

STAX 2 : A Computer Program
for Calculating Neutron Elastic
and Inelastic Scattering Cross
Sections by Means of the Opti-
cal Model and Moldauer's
Theory

July 1970

日本原子力研究所

Japan Atomic Energy Research Institute

日本原子力研究所は、研究成果、調査結果などを JAERI レポートとして、つぎの4種に分けそれぞれの通し番号を付し、不定期に刊行しております。

- | | | |
|---------|--------------------------------|-------------|
| 1. 研究報告 | まとまった研究の成果あるいはその一部における重要な結果の報告 | JAERI 1001- |
| 2. 調査報告 | 総説・展望・調査の結果などをまとめたもの | JAERI 4001- |
| 3. 年報 | 研究・開発その他の活動状況などの報告 | JAERI 5001- |
| 4. 資料 | 施設の概要や手引きなど | JAERI 6001- |

このうち既刊分については「JAERI レポート一覧」にタイトル・要旨をまとめて掲載し、また新刊レポートは「研究成果要旨集」(隔月刊)で逐次紹介しています。

これらのリスト・研究報告書の入手および複写・翻訳などのご要求は日本原子力研究所技術情報部(茨城県那珂郡東海村)に申しこんでください。

Japan Atomic Energy Research Institute publishes the nonperiodical reports with the following classification numbers:

1. JAERI 1001- Research reports
2. JAERI 4001- Survey reports and reviews
3. JAERI 5001- Annual reports
4. JAERI 6001- Manuals etc.

Requests for the above publications, and reproduction and translation should be addressed to Division of Technical Information, Japan Atomic Energy Research Institute, Tokai-mura, Naka-gun, Ibaraki-ken, Japan

STAX 2
A Computer Program for Calculating
Neutron Elastic and Inelastic Scattering Cross Sections
by Means of
the Optical Model and Moldauer's Theory

Summary

This program computes the cross sections for neutron elastic and inelastic scattering using the optical model and Moldauer's theory, and searches for potential parameters which reproduce experimental cross sections.

Improvement is made in the treatment of the effect of resonance interference. The quantity "Q" characterizing this effect has been regarded as an arbitrary parameter in the calculations hitherto published. In the present program this quantity can be calculated considering its dependence on the strength function.

Calculation using the Hauser-Feshbach theory is also possible.

November 1969

Yoshiaki TOMITA
Physics Division Tokai Research Establishment
Japan Atomic Energy Research Institute

STAX 2
光学模型と Moldauer の理論による中性子弾性
非弾性散乱断面積計算コード

要 旨

このコードは、光学模型と Moldauer の理論によって、中性子の弾性非弾性散乱の断面積を計算し、実験値を再現するポテンシャルパラメーターを求める。

これまでの計算では任意なパラメーターとしてあつかわれてきた、共鳴準位の干渉の効果を示す量 "Q" を、強度関数への依存性を考慮することによって計算できるよう改良した。

Hauser-Feshbach の理論による計算も可能になっている。

1969年11月

日本原子力研究所東海研究所
物理部 富 田 芳 明

Contents

1. Introduction	1
2. Theory	1
3. Fitting to experimental cross sections	4
4. Methods of numerical calculations.....	6
5. Subroutines	9
6. Input	9
7. Requirement for machine	12
References.....	13
Appendix A Sample calculation	15
Appendix B Symbolic listing of STAX2	:

目 次

1. 序 論	1
2. 理 論	1
3. 実験値への適合	4
4. 数値計算の方法	6
5. サブルーチン	9
6. インプット	9
7. 計算機に対する要求	12
文 献	13
附録 A 計算例.....	15
附録 B プログラムのリスト.....	:

1. Introduction

STAX2 is a computer program for calculating neutron elastic and inelastic scattering cross sections using the optical model and Moldauer's theory, and can search for potential parameters which can reproduce experimental cross sections. Search is made on all potential parameters with respect to any combination of the following cross sections:

- 1) total cross section,
- 2) elastic scattering cross section (integral or differential),
- 3) inelastic scattering cross section for the first excited level (integral or differential).

Although there exist many optical model programs^{1),2)}, none of them seems to use Moldauer's theory for the calculation of compound nuclear process. The program NEARREX³⁾, often used to compute reaction cross sections using Moldauer's theory, can not compute angular distributions and does not include an optical model routine. These are the reasons why the present program has been written.

In this program improvement is made in the treatment of the resonance interference by considering the dependence of the quantity $Q^{J\pi}$, characterizing this effect, on the strength function.

2. Theory

We consider the scattering of neutrons by a target nucleus of mass M . The cross section σ_{el} for elastic scattering is composed of the shape elastic cross section σ_{se} and the compound elastic cross section σ_{c0} :

$$\sigma_{el} = \sigma_{se} + \alpha\sigma_{c0}. \quad (1)$$

The quantity α is a correction factor equal to or less than unity and is used to approximately correct for the effect of other open channels than neutron channels, because this program can treat only neutrons. This quantity is also used when higher energy levels are not well known. For the inelastic channels we neglect the direct process and the cross sections are given by

$$\sigma_n = \alpha\sigma_{cn}, \quad (2)$$

where n indicates the n -th excited state.

Cross sections computed in the present program are those in the center-of-mass system.

2.1 Shape Elastic Scattering—Optical Model

The shape elastic scattering is treated by the optical model. The optical potential employed in the present program is of the form

$$-Vf(r) - iWg(r) - V_{so}(\mathbf{l} \cdot \boldsymbol{\sigma})h(r), \quad (3)$$

where

$$f(r) = \frac{1}{1 + \exp[(r - r_0 A^{1/3})/a]}, \quad (4)$$

$$g(r) = \frac{4 \exp[(r - r_s A^{1/3})/b]}{\{1 + \exp[(r - r_s A^{1/3})/b]\}^2}, \quad (5)$$

$$h(r) = -\lambda_\pi^2 \frac{1}{r} \frac{d}{dr} g(r), \quad (6)$$

$$\lambda_\pi = 2.04264 \text{ fm}. \quad (7)$$

Each of the seven parameters V , W , V_{so} , r_s , a and b may have an energy dependence

$$V = \sum_{n=0}^5 V^{(n)} E^n \quad \text{etc.}, \quad (8)$$

where E is the kinetic energy in the center-of-mass system.

The radial Schroedinger equation describing the scattering is

$$\left[-\frac{\hbar^2}{2\mu} \left\{ \frac{d^2}{dr^2} - \frac{l(l+1)}{r^2} \right\} - Vf(r) - iWg(r) - C_l^{(\pm)} V_{so} h(r) - E \right] \varphi_l^{(\pm)} = 0, \quad (9)$$

where (\pm) corresponds to the value of coupled angular momentum $j = l \pm 1/2$ of \mathbf{l} and $\boldsymbol{\sigma}$, and

$$C_l^{(+)} = -l, \quad C_l^{(-)} = l+1, \quad (10)$$

$$\mu = \frac{mM}{m+M}. \quad (11)$$

Here m is the neutron mass, which is 1.008665 amu⁻¹²C.

