), find the unique 'icl'

```

```

> 110 continue
>      jrtry=jrtry+1
>      if(jrtry.gt.10) go to 120
>      kdb = 0
>      xxx=xxx+u_ww*1.d+00
>      yyy=yyy+v_ww*1.d+00
>      zzz=zzz+w_ww*1.d+00
>      n_pl=0
>      do i=1,mtx
>          call chkcel(i,0,k)
>          if(k.eq.0) then
>              j2=abs(lca(i))
>              j3=abs(lca(i+1))
>              do j=1,j3-j2+1
>                  lgcwk(j)=lgc(j)
>              enddo
>              n_pl=n_pl+1
>              ncel_pl(n_pl)=i
>              go to 112
>          endif
>      enddo
> 112 continue
>      if(n_pl.ne.1) go to 110      !*retry 'shift (x,y,z)'
>      icl=ncel_pl(1)
>      ic9=icl
>
> !* change 'jsu'
>      j2=abs(lca(ic9))
>      j3=abs(lca(ic9+1))
>      do j=1,j3-j2+1
>          lgc(j)=lgcwk(j)
>      enddo
>      call track(ic9)
>      if(kdb.ne.0) go to 110      !*retry 'shift (x,y,z)'
> !
>      uuu=-u_ww
>      vvv=-v_ww
>      www=-w_ww

```

```

>      do j=1,j3-j2+1
>          lgc(j)=lgcwk(j)
>      enddo
>      call track(ic9)
> !*???*B
>      if(kdb.ne.0) then
>          uuu=u_ww
>          vvv=v_ww
>          www=w_ww
>          go to 110      !*retry 'shift (x,y,z)'
>      endif
> ! write(6,*) 'HISTORY-R: go to the next cell process.'
> !*???*E
>      jsu=jap
>      uuu=u_ww
>      vvv=v_ww
>      www=w_ww
>      do j=1,j3-j2+1
>          lgc(j)=lgcwk(j)
>      enddo
>      d  =d_ww
>      go to 50          !* go to the next cell process.
>
> 120 continue
>      kdb=kdbw
>      jsu=jsuwk
>      icl=iclwk
>      d  =d_ww
>      uuu=u_ww
>      vvv=v_ww
>      www=w_ww
>      xxx=x_ww
>      yyy=y_ww
>      zzz=z_ww
> !*Lower*E
>      jerr=7          !*Debug*
>      go to 390
>      endif

```

380c509,512

< if(kdb/=0) go to 390

> if(kdb/=0) then

> jerr=8 !*Debug*

> go to 390

> endif

391c523,526

< if(kdb/=0) go to 390

> if(kdb/=0) then

> jerr=9 !*Debug*

> go to 390

> endif

396c531,534

< if(kdb/=0) go to 390

> if(kdb/=0) then

> jerr=10 !*Debug*

> go to 390

> endif

503a642

> jerr=11 !*Debug*

541a681,690

> !*Debug*B

> write(6,'(53hHISTORY-R: L390. jerr,nps,icl,kdb,loop_ctr,mynum,nbnk=, &

> & i2,i10,i8,3i4,i8)') jerr,nps,icl,kdb,loop_ctr,mynum,nbnk

> !*Debug*E

> !

> !*Lower*B

> kdb=0

> loop_ctr=0

> go to 260 !* for unbanking.

> !*Lower*E

543d691

<

(3) Modification of "history.F90" for mod2-2

53a54

> if(kdb/=0) ipath =1

75a77

> if(kdb/=0) ipath =2

93a96

> if(kdb/=0) ipath =3

151c154,155

< call forcol

> call forcol

> if(kdb/=0) ipath =4

230a235

> if(kdb/=0) ipath =5

236a242

> if(kdb/=0) ipath =6

339a346

> if(kdb/=0) ipath =7

379a387

> if(kdb/=0) ipath =8

390a399

> if(kdb/=0) ipath =9

395a405

> if(kdb/=0) ipath =10

503a514

> ipath =11

541a553,636

> ! write(iuo,16) ipath,kdb,npa,mxa

> ! 16 format("ipath = ", i6, "kdb= ", i6, "npa = ", i6, "mx = ", i6)

> if(kdb<=0.or.kdb>=11) go to 580

> if(kdb==2 .or. kdb==4 .or. kdb==6) go to 580

> if(icl > mxa) go to 580

> ! determine whether the particle really is in cell icl.

> uuu = -uuu

> vvv = -vvv

> www = -www

> call chkcel(icl,0,j1)

> if(j1/=0) then