The scattering amplitude $\eta_l^{(\pm)}$ is obtained from the asymptotic behavior of the wave function :

$$\varphi_l^{(\pm)}(r) \underset{r \rightarrow \infty}{\sim} r(h_l^{(2)}(kr) + \eta_l^{(\pm)} h_l^{(1)}(kr)), \quad (12)$$

where $h_l^{(1)}$ and $h_l^{(2)}$ are the spherical Hankel functions and

$$k = \frac{1}{\lambda} = \sqrt{\frac{2\mu E}{\hbar^2}}. \quad (13)$$

The shape elastic cross section and the cross section for the formation of compound nucleus are given by

$$\sigma_{se} = \pi \lambda^2 \sum_l \{ (l+1) |1 - \eta_l^{(+)}|^2 + l |1 - \eta_l^{(-)}|^2 \}, \quad (14)$$

$$\frac{d\sigma_{sc}}{d\Omega} = \frac{\lambda^2}{4} \left[\sum_l \{ (l+1)(1 - \eta_l^{(+)} + l(1 - \eta_l^{(-)}) \} P_l(\cos \theta) \right]^2 + \left[\sum_l (\eta_l^{(+)} - \eta_l^{(-)}) P_l^1(\cos \theta) \right]^2, \quad (15)$$

$$\sigma_c = \pi \lambda^2 \sum_l \{ (l+1) T_l^{(+)} + l T_l^{(-)} \}, \quad (16)$$

where transmission coefficients $T_l^{(\pm)}$ are given by

$$T_l^{(\pm)} = 1 - |\eta_l^{(\pm)}|^2. \quad (17)$$

2.2 Compound Nuclear Process—Moldauer's Theory

By using Moldauer's theory⁽⁴⁾ differential cross section for the n -th excited state can be written in the following form :

$$\frac{d\sigma_{en}}{d\Omega} = \frac{\lambda^2}{2(2I_0+1)} \sum_L B_L P_L(\cos \theta), \quad (18)$$

$$B_L = B_L^{(1)} + B_L^{(2)} + B_L^{(3)}, \quad (19)$$

$$B_L^{(1)} = \frac{(-)^{I_n - I_0}}{4} \sum_{J \Pi I_j I_j'} (2J+1)^2 \delta(\pi_0(-)^I, \Pi) \delta(\pi_n(-)^{I'}, \Pi) \\ \times \bar{Z}(ljlj; \frac{1}{2}L) W(jJjJ; I_0L) \bar{Z}(l'j'l'j'; \frac{1}{2}L) W(j'Jj'J; I_nL) \\ \times \left\langle \frac{\Theta_{\mu, 0l_j} \Theta_{\mu, n l_j'}}{\Theta_{\mu}} \right\rangle_{\mu}^{J \Pi}, \quad (20)$$

$$B_L^{(2)} = \frac{\delta(n, 0)}{4} \sum_{J \Pi (I_j) \neq (I_j')} (2J+1)^2 \left\{ \bar{Z}(lj'l'j'; \frac{1}{2}L) W(jJj'J; I_0L) \right\}^2 \\ \times \delta(\pi_0(-)^I, \Pi) \delta(\pi_0(-)^{I'}, \Pi) R_{I_j, I_j'} B \left\langle \frac{\Theta_{\mu, 0l_j} \Theta_{\mu, 0l_j'}}{\Theta_{\mu}} \right\rangle_{\mu}^{J \Pi}, \quad (21)$$

$$B_{L}^{(3)} = \frac{-\delta(n, 0)}{4} \sum_{JIII l' j'} (2J+1)^2 \left[\bar{Z}(lj l' j'; \frac{1}{2}L) W(jJ j' J; I_0 L) \right]^2 \\ \times \delta(\pi_0(-)^l, II) \delta(\pi_0(-)^{l'}, II) R_{lj, l' j'} Q_{\mu}^{JIII} \langle \Theta_{\mu, 0lj} \rangle_{\mu}^{JIII} \langle \Theta_{\mu, 0l' j'} \rangle_{\mu}^{JIII}, \quad (22)$$

where I_n and π_n are the spin and parity of the n -th excited state (the ground state is denoted by $n=0$), J and II are the total angular momentum and parity, and $\delta(n, m)$ is Kronecker's symbol. The coefficients \bar{Z} and W are the Blatt-Biedenharn and the Racah coefficients, and

$$Q^{JIII} = \frac{2B}{N^2} \left[1 - \Phi_0 \left(\frac{\langle \Theta_{\mu} \rangle_{\mu}^{JIII}}{2N^2} \right) \right], \quad (23)$$

$$B = |b_{nlj}|^2, \quad (24)$$

$$R_{lj, l' j'} = \text{Re}(b_{0lj} b_{0l' j'}^*) / B. \quad (25)$$

Here B is assumed to be independent of (nlj) .

The function Φ_0 is defined by

$$\Phi_0(x) = 1 - \frac{1}{x} \left[1 - \frac{1}{x} e^{-x} \sinh x \right] - \frac{1}{x} \text{Ei}(-x) \left[\cosh x - \frac{1}{x} \sinh x \right], \quad (26)$$

and

$$\Theta_{\mu} = \sum_{nlj} \Theta_{\mu, nlj}. \quad (27)$$

Other quantities $\Theta_{\mu, nlj}$, b_{nlj} and N are defined in ref.⁴⁾ and the average $\langle \rangle_{\mu}^{JIII}$ is taken over the compound levels μ which have a spin and parity of $JIII$. The quantity $\langle \Theta_{\mu, nlj} \rangle_{\mu}^{JIII}$ is proportional to the strength function $\langle \Gamma_{\mu, nlj} / D \rangle_{\mu}^{JIII}$, and is related to the transmission coefficient T_{nlj} by

$$T_{nlj} = \langle \Theta_{\mu, nlj} \rangle_{\mu}^{JIII} - \frac{Q^{JIII}}{4} \{ \langle \Theta_{\mu, nlj} \rangle_{\mu}^{JIII} \}^2. \quad (28)$$

Assuming a χ^2 -distribution with ν degrees of freedom for $\Theta_{\mu, nlj}$ we obtain⁵⁾

$$\left\langle \frac{\Theta_{\mu, nlj} \Theta_{\mu, n' l' j'}}{\Theta_{\mu}} \right\rangle_{\mu}^{JIII} = \frac{\langle \Theta_{\mu, nlj} \rangle_{\mu}^{JIII} \langle \Theta_{\mu, n' l' j'} \rangle_{\mu}^{JIII}}{\langle \Theta_{\mu} \rangle_{\mu}^{JIII}} S_{nlj, n' l' j'}^{JIII}, \quad (29)$$

$$S_{nlj, n' l' j'}^{JIII} = \int_0^{\infty} dt \frac{1 + 2\delta_{nn'} \delta_{ll'} \delta_{jj'}}{f_{nlj} f_{n' l' j'} \Pi_{n' l' j'}^{JIII} f_{n' l' j'}^{\nu/2}}, \quad (30)$$

where the product in the denominator is taken over all channels which have a total angular momentum and parity of $JIII$, and

$$f_{nlj} = 1 + \frac{\langle \Theta_{\mu, nlj} \rangle_{\mu}^{JIII} 2t}{\langle \Theta_{\mu} \rangle_{\mu}^{JIII} \nu}. \quad (31)$$

Various quantities given above have limited ranges of magnitude:

$$1 \geq B \geq 0, \quad (32)$$

$$N \geq 1, \quad (33)$$

$$1 \geq R_{lj, l' j'} \geq -1, \quad (34)$$

$$R_{lj, lj} = 1, \quad (35)$$

$$2 \geq Q^{JIII} \geq 0, \quad (36)$$

and in the limit of $\Gamma/D \rightarrow 0$ they become

$$B = 1, \quad (37)$$

$$N = 1, \quad (38)$$

$$R_{lj, l' j'} = 1, \quad (39)$$

$$Q^{JIII} = 2. \quad (40)$$

The quantity $R_{lj, l' j'}$ for $(lj) \neq (l' j')$ is assumed to be independent of $(lj; l' j')$ in the present program. This quantity does not affect the cross sections when the spin of the target nucleus is zero.

In the calculations hitherto published, Q^{JIII} was treated as a constant neglecting its dependence on $\langle \Theta_{\mu, nlj} \rangle_{\mu}^{JIII}$. However, this treatment sometimes results in negative values⁶⁾ for the compound

elastic cross sections for some partial waves, when many channels are open and the constant is not taken sufficiently small. This difficulty is avoided if Eq. (28) is solved for $\langle \Theta_{\mu, nlj} \rangle_{\mu}^{JII}$ by considering the functional dependence of Q^{JII} on $\langle \Theta_{\mu, nlj} \rangle_{\mu}^{JII}$ given by Eq. (23). Both treatment of Q^{JII} are possible in the present program.

The largest possible Q^{JII} is given for $B=1$ and $N=1$.

By putting b_{nlj} to zero and $S_{nlj; n'l'j}^{JII}$ to unity, B and Q^{JII} become zero and the Hauser-Feshbach formula is obtained in Eqs. (18)-(22).

3. Fitting to Experimental Cross Sections

Fitting the calculation to experimental cross sections is made by minimizing the quantity χ^2 defined by

$$\chi^2 = \left\{ W_t \left(\frac{\sigma_t^{\text{cal}} - \sigma_t^{\text{exp}}}{\Delta \sigma_t} \right)^2 + W_{\text{el}} \left(\frac{\sigma_{\text{el}}^{\text{cal}} - \sigma_{\text{el}}^{\text{exp}}}{\Delta \sigma_{\text{el}}} \right)^2 + W_1 \left(\frac{\sigma_1^{\text{cal}} - \sigma_1^{\text{exp}}}{\Delta \sigma_1} \right)^2 \right. \\ \left. + W_{\text{ela}} \sum_{\theta} \left(\frac{\frac{d\sigma_{\text{el}}^{\text{cal}}(\theta)}{d\Omega} - \frac{d\sigma_{\text{el}}^{\text{exp}}(\theta)}{d\Omega}}{\Delta \frac{d\sigma_{\text{el}}(\theta)}{d\Omega}} \right)^2 + W_{1a} \sum_{\theta} \left(\frac{\frac{d\sigma_1^{\text{cal}}(\theta)}{d\Omega} - \frac{d\sigma_1^{\text{exp}}(\theta)}{d\Omega}}{\Delta \frac{d\sigma_1(\theta)}{d\Omega}} \right)^2 \right\} \\ / (W_t + W_{\text{el}} + W_1 + \sum_{\theta} W_{\text{ela}} + \sum_{\theta} W_{1a}), \quad (41)$$

where σ_t is the total cross section, "cal" and "exp" indicate the calculated and the experimental cross sections, Δ indicates the experimental uncertainty, and W_t etc. are weighting factors which may be arbitrarily chosen. Search for the "best" parameters is made on any number of parameters $V^{(0)}$, $W^{(0)}$, $V_{\text{so}}^{(0)}$, $r_0^{(0)}$, $r_s^{(0)}$, $a^{(0)}$ and $b^{(0)}$ by means of Gauss-Newton's method.

For brevity we denote the set of parameters to be searched for by $\{x_i\}$ and the cross sections to be fitted by $\{\sigma_{\alpha}\}$. Around $\{x_i^{(n)}\}$ the cross sections are approximated by

$$\sigma_{\alpha}^{\text{cal}}(\{x_i^{(n)} + \delta x_i\}) = \sigma_{\alpha}^{\text{cal}}(\{x_i^{(n)}\}) + \sum_i \frac{\partial \sigma_{\alpha}^{\text{cal}}(\{x_i^{(n)}\})}{\partial x_i} \delta x_i, \quad (42)$$

where $\{x_i^{(n)}\}$ indicates the parameters obtained after the n-th iteration. The parameters which minimize χ^2 can be obtained by

$$\sum_j \frac{W_{\alpha}}{(\Delta \sigma_{\alpha})^2} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_i} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_j} \delta x_j = \sum_{\alpha} \frac{W_{\alpha} (\sigma_{\alpha}^{\text{exp}} - \sigma_{\alpha}^{\text{cal}})}{(\Delta \sigma_{\alpha})^2} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_i}. \quad (43)$$

This equation can be solved easily. However, since we neglected higher order terms of $\{\delta x_i\}$ in Eq. (42), this equation is not very accurate when some of the eigenvalues of the matrix

$$\left(\sum_{\alpha} \frac{W_{\alpha}}{(\Delta \sigma_{\alpha})^2} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_i} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_j} \right)$$

are small, and such circumstances may often happen because of the well known optical potential ambiguities. For this reason Eq. (43) is not solved as it is. We define a symmetric matrix and column vectors in the parameter space:

$$(A_{ij}) = \left(\sum_{\alpha} \frac{W_{\alpha}}{(\Delta \sigma_{\alpha})^2} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_i} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_j} \delta x_{i, \text{max}} \delta x_{j, \text{max}} \right), \quad (44)$$

$$(b_i) = \left(\sum_{\alpha} \frac{W_{\alpha} (\sigma_{\alpha}^{\text{cal}} - \sigma_{\alpha}^{\text{exp}})}{(\Delta \sigma_{\alpha})^2} \frac{\partial \sigma_{\alpha}^{\text{cal}}}{\partial x_i} \delta x_{i, \text{max}} \right), \quad (45)$$

$$(\delta \xi_i) = \left(\frac{\delta x_i}{\delta x_{i, \text{max}}} \right), \quad (46)$$

where $\{\delta x_{i, \text{max}}\}$ is introduced in order to make the magnitude of A_{ij} uniform and also to set

a limit for $\{\delta x_i\}$, and is given as follows:

$$\left. \begin{aligned} \delta V^{\max} = \delta W^{\max} = \delta V_{s_0}^{\max} &= 3.0 \text{ MeV}, \\ \delta r_0^{\max} &= 0.05 \text{ fm}, \\ \delta r_s^{\max} &= 0.1 \text{ fm}, \\ \delta a^{\max} &= 0.07 \text{ fm}, \\ \delta b^{\max} &= 0.1 \text{ fm}. \end{aligned} \right\} \quad (47)$$

Then Eq. (43) can be written as

$$\mathbf{A}\delta\boldsymbol{\xi} = \mathbf{b}. \quad (48)$$

The matrix \mathbf{A} is diagonalized by a orthogonal matrix \mathbf{U} :

$$\mathbf{U}\mathbf{A}\mathbf{U}' = \mathbf{D} = \begin{pmatrix} \lambda_1 & 0 \\ & \lambda_2 \\ 0 & \ddots \end{pmatrix}. \quad (49)$$

Solution of Eq. (48) is given by

$$\delta\boldsymbol{\xi} = \mathbf{U}'\delta\boldsymbol{\eta}, \quad (50)$$

$$\delta\boldsymbol{\eta} = \mathbf{D}^{-1}\mathbf{U}\mathbf{b}. \quad (51)$$

Because of the reason mentioned above, we neglect small eigenvalues and replace \mathbf{D}^{-1} by $\tilde{\mathbf{D}}^{-1}$, the elements of which are given by

$$\tilde{D}_{ij}^{-1} = \begin{cases} \frac{\delta_{ij}}{\lambda_i}, & \text{if } |\lambda_i| \geq \varepsilon \\ 0, & \text{if } |\lambda_i| < \varepsilon, \end{cases} \quad (52)$$

where ε is a small number defined below in Eq. (57). Furthermore, since the approximation (42) is valid only in the neighbourhood of $\{x_i^{(n)}\}$, $\delta\boldsymbol{\eta}$ is truncated as follows:

$$\delta\eta_i = \begin{cases} \tilde{\delta\eta}_i, & \text{if } |\tilde{\delta\eta}_i| \leq 1 \\ \frac{\tilde{\delta\eta}_i}{|\tilde{\delta\eta}_i|}, & \text{if } |\tilde{\delta\eta}_i| > 1, \end{cases} \quad (53)$$

where

$$\tilde{\delta\boldsymbol{\eta}} = \tilde{\mathbf{D}}^{-1}\mathbf{U}\mathbf{b}. \quad (54)$$

At $\{x_i^{(n)} + \delta x_i\}$, χ^2 is expected to decrease by

$$\Delta\chi^2 = \sum_{i, |\lambda_i| \geq |\varepsilon|} \lambda_i (\delta\eta_i)^2 \quad (55)$$

in the same approximation as above. When $\chi^2(\{x_i^{(n)}\}) - \chi^2(\{x_i^{(n)} + \delta x_i\}) \geq \frac{\Delta\chi^2}{4}$, $\{x_i^{(n)} + \delta x_i\}$ is taken as $\{x_i^{(n+1)}\}$. When $\chi^2(\{x_i^{(n)}\}) - \chi^2(\{x_i^{(n)} + \delta x_i\}) \leq \frac{\Delta\chi^2}{4}$, we consider the approximation (42) to be not good and compute χ^2 at $\{x_i^{(n)} + \frac{1}{2}\delta x_i\}$. Then $\{x_i^{(n+1)}\}$ is obtained by approximating $\chi^2(x_i^{(n)} + \alpha\delta x_i)$ by a quadratic polynomial of α .