```

> do iq = 1, mxa
>   call chkcel(iq, 0, j2)
>   if( j2==0 ) exit
> enddo
> icl=iq
> endif
> uu = -uu
> vv = -vv
> ww = -ww
> do 261 iida=1, 10
>   aiida=iida
>   dskip=rdum(5)/5*aiida
> !   dskip=rdum(5)
> !   write(iuo, 14) xxx, yyy, zzz, uu, vv, ww, icl, iii, jjj, kkk, nbk
> ! 14   format( "yyy)", 6e12.5, 5i6)
>   xxx = xxx + uu*dskip
>   yyy = yyy + vv*dskip
>   zzz = zzz + ww*dskip
> do iq = 1, mxa
>   call chkcel(iq, 0, j2)
>   if( j2==0 ) exit
> enddo
> if( iq > mxa ) cycle
> if( iq==icl ) cycle
> ! function wwval(ny, nc, nb, na, ix)
>   nb=1
>   wwva=wwval(ipt, iq, nb, 0, 0)
>   if(wwva<0.0) then
>     xxx = xxx - uu*dskip
>     yyy = yyy - vv*dskip
>     zzz = zzz - ww*dskip
>     go to 580
>   endif
>   iid=icl
>   icl=iq
>   iicl=ncl(icl)
>   iiid=ncl(iid)
>

```