This parameter search is stopped when one of the following conditions is encountered:

- 1) a specified number of iteration has been made,
- 2) specified calculation time has elapsed,
- 3) χ^2 has become smaller than a specified value χ_{\min}^2 ,
- 4) $\chi^2(\{x_i^{(n)}\}) - \chi^2(\{x_i^{(n+1)}\}) \leq 0.01 \chi_{\min}^2$.

The matrix \mathbf{A} is diagonalized by means of Jacobi's method and the accuracy of diagonalization is taken as

$$\varepsilon' = \frac{10^{-K}}{N_p} \sum A_{ii}, \quad (56)$$

where N_p is the number of the parameter to be searched for and K is usually set equal to five unless otherwise specified. The ε in Eq. (52) is given by

$$\varepsilon = 10 N_p \varepsilon'. \quad (57)$$

Derivative $\frac{\partial \sigma_\alpha}{\partial x_i}$ is approximated by

$$\frac{\sigma_\alpha(x_i^{(n)} + \Delta x_i) - \sigma_\alpha(x_i^{(n)})}{\Delta x_i}, \quad (58)$$

where $\{\Delta x_i\}$ is given as follows:

$$\left. \begin{aligned} \Delta V = \Delta W = \Delta V_{s_0} = 0.1 \text{ MeV}, \\ \Delta r_0 = r_s = \Delta a = \Delta b = 0.01 \text{ fm}. \end{aligned} \right\} \quad (59)$$

In computing $\sigma_\alpha(\{x_i^{(n)} + \Delta x_i\})$ two methods are employed:

- (A) $\sigma_\alpha(\{x_i^{(n)} + \Delta x_i\})$ is calculated rigorously,
- (B) only the shape elastic cross section and the cross section for the formation of compound nucleus are calculated, and the cross sections through compound process are approximated as follows:

$$\sigma_{cn}(x_i^{(n)} + \Delta x_i) = \frac{\sigma_c(x_i^{(n)} + \Delta x_i)}{\sigma_c(x_i^{(n)})} \sigma_{cn}(x_i^{(n)}), \quad (60)$$

$$\frac{d\sigma_{cn}(x_i^{(n)} + \Delta x_i)}{d\Omega} = \frac{\sigma_c(x_i^{(n)} + \Delta x_i)}{\sigma_c(x_i^{(n)})} \frac{d\sigma_{cn}(x_i^{(n)})}{d\Omega}. \quad (61)$$

Method (B) is inferior in accuracy to method (A), but requires much shorter calculation time especially when many channels are open. In most cases method (B) is preferred.

4. Methods of Numerical Calculation

4.1 Solution of the Schroedinger Equation

Equation (9) is solved by means of Noumerov's method⁷⁾ in the internal region. Rewriting Eq. (9) as

$$\frac{d^2}{dr^2} \varphi_l^{(\pm)} = D_l^{(\pm)} \varphi_l^{(\pm)}, \quad (62)$$

and defining the quantity

$$\Psi_l^{(\pm)} = \left(1 - \frac{\hbar^2}{12} D_l^{(\pm)} \right) \varphi_l^{(\pm)}, \quad (63)$$

we can replace Eq. (9) by the recurrence relation

$$\Psi_l^{(\pm)}(r_n) = \left[2 + \frac{\hbar^2 D_l^{(\pm)}(r_{n-1})}{1 - \frac{\hbar^2}{12} D_l^{(\pm)}(r_{n-1})} \right] \Psi_l^{(\pm)}(r_{n-1}) - \Psi_l^{(\pm)}(r_{n-2}), \quad (64)$$

where

$$r_n = r_1 + (n-1)h \quad (65)$$

and $h=0.25$ fm unless otherwise specified. Two values $\Psi_l^{(\pm)}(r_1)$ and $\Psi_l^{(\pm)}(r_2)$ are necessary for solving Eq. (65). These values are obtained by expanding the functions in power series of r :

$$\varphi_l^{(\pm)}(r) = \sum_{n=0}^{\infty} b_{ln}^{(\pm)} r^{n+l+1}, \quad (66)$$

$$D_l^{(\pm)}(r) = \frac{l(l+1)}{r^2} - \frac{K_l^{(\pm)}}{r} - \sum_{n=0}^{\infty} a_{ln}^{(\pm)} r^n, \quad (67)$$

where

$$K_l^{(\pm)} = \frac{2\mu \lambda_\pi^2 V_{s_0} C_l^{(\pm)}}{\hbar^2} \frac{\exp(-r_0 A^{1/3}/a)}{\{1 + \exp(-r_0 A^{1/3}/a)\}^2}, \quad (68)$$

$$a_{ln}^{(\pm)} = \frac{2\mu}{\hbar^2} [E\delta_{n0} + f_n V + 4ig_n W + (n+2)\lambda_\pi^2 V_{so} C_l^{(\pm)} f_{n+2}]. \tag{69}$$

The coefficients f_n and g_n are the derivatives of potentials and are given by the following equations:

$$f_n = F_n(r_0, a) \tag{70}$$

$$g_n = -b(n+1)F_n(r_s, b), \tag{71}$$

$$F_0(r, c) = [1 + \exp(rA^{1/3}/c)]^{-1}, \tag{72}$$

$$F_{n+1}(r, c) = \frac{1}{c(n+1)} \left[-F_n(r, c) + \sum_{k=0}^n F_k(r, c) F_{n-k}(r, c) \right] \tag{73}$$

The expansion coefficients of $\varphi_l^{(\pm)}$ are given by

$$b_{l0}^{(\pm)} = 1, \tag{74}$$

$$b_{l1}^{(\pm)} = -\frac{K_l^{(\pm)}}{2(l+1)}, \tag{75}$$

$$b_{ln}^{(\pm)} = \frac{-1}{n(2l+n+1)} \left[K_l^{(\pm)} b_{l, n-1}^{(\pm)} + \sum_{k=0}^{n-2} b_{lk}^{(\pm)} a_{l, n-k-2}^{(\pm)} \right]. \tag{76}$$

Only the first eight terms are taken in the expansion of $\varphi_l^{(\pm)}$.

Because of the singularity of $l(l+1)/r^2$, Eq. (64) is not accurate in the neighbourhood of $r=0$. Therefore r_1 and r_2 should be taken as large as possible. Error of $\varphi_l^{(\pm)}$ due to the truncation of the power series (66) is of the order of

$$\frac{(Kr)^8}{9!} \varphi_l^{(\pm)}(r), \tag{77}$$

where

$$K = \sqrt{\frac{2\mu}{\hbar^2}(E+V)}. \tag{78}$$

We require $\varphi_l^{(\pm)}(r_2)$ to be accurate to 10^{-5} and take r_2 as

$$r_2 = \frac{1.2}{K}. \tag{79}$$

4.2 Determination of Scattering Amplitude

Scattering amplitude $\eta_l^{(\pm)}$ is obtained by matching the internal wave function with the external wave function (12):

$$\frac{\varphi_l^{(\pm)}(r_M)}{\varphi_l^{(\pm)}(r_{M-1})} = \frac{r_M}{r_{M-1}} \frac{h_l^{(2)}(kr_M) + \eta_l^{(\pm)} h_l^{(1)}(kr_M)}{h_l^{(2)}(kr_{M-1}) + \eta_l^{(\pm)} h_l^{(1)}(kr_{M-1})}. \tag{80}$$

Matching radius r_M is determined by the requirement that at $r > r_M$ influence of the optical potential is negligible. By rewriting Eq. (9) as

$$\frac{d^2}{dr^2} \varphi_l^{(\pm)} + \left\{ k^2 - \frac{l(l+1)}{r^2} - v(r) \right\} \varphi_l^{(\pm)} = 0, \tag{81}$$

relative error of $\varphi_l^{(\pm)}(r_M)$ introduced by neglecting $v(r)$ at $r > r_M$ can be shown to be of the order of

$$\frac{1}{k} \int_{r_M}^{\infty} dr r^2 v(r) h_l^{(1)}(kr) h_l^{(2)}(kr) \lesssim \frac{1}{k} \int_{r_M}^{\infty} dr |v(r)|. \tag{82}$$

The right hand side of the inequality becomes

$$\frac{|C|}{\alpha k} e^{-\alpha r_M} \tag{83}$$

for $v(r) = Ce^{-\alpha r}$. Therefore we require r_M to satisfy the following conditions

$$\left. \begin{aligned} \frac{akV}{E} e^{\frac{r_0 A^{1/3} - r_M}{a}} < 10^{-6}, \\ \frac{4bkW}{E} e^{\frac{r_2 A^{1/3} - r_M}{a}} < 10^{-6}, \\ \frac{10k\lambda\pi^2 V_{s0}}{a A^{1/3} E} e^{\frac{r_0 A^{1/3} - r_M}{a}} < 10^{-6}. \end{aligned} \right\} \quad (84)$$

4.3 Calculation of Q^{JII} and $\langle \theta_{\mu, nlj} \rangle_{\mu}^{JII}$

Equation (28) can be rewritten as

$$\theta_c = \frac{2}{Q} \{1 - \sqrt{Q T_c}\} \quad (85)$$

and θ_c may be regarded as a function of Q , where we denote $\langle \theta_{\mu, nlj} \rangle_{\mu}^{JII}$, T_{nlj} and Q^{JII} by θ_c , T_c and Q for brevity. Then Eq. (23) becomes an equation with only one variable Q . This equation is solved by Newton's iteration method. Starting values are taken as

$$\theta_c^{(0)} = T_c, \quad (86)$$

$$Q^{(0)} = \frac{2B}{N^2} \left(1 - \Phi_0 \left(\frac{\sum \theta_c^{(0)}}{2N^2} \right) \right), \quad (87)$$

and the next value is given by

$$Q^{(n+1)} = Q^{(n)} + \frac{\frac{2B}{N^2} \left\{ 1 - \Phi_0 \left(\frac{\sum \theta_c^{(n)}}{2N^2} \right) \right\} - Q^{(n)}}{1 + \frac{B}{4N^2} \sum_c \frac{\{\theta_c^{(n)}\}^2}{\sqrt{1 - Q_c^{(n)} T_c}} \Phi_0' \left(\frac{\sum \theta_c^{(n)}}{2N^2} \right)}, \quad (88)$$

where $\Phi_0'(x)$ is replaced by

$$100 \{ \Phi_0(x + 0.01) - \Phi_0(x) \}. \quad (89)$$

Iteration is continued until

$$|\theta_c^{(n+1)} - \theta_c^{(n)}| < 10^{-6} \quad (90)$$

is reached for all c .

4.4 Calculation of Width Fluctuation Factor

For brevity we denote the integral in Eq. (30) by

$$I = \int_0^{\infty} \frac{dt}{\prod_k (1 + \alpha_k t)^{\nu/2}}. \quad (91)$$

By the transformation

$$y = (1 + \bar{\alpha}t)^{1-n} \quad (92)$$

the integral becomes

$$I = \frac{1}{(n-1)\bar{\alpha}} \int_0^1 \frac{dy}{\prod_k Y_k}, \quad (93)$$

where

$$n = \frac{1}{2} \sum_k \nu, \quad (94)$$

$$\bar{\alpha} = \frac{1}{n} \sum_k \alpha_k, \quad (95)$$

$$Y_k = \left(1 - \frac{\alpha_k}{\bar{\alpha}} \right) y^{\bar{n}-1} + \frac{\alpha_k}{\bar{\alpha}}. \quad (96)$$

The integrand is a monotonously decreasing function of y , and when values of α_k are scattered widely around their average $\bar{\alpha}$ it has a large derivative in the neighbourhood of $y=0$. In order to avoid this rapid variation further transformation

$$y = z^5 \quad (97)$$

is made and the integral becomes

$$I = \frac{5}{(n-1)\bar{\alpha}} \int_0^1 \frac{z^4}{\prod_k Y_k} dz, \quad (98)$$

which is evaluated by means of the 16-point Legendre-Gauss formula.

5. Subroutines

The following subroutines are used in the present program.

1. SUBROUTINE SEARCH searches for the "best" potential parameters.
2. SUBROUTINE EIGEN diagonalizes the matrix A in Eq. (49).
3. SUBROUTINE XSECT computes various cross sections.
4. SUBROUTINE DIFEQ solves the Schroedinger equation.
5. SUBROUTINE BESSEL AND BESSE2 compute the spherical Bessel and Neumann functions.
6. SUBROUTINE SHAPE computes the shape elastic cross section.
7. SUBROUTINE LEGND1 and LEGND2 compute the Legendre functions.
8. SUBROUTINE TRANS computes Q^{JM} and $\langle \Theta_{\mu, nl, j} \rangle_{\mu}^{JM}$.
9. SUBROUTINE PHAI computes the function Φ_0 .
10. SUBROUTINE EI computes the exponential integral.
11. SUBROUTINE BLL computes B_L in Eq. (18).
12. FUNCTION S computes the width fluctuation factor.
13. SUBROUTINE FACTRL, FUNCTION DELT and EXDELT are used for computing the Clebsch-Gordan coefficient and the Racah coefficient.
14. FUNCTION CG0 computes the Clebsch-Gordan coefficient.
15. FUNCTION WRAC computes the Racah coefficient.

6. Input

The meaning and FORMAT of input cards are described below. Units of energy and length are MeV and fm, respectively.

CARD 1

(DATE(N), N=1, 3), (NAME(N), N=1, 3)

FORMAT (6 A 4)

Date and name.

CARD 2

(TITLE(N), N=1, 18)

FORMAT (18 A 4)

Title of the calculation.

CARD 3

FMASS, NOMASS, NLEV, NLOUT, ELAB, PHIPRE, DR, ITEST, IFDPR, NPRES.

FORMAT (F 12. 0, 3 I 6, F 6. 0, E 12. 3, F 6. 0, 3 I 6)

FMASST : Mass of the target nucleus in amu-C¹².

NOMASS : Mass number of the target nucleus.

NLEV : Number of levels. (≤ 25)

NLOUT : Cross sections are calculated only for the lowest NLOUT levels.

ELAB : Energy of the incident neutrons in the laboratory system.

PHIPRE : The factor 10^{-6} in (84) is replaced by PHIPRE, if PHIPRE $\neq 0$.

DR : If ITEST $\neq 0$, matrices **A**, **b**, **D** and **U** in Eqs. (44)-(49) are printed as output.

IFDPR : If IFDPR=1, $\{\Delta x_i\}$ in Eqs. (59) are replaced by the values on CARD 4.
If IFDPR=2, $\{\delta x_i^{\max}\}$ in Eqs. (47) are replaced by the values on CARD 5.
If IFDPR=3, both quantities are replaced.

NPRE : If NPRE $\neq 0$, the factor K in Eq. (56). is replaced by NPRE.