```

> !      write(iuo,542) lev,iiid,icl,iicl,iida,ipt,wwva,rdum(5),dskip,erg,wgt
> ! 542      format("OK!",6i6,6e9.2)
> !          write(iuo ,13) xxx,yyy,zzz,uuu,vvv,www,icl,iii,jjj,kkk
> ! 13      format( "xxx)", 6e12.5,4i7)
>          do j=0,lev
>              udt(1, j ) = xxx
>              udt(2, j ) = yyy
>              udt(3, j ) = zzz
>              udt(4, j ) = uuu
>              udt(5, j ) = vvv
>              udt(6, j ) = www
>              udt(7, j ) = icl
>              udt(8, j ) = iii
>              udt(9, j ) = jjj
>              udt(10, j ) = kkk
>          end do
> !          write(iuo ,11) iicl, iiid, lev, ji
> ! 11      format( "iicl,iiid,lev,ji,levchk", 5i6  )
> !          write(iuo ,12) (udt(i,lev),i=1,10),nbnk,iida
> ! 12      format( "udt(1-10,lev)", 6e12.5,4f7.1,2i6)
> !          kdb = 0
> !          if( npa == 0 ) npa = 1          !iida
>          npa = 1          !iida
>          if (ipt==1) call bankit(7)
>          if (ipt==2) call bankit(8)
>          kdb = 0
> !          write(iuo ,15) nbnk
> ! 15      format( "nbnk  =", i6)
>          go to 260
> 261      continue
> 580 continue
> !          write(iuo,541)  icl,iq,ncl(icl), xxx, yyy, zzz,kdb,wwva
> ! 541      format("580 icl,iq x,y,z kdb wwva",3i6,3e12.5,i6,e9.2)
>          kdb = 0
>          go to 260
~

```


国際単位系 (SI)

表 1. SI 基本単位

基本量	SI 基本単位	
	名称	記号
長さ	メートル	m
質量	キログラム	kg
時間	秒	s
電流	アンペア	A
熱力学温度	ケルビン	K
物質量	モル	mol
光度	カンデラ	cd

表 2. 基本単位を用いて表されるSI組立単位の例

組立量	SI 基本単位	
	名称	記号
面積	平方メートル	m ²
体積	立方メートル	m ³
速度	メートル毎秒	m/s
加速度	メートル毎秒毎秒	m/s ²
波数	毎メートル	m ⁻¹
密度 (質量密度)	キログラム毎立方メートル	kg/m ³
質量体積 (比体積)	立法メートル毎キログラム	m ³ /kg
電流密度	アンペア毎平方メートル	A/m ²
磁界の強さ	アンペア毎メートル	A/m
(物質量の) 濃度	モル毎立方メートル	mol/m ³
輝度	カンデラ毎平方メートル	cd/m ²
屈折率	(数の) 1	1

表 5. SI 接頭語

乗数	接頭語	記号	乗数	接頭語	記号
10 ²⁴	ヨタ	Y	10 ⁻¹	デシ	d
10 ²¹	ゼタ	Z	10 ⁻²	センチ	c
10 ¹⁸	エクサ	E	10 ⁻³	ミリ	m
10 ¹⁵	ペタ	P	10 ⁻⁶	マイクロ	μ
10 ¹²	テラ	T	10 ⁻⁹	ナノ	n
10 ⁹	ギガ	G	10 ⁻¹²	ピコ	p
10 ⁶	メガ	M	10 ⁻¹⁵	フェムト	f
10 ³	キロ	k	10 ⁻¹⁸	アト	a
10 ²	ヘクト	h	10 ⁻²¹	ゼプト	z
10 ¹	デカ	da	10 ⁻²⁴	ヨクト	y

表 3. 固有の名称とその独自の記号で表されるSI組立単位

組立量	SI 組立単位			
	名称	記号	他のSI単位による表し方	SI基本単位による表し方
平面角	ラジアン ^(a)	rad		m・m ⁻¹ =1 ^(b)
立体角	ステラジアン ^(a)	sr ^(c)		m ² ・m ⁻² =1 ^(b)
周波数	ヘルツ	Hz		s ⁻¹
力	ニュートン	N		m・kg・s ⁻²
圧力, 応力	パスカル	Pa	N/m ²	m ⁻¹ ・kg・s ⁻²
エネルギー, 仕事, 熱量	ジュール	J	N・m	m ² ・kg・s ⁻²
工率, 放射束	ワット	W	J/s	m ² ・kg・s ⁻³
電荷, 電気量	クーロン	C		s・A
電位差 (電圧), 起電力	ボルト	V	W/A	m ² ・kg・s ⁻³ ・A ⁻¹
静電容量	ファラド	F	C/V	m ⁻² ・kg ⁻¹ ・s ⁴ ・A ²
電気抵抗	オーム	Ω	V/A	m ² ・kg・s ⁻³ ・A ⁻²
コンダクタンス	ジーメン	S	A/V	m ⁻² ・kg ⁻¹ ・s ³ ・A ²
磁束	ウェーバ	Wb	V・s	m ² ・kg・s ⁻² ・A ⁻¹
磁束密度	テスラ	T	Wb/m ²	kg・s ⁻² ・A ⁻¹
インダクタンス	ヘンリー	H	Wb/A	m ² ・kg・s ⁻² ・A ⁻²
セルシウス温度	セルシウス度 ^(d)	°C		K
光束	ルーメン	lm	cd・sr ^(c)	m ² ・m ⁻² ・cd=cd
照射 (放射性核種の) 放射能	ベクレル	Bq	lm/m ²	m ² ・m ⁻⁴ ・cd=m ⁻² ・cd
吸収線量, 質量エネルギー分与, カーマ線量当量, 周辺線量当量, 方向性線量当量, 個人線量当量, 組織線量当	グレイ	Gy	J/kg	m ² ・s ⁻²
	シーベルト	Sv	J/kg	m ² ・s ⁻²