CARD 4 (Necessary only when IFDPR=1 or 3.)

(DPARA(N), N=1, 7)

FORMAT (7 F 6. 0)

: ΔV , ΔW , ΔV_{s_0} , Δr_0 , Δr_s , Δa and Δb in Eqs. (59).

CARD 5 (Necessary only when IFDPR=2 or 3.)

(DRPMAX(N), N=1, 7)

FORMAT (7 F 6. 0)

: δV^{\max} , δW^{\max} , $\delta V_{s_0}^{\max}$, δr_0^{\max} , δr_s^{\max} , δa^{\max} and δb^{\max} in Eqs. (47).

CARD 6

(EX(N), N=1, NLEV)

FORMAT (12 F 6. 0)

: Excitation energy of each levels in ascending order.

CARD 7

(SPIN(N), PARITY(N), N=1, NLEV)

FORMAT (12(F 5. 0, A 1))

: Spin and parity of each levels. Parities are punched as + or -.

CARD 8

NANGL, IFCM, IFCOS

FORMAT (3 I 6)

NANGL : Number of angles for which differential cross sections are calculated. (≤ 20)

IFCM : If IFCM=1, the values of angles are those of in the center-of-mass system,
and if IFCM=0, they are in the laboratory system.

IFCOS : If IFCOS=1, angles are given by cosine, and if IFCOS=0, they are given
in degrees.

CARD 9

(ANG(N) or COSIN(N), N=1, NANGL)

FORMAT (12 F 6. 0)

: Values of angles.

CARD 10

NSRCH0, NSRCH, CHIMIN, TIME, WGTOT, WGEL, WG 1, WGELA, WG 1 A.

FORMAT (2 I 6, 7 F 6. 0)

NSRCH0 and NSRCH: Number of iteration of the parameter search by method (B)
and by method (A). Usually method (B) is tried first and then method (A),
but if NSRCH < 0, method (A) is tried first. Either may be zero.

CHIMIN : Search ends when χ^2 has become smaller than CHIMIN.

TIME : Search is interrupted when calculation time has exceeded TIME.
TIME is given in minutes.

WGTOT, WGEL, WG1, WGELA and WG1A:

Weighting factors W_t , W_{el} , W_1 , W_{ela} and W_{1a} in Eq. (41).

CARD 11

IFV, VO, (VE(N), N=1, 5)

FORMAT (A 1, F 11.0, 5 F 12.0)

IFV : When search is made for the depth V of real potential, * is punched for IFV.

VO : $V^{(0)}$ in Eq. (8).

VE(N) : $V^{(N)}$ in Eq. (8).

CARD 12

IFW, WO (WE(N) N=1, 5)

FORMAT (A 1, F 11.0, 5 F 12.0)

: Same as CARD 11 for the depth W of imaginary potential.

CARD 13

IFVSO, VSOO, (VSOE(N) N=1, 5)

FORMAT (A 1, F 11.0, 5 F 12.0)

: Same as CARD 11 for the depth V_{so} of spin-orbit potential.

CARD 14

IFRR, RRO, (RRE(N), N=1, 5)

FORMAT (A 1, F 11.0, 5 F 12.0)

: Same as CARD 11 for the radius parameter r_0 of real potential

CARD 15

IFRS, RSO, (RSE(N), N=1, 5)

FORMAT (A 1, F 11.0, 5 F 12.0)

: Same as CARD 11 for the radius parameter r_s of imaginary potential.

CARD 16

IFA, AO, (AE(N), N=1, 5)

FORMAT (A 1, F 11.0, 5 F 12.0)

: Same as CARD 11 for the diffuseness a of real potential.

CARD 17

IFB, BO, (BE(N), N=1, 5)

FORMAT (A 1, F 11.0, 5 F 12.0)

: Same as Card 11 for the diffuseness b of imaginary potential.

CARD 18 (Necessary only when at least one of WGTOT, WGEL and WG1 is not zero.)

STOTEX, ERTOT, SELEX, EREL, S1STEX, ERIST

FORMAT (6 F 6.0)

STOTEX and ERTOT: Experimental total cross section and its uncertainty.

SELEX and EREL : Same as above for the elastic scattering cross section.

S1STEX and ER1ST : Same as above for the first level cross section.

(CARD 19-21 are for the experimental values of differential elastic scattering cross sections, and are necessary only when WGELA \neq 0.)

CARD 19

ERRAT

FORMAT (F 6.0)

: Uncertainly in cross section in percent. When ratios of uncertainty to cross section

are different for various angles, ERRAT is set equal to zero and uncertainties are given on CARD 21.

CARD 20

(SELAEX(N), N=1, NANGL)

FORMAT (12 F 6. 0)

: Values of the differential cross section in mb/sr. For angles where experimental values are lacking, SELAEX (N) are set equal to zero. They are not included in the χ^2 calculation.

CARD 21 (Necessary only when ERRAT=0.)

(ERELA(N), N=1, NANGL)

FORMAT (12 F 6. 0)

: Uncertainty in the above values in mb/sr.

(CARD 22-24 are for the experimental values of differential inelastic scattering cross sections leading to the first excited state, and are necessary only when $WG1A \neq 0$. The meaning of these cards are same as CARD 19-21.)

CARD 22

ERRAT

FORMAT (F 6. 0)

CARD 23

(S1AEX(N), N=1, NANGL)

FORMAT (12 F 6. 0)

CARD 24 (Necessary only when ERRAT=0.)

(ER1A(N), N=1, NANGL)

FORMAT (12 F 6. 0)

CARD 25

FCHI, IQCONS, QCONST, BC, REBB, FNR, COMP

FORMAT (F 6. 0, I 6, 5 F 6. 0)

FCHI : Degrees of freedom of the χ^2 -distribution for $\theta_{\mu, n l j}$. IF FCHI=0, cross sections are calculated by means of the Hauser-Feshbach theory.

IQCONS and QCONST: If IQCONS=0, Q^{Jn} is calculated as described in Sec. 2. 2. If IQCONS=1, Q^{Jn} is set equal to QCONST.

BC : B in Eq. (24).

REBB : $R_{l j, l' j'}$ for $(l' j) \neq (l' j')$ in Eq. (25).

FNR : N in Eq. (23).

COMP : α in Eqs. (1) and (2).

7. Requirement for Machine

This program is coded in FORTAN for a FACOM-230/60 computer which has a library subroutine CLOCK which gives the elapsed time in units of Second. Memory size required for this program is about 40K words.

The author is indebted to S. IGARASHI for making available the code ELIESE-2 and for advices. Comparison with the calculation by ELIESE-2 greatly facilitated the error checking of STAX2. He also wishes to express his thanks to K. TSUKADA, K. OKAMOTO and M. MARUYAMA for careful reading of the manuscript.

References

- 1) AUERBACH E. H.: BNL 6562, (1962) (*unpublished*).
- 2) IGARASI S.: Program ELIESE-2, A FORTRAN-IV Program for Calculation of the Nuclear Cross Sections by Use of the Optical Model and Hauser-Feshbach's Method JAERI 1169 (1968).
- 3) MOLDAUER P. A., ENGELBRECHT C. A and DUFFY G. J.: ANL-6978, (1964) (*unpublished*).
- 4) MOLDAUER P. A.: *Phys. Rev.* **135**, B 642 (1964)
- 5) MOLDAUER P. A.: *Rev. Mod. Phys.* **36**, 1079 (1964)
- 6) TSUKADA K., TANAKA Y., TOMITA Y. and MARUYAMA M.: *Nucl. Phys.* **A 125**, 641 (1969).
- 7) MELKANOFF M. A., SAWADA T. and RAYNAL J.: "Method in Computational Physics", Vol. 6, Academic Press, New York.

This is a blank page.

Appendix A Sample Calculation

For illustration, the input and output of a calculation is shown below. In this calculation fitting is tried to the data of Zn at $E_n=2.24$ MeV by using the largest possible values of Q .

氏名		日付		プログラム名		JOB NO.		カード色指定		PUNCH		
所属		電話		研究テーマ番号		IBJOB DECKNAME		備考		YES <input type="checkbox"/> NO <input type="checkbox"/>		
1	2	3	4	5	6	7	8	9	0	1	2	
1969	10	20	Y. TOMITA		ZN-2.24MEV		G-MAX					
		65.289	65	4	4	2.24						
0.0	1.02	1.835	1.90									
0+	2+	2+	0+									
19												
0.	10.	20.	30.	40.	50.	60.	70.	80.	90.	100.	110.	
120.	130.	140.	150.	160.	170.	180.						
20		25	5.									
*	50.											
*	10.											
	10.											
	1.2											
	1.2											
*	.65											
*	.50											
			.690	.013								
0.												
		550.0	368.1	230.1	137.2	88.0	82.0	90.5	104.8	110.6		
112.0	95.3	92.0	107.0									
		116.0	13.0	10.0	7.8	5.4	5.1	5.4	6.0	6.0		
6.8	5.7	5.4	6.0									
1.0	0											

ZN=2.24MEV G-MAX 1969-10-20 Y.TOMITA

MASS OF TARGET NUCLEUS..... 65.2890
 MASS NUMBER..... 65
 ENERGY IN LAB. SYSTEM..... 2.2400 MEV
 ENERGY IN C.M. SYSTEM..... 2.2059 MEV

	GROUND	1ST	2ND	3RD
E	0.000	1.020	1.835	1.900
SPIN	0 +	2 +	2 +	0 +

INITIAL POTENTIAL PARAMETERS

POWERS OF E	0	1	2	3	4	5
* V	50.000	-0.000	-0.000	-0.000	-0.000	-0.000
* W	10.000	-0.000	-0.000	-0.000	-0.000	-0.000
VSD	10.000	-0.000	-0.000	-0.000	-0.000	-0.000
RR	1.200	-0.000	-0.000	-0.000	-0.000	-0.000
RS	1.200	-0.000	-0.000	-0.000	-0.000	-0.000
* A	0.650	-0.000	-0.000	-0.000	-0.000	-0.000
* B	0.500	-0.000	-0.000	-0.000	-0.000	-0.000

*'S INDICATE THE PARAMETERS TO BE SEARCHED FOR

EXPERIMENTAL CROSS SECTIONS (AND ERRORS)

1ST LEVEL CROSS SECTION... 0.690(0.013) BARN WEIGHT..... 1.000

ELASTIC ANGULAR DISTRIBUTION (MILLIBARN/STERAD.) WEIGHT..... 1.000

ANGLE	CROSS SECTION	ANGLE	CROSS SECTION	ANGLE	CROSS SECTION	ANGLE	CROSS SECTION
0.0	-0.0(-0.0)	10.0	-0.0(-0.0)	20.0	-0.0(-0.0)	30.0	550.0(16.0)
40.0	368.1(13.0)	50.0	230.1(10.0)	60.0	137.2(7.8)	70.0	88.0(5.4)
80.0	82.0(5.1)	90.0	90.5(5.4)	100.0	104.8(6.0)	110.0	110.6(6.0)
120.0	112.0(6.3)	130.0	95.3(5.7)	140.0	92.0(5.4)	150.0	102.2(6.0)
160.0	-0.0(-0.0)	170.0	-0.0(-0.0)	180.0	-0.0(-0.0)		

M0 UNDERFLOW AT 060067

ZN=2.24MEV G-MAX 1969-10-20 Y.TOMITA

PARAMETER SEARCH

NO.	*V	**	VSD	RR	RS	*A	*B	TOTAL CROSS SECTION	COMP.FORM. CROSS SECTION	ELASTIC CROSS SECTION	1ST LEVEL CROSS SECTION	CHI-SQUARE	TIME	
0	50.000	10.000	10.000	1.200	1.200	0.650	0.500	3.324	1.863	2.243	1.460	0.783	7.305* 22.7	
1*	53.231	10.413	10.000	1.200	1.200	0.615	0.611	3.467	1.953	2.346	1.514	0.832	2.689* 53.2	
2*	54.885	9.091	10.000	1.200	1.200	0.586	0.604	3.309	1.848	2.260	1.461	0.799	0.624* 83.2	
M0 UNDERFLOW AT	060067													
3*	54.801	9.338	10.000	1.200	1.200	0.597	0.594	3.312	1.865	2.252	1.447	0.805	0.688	0.610* 113.0
M0 UNDERFLOW AT	060067													
4*	54.801	9.299	10.000	1.200	1.200	0.596	0.595	3.312	1.863	2.254	1.449	0.805	0.687	0.609* 143.0
M0 UNDERFLOW AT	060067													

ZN=2.24MEV @-MAX

1969-10-20 Y.TOMITA

FINAL POTENTIAL PARAMETERS

POWERS OF E	0	1	2	3	4	5
* V	54.801	-0.000	-0.000	-0.000	-0.000	-0.000
* W	9.299	-0.000	-0.000	-0.000	-0.000	-0.000
VSO	10.000	-0.000	-0.000	-0.000	-0.000	-0.000
WR	1.200	-0.000	-0.000	-0.000	-0.000	-0.000
WS	1.200	-0.000	-0.000	-0.000	-0.000	-0.000
* A	0.596	-0.000	-0.000	-0.000	-0.000	-0.000
* R	0.595	-0.000	-0.000	-0.000	-0.000	-0.000

**S INDICATE THE PARAMETERS TO BE SEARCHED FOR

TOTAL CROSS SECTION..... 3.312 BARN
 ABSORPTION CROSS SECTION..... 1.863 BARN
 SHAPE ELASTIC CROSS SECTION.. 1.449 BARN
 CHI-SQUARE VALUE..... 0.609

TRANSMISSION COEFFICIENTS

LEVEL	GROUND	1ST	2ND	3RD
ECM (EXIT)	2.206	1.186	0.371	0.306
MESH	0.250	0.250	0.250	0.250
MATCH. RADIIIS	14.979	15.246	15.991	16.441
L	L-S			
0 +	7.2579E-01	6.7528E-01	5.1629E-01	4.8683E-01
1 -	6.3922E-01	4.4889E-01	1.5751E-01	1.2659E-01
1 +	7.8605E-01	5.8511E-01	2.1475E-01	1.7777E-01
2 -	5.5007E-01	2.9226E-01	3.1704E-02	2.0422E-02
2 +	3.4701E-01	1.6961E-01	1.9508E-02	1.2748E-02
3 -	8.5040E-02	1.2634E-02	2.6281E-04	1.3600E-04
3 +	2.0112E-01	2.7640E-02	5.1735E-04	2.6552E-04
4 -	9.7385E-03	7.5084E-04	0.	0.
4 +	7.0109E-03	5.5110E-04	0.	0.
5 -	3.0413E-04	1.2353E-05	0.	0.
5 +	6.7813E-04	2.6867E-05	0.	0.
6 -	1.0997E-05	0.	0.	0.
6 +	1.5602E-05	0.	0.	0.