- (a) ラジアン及びステラジアンの使用は、同じ次元であっても異なった性質をもった量を区別するときの組立単位の表し方として利点がある。組立単位を形作るときいくつかの用例は表 4 に示されている。
- (b) 実際には、使用する時には記号rad及びsrが用いられるが、習慣として組立単位としての記号“1”は明示されない。
- (c) 測光学では、ステラジアンの名称と記号srを単位の表し方の中にそのまま維持している。
- (d) この単位は、例としてミリセルシウス度m°CのようにSI接頭語を併って用いても良い。

表 4. 単位の中に固有の名称とその独自の記号を含むSI組立単位の例

組立量	SI 組立単位		
	名称	記号	SI 基本単位による表し方
粘力のモーメント	パスカル秒	Pa・s	m ⁻¹ ・kg・s ⁻¹
表面張力	ニュートンメートル	N・m	m ² ・kg・s ⁻²
角速度	ニュートン毎メートル	N/m	kg・s ⁻²
角加速度	ラジアン毎秒	rad/s	m・m ⁻¹ ・s ⁻¹ =s ⁻¹
熱流密度, 放射照度	ラジアン毎平方秒	rad/s ²	m・m ⁻¹ ・s ⁻² =s ⁻²
熱容量, エントロピー	ワット毎平方メートル	W/m ²	kg・s ⁻³
質量熱容量 (比熱容量), 質量エントロピー	ジュール毎平方メートル	J/K	m ² ・kg・s ⁻² ・K ⁻¹
質量エネルギー (比エネルギー)	ジュール毎キログラム	J/(kg・K)	m ² ・s ⁻² ・K ⁻¹
熱伝導率	ジュール毎キログラム毎平方メートル	J/(m ² ・K)	m ² ・s ⁻² ・K ⁻¹
体積エネルギー	ワット毎メートル毎立方メートル	W/(m・K)	m・kg・s ⁻³ ・K ⁻¹
電界の強さ	ジュール毎立方メートル	J/m ³	m ⁻¹ ・kg・s ⁻²
体積電荷	ボルト毎メートル	V/m	m・kg・s ⁻³ ・A ⁻¹
電気変位	クーロン毎立方メートル	C/m ³	m ⁻³ ・s・A
誘電率	クーロン毎平方メートル	C/m ²	m ⁻² ・s・A
透磁率	ファラド毎メートル	F/m	m ⁻³ ・kg ⁻¹ ・s ⁴ ・A ²
モルエネルギー	ヘンリー毎メートル	H/m	m・kg・s ⁻² ・A ⁻²
モルエントロピー, モル熱容量	ジュール毎モル	J/mol	m ² ・kg・s ⁻² ・mol ⁻¹
照射線量 (X線及びγ線)	ジュール毎モル毎ケルビン	J/(mol・K)	m ² ・kg・s ⁻² ・K ⁻¹ ・mol ⁻¹
吸収線量	クーロン毎キログラム	C/kg	kg ⁻¹ ・s・A
放射強度	グレイ毎秒	Gy/s	m ² ・s ⁻³
放射輝度	ワット毎ステラジアン	W/sr	m ⁴ ・m ⁻² ・kg・s ⁻³ =m ² ・kg・s ⁻³
	ワット毎平方メートル毎ステラジアン	W/(m ² ・sr)	m ² ・m ⁻² ・kg・s ⁻³ =kg・s ⁻³

表 6. 国際単位系と併用されるが国際単位系に属さない単位

名称	記号	SI 単位による値
分	min	1 min=60s
時	h	1h =60 min=3600 s
日	d	1 d=24 h=86400 s
度	°	1° =(π/180) rad
分	′	1′ =(1/60)° =(π/10800) rad
秒	″	1″ =(1/60)′ =(π/648000) rad
リットル	l, L	1l=1 dm ³ =10 ⁻³ m ³
トン	t	1t=10 ³ kg
ネーパ	Np	1Np=1
ベル	B	1B=(1/2) ln10 (Np)

表 7. 国際単位系と併用されこれに属さない単位で SI 単位で表される数値が実験的に得られるもの

名称	記号	SI 単位であらわされる数値
電子ボルト	eV	1eV=1. 60217733 (49) ×10 ⁻¹⁹ J
統一原子質量単位	u	1u=1. 6605402 (10) ×10 ⁻²⁷ kg
天文単位	ua	1ua=1. 49597870691 (30) ×10 ¹¹ m

表 8. 国際単位系に属さないが国際単位系と併用されるその他の単位

名称	記号	SI 単位であらわされる数値
海里		1 海里=1852m
ノット		1 ノット=1 海里毎時=(1852/3600) m/s
アール	a	1 a=1 dam ² =10 ² m ²
ヘクタール	ha	1 ha=1 hm ² =10 ⁴ m ²
バール	bar	1 bar=0. 1MPa=100kPa=1000hPa=10 ⁵ Pa
オングストローム	Å	1 Å=0. 1nm=10 ⁻¹⁰ m
バイン	b	1 b=100fm ² =10 ⁻²⁸ m ²

表 9. 固有の名称を含むCGS組立単位

名称	記号	SI 単位であらわされる数値
エルグ	erg	1 erg=10 ⁻⁷ J
ダイン	dyn	1 dyn=10 ⁻⁵ N
ポアズ	P	1 P=1 dyn・s/cm ² =0. 1Pa・s
ストークス	St	1 St =1cm ² /s=10 ⁻⁴ m ² /s
ガウス	G	1 G ≡10 ⁻⁴ T
エルステッド	Oe	1 Oe ≡(1000/4π) A/m
マクスウェル	Mx	1 Mx ≡10 ⁻⁸ Wb
スチルブ	sb	1 sb =1cd/cm ² =10 ⁴ cd/m ²
ホト	ph	1 ph=10 ⁴ lx
ガリ	Gal	1 Gal =1cm/s ² =10 ⁻² m/s ²

表10. 国際単位に属さないその他の単位の例

名称	記号	SI 単位であらわされる数値
キュリー	Ci	1 Ci=3. 7×10 ¹⁰ Bq
レントゲン	R	1 R = 2. 58×10 ⁻⁴ C/kg
ラド	rad	1 rad=1cGy=10 ⁻² Gy
レム	rem	1 rem=1 cSv=10 ⁻² Sv
X線単位	1X unit	1X unit=1. 002×10 ⁻⁴ nm
ガッ	γ	1γ=1 nT=10 ⁻⁹ T
ジャンスキー	Jy	1 Jy=10 ⁻²⁶ W・m ⁻² ・Hz ⁻¹
フェルミ	1 fermi	1 fermi=1 fm=10 ⁻¹⁵ m
メートル系カラット	1 metric carat	1 metric carat = 200 mg = 2×10 ⁻⁴ kg
トル	1 Torr	1 Torr = (101 325/760) Pa
標準気圧	1 atm	1 atm = 101 325 Pa
カロリ	1 cal	
マイクロン	1 μ	1 μ =1μm=10 ⁻⁶ m