ZN=2.24MEV @-MAX

1969-10-20 Y.TOMITA

S-MATRIX ELEMENTS FOR THE ELASTIC CHANNELS

L	L-S -		L-S +	
	REAL PART	IMAGINARY PART	REAL PART	IMAGINARY PART
0			-2.18351E-01	4.75951E-01
1	3.19318E-01	-5.08737E-01	3.17079E-01	-3.36771E-01
2	5.84469E-01	-3.29129E-01	7.22493E-01	-3.62065E-01
3	9.56496E-01	8.69566E-03	8.93220E-01	3.20154E-02
4	9.95119E-01	1.76826E-04	9.96486E-01	2.08633E-03
5	9.99848E-01	1.01511E-04	9.99641E-01	1.73409E-04
6	9.99994E-01	4.33369E-06	9.99992E-01	1.34409E-05

VARIOUS STATISTICAL QUANTITIES

CALCULATION OF COMPOUND PROCESS WAS CARRIED OUT BY USING MOLDAUER'S THEORY

N = 1.000

BC = 1.000

REAL PART OF THE RELATIVE PHASE OF SQUARES OF REDUCED WIDTH AMPLITUDES BETWEEN DIFFERENT CHANNELS.... 1.000
 DEGREE OF FREEDOM OF CHI-SQUARE DISTRIBUTION FOR SQUARES OF REDUCED WIDTH AMPLITUDE..... 1.000

TOTAL SPIN PARITY	TOTAL STRENGTH FUNCTION	COMPOUND FORMATION CROSS SECTION	W
1/2-	1.8829	0.192	1.002
1/2+	2.0458	0.217	0.963
3/2-	2.7752	0.471	0.822
3/2+	2.6063	0.330	0.851
5/2-	1.7579	0.076	1.033
5/2+	2.3614	0.312	0.897
7/2-	1.2413	0.241	1.189
7/2+	0.5782	0.012	1.483
9/2-	0.0419	0.000	1.926
9/2+	0.2124	0.011	1.740
11/2-	0.0293	0.001	1.945
11/2+	0.0013	0.000	1.996
13/2-	0.0000	0.000	2.000
13/2+	0.0006	0.000	1.998

ZN-2.24MEV

0-MAX

1969-10-20 Y.TOMITA

GROUND 0.000 MEV. 0 +

LAB. ANGLE	LAB. COSINE	C.M. ANGLE	C.M. COSINE	COMPOUND ELASTIC	SHAPE LLASTIC	TOTAL ELASTIC
0.00	1.000	0.00	1.000	135.10	852.59	987.69
10.00	0.985	10.15	0.984	120.40	800.28	928.68
20.00	0.940	20.30	0.930	111.45	660.24	771.69
30.00	0.866	30.44	0.862	91.37	475.11	566.48
40.00	0.766	40.57	0.760	74.36	294.34	368.70
50.00	0.643	50.68	0.634	62.56	155.05	217.61
60.00	0.500	60.77	0.488	54.77	72.00	126.77
70.00	0.342	70.83	0.320	49.18	39.08	88.25
80.00	0.174	80.87	0.154	45.38	38.12	83.50
90.00	0.000	90.89	-0.015	44.15	49.10	93.25
100.00	-0.174	100.87	-0.189	45.88	57.26	103.14
110.00	-0.342	110.83	-0.356	49.99	55.88	105.87
120.00	-0.500	120.77	-0.512	55.78	45.45	101.23
130.00	-0.643	130.68	-0.654	63.87	30.70	94.57
140.00	-0.766	140.57	-0.772	76.01	17.09	93.10
150.00	-0.866	150.44	-0.870	93.07	8.11	101.18
160.00	-0.940	160.30	-0.941	112.63	4.20	116.83
170.00	-0.985	170.15	-0.985	120.78	3.49	132.27
180.00	-1.000	180.00	-1.000	135.10	3.60	138.70

INTEGRATED CROSS SECTION..... 0.805 1,449 2,254

1ST 1.020 MEV. 2 +

LAB. ANGLE	LAB. COSINE	C.M. ANGLE	C.M. COSINE	CROSS SECTION	LAB. ANGLE	LAB. COSINE	C.M. ANGLE	C.M. COSINE	CROSS SECTION
0.00	1.000	0.00	1.000	54.17	100.00	-0.174	101.19	-0.194	54.13
10.00	0.985	10.21	0.984	54.33	110.00	-0.342	111.13	-0.361	54.46
20.00	0.940	20.41	0.937	54.73	120.00	-0.500	121.05	-0.516	54.87
30.00	0.866	30.60	0.861	55.12	130.00	-0.643	130.92	-0.655	55.19
40.00	0.766	40.78	0.757	55.28	140.00	-0.766	140.78	-0.775	55.27
50.00	0.643	50.92	0.630	55.14	150.00	-0.866	150.60	-0.871	55.08
60.00	0.500	61.05	0.484	54.78	160.00	-0.940	160.41	-0.942	54.69
70.00	0.342	71.13	0.323	54.37	170.00	-0.985	170.21	-0.985	54.32
80.00	0.174	81.19	0.153	54.07	180.00	-1.000	180.00	-1.000	54.17
90.00	0.000	91.21	-0.021	53.98					

INTEGRATED CROSS SECTION..... 0.687

ZN-2.24MEV

0-MAX

1969-10-20 Y.TOMITA

2ND 1.835 MEV. 2 +

LAB. ANGLE	LAB. COSINE	C.M. ANGLE	C.M. COSINE	CROSS SECTION	LAB. ANGLE	LAB. COSINE	C.M. ANGLE	C.M. COSINE	CROSS SECTION
0.00	1.000	0.00	1.000	23.01	100.00	-0.174	102.13	-0.210	22.26
10.00	0.985	10.37	0.984	22.99	110.00	-0.342	112.03	-0.375	22.36
20.00	0.940	20.74	0.935	22.94	120.00	-0.500	121.87	-0.528	22.50
30.00	0.866	31.08	0.856	22.86	130.00	-0.643	131.65	-0.665	22.64
40.00	0.766	41.39	0.750	22.74	140.00	-0.766	141.39	-0.781	22.77
50.00	0.643	51.62	0.620	22.59	150.00	-0.866	151.08	-0.875	22.88
60.00	0.500	61.87	0.471	22.44	160.00	-0.940	160.74	-0.944	22.95
70.00	0.342	72.03	0.309	22.32	170.00	-0.985	170.37	-0.986	22.99
80.00	0.174	82.13	0.137	22.24	180.00	-1.000	180.00	-1.000	23.01
90.00	0.000	92.16	-0.030	22.22					

INTEGRATED CROSS SECTION..... 0.283

3RD 1.900 MEV. 0 +

LAB. ANGLE	LAB. COSINE	C.M. ANGLE	C.M. COSINE	CROSS SECTION	LAB. ANGLE	LAB. COSINE	C.M. ANGLE	C.M. COSINE	CROSS SECTION
0.00	1.000	0.00	1.000	9.83	100.00	-0.174	102.34	-0.214	5.93
10.00	0.985	10.41	0.984	9.68	110.00	-0.342	112.23	-0.378	6.28
20.00	0.940	20.81	0.935	9.25	120.00	-0.500	122.06	-0.531	6.79
30.00	0.866	31.19	0.855	8.62	130.00	-0.643	131.82	-0.667	7.43
40.00	0.766	41.53	0.749	7.90	140.00	-0.766	141.53	-0.783	8.11
50.00	0.643	51.82	0.618	7.18	150.00	-0.866	151.19	-0.876	8.78
60.00	0.500	62.08	0.469	6.56	160.00	-0.940	160.81	-0.944	9.33
70.00	0.342	72.23	0.305	6.10	170.00	-0.985	170.41	-0.986	9.70
80.00	0.174	82.34	0.135	5.83	180.00	-1.000	180.00	-1.000	9.83
90.00	0.000	92.38	-0.041	5.78					

INTEGRATED CROSS SECTION..... 0.089

TIME REQUIRED 176.5 SEC.

END-OF-DATA ENCOUNTERED ON SYSTEM INPUT FILE.

Appendix B Symbolic listing of STAX 2


```

LINE=LINE+1
IF (LINE.LE.50) GO TO 1698
WRITE (6,600) TDN
LINE=1
1698 WRITE (6,1700) EXSW
1700 FORMAT (9X,95HEXPECTED VALUE OF CHI-SQUARE .....F12.1)
C
1910 CALL XSECT(NLO,NLD,PARA1,SOC,1,2)
1920 NDO=2
MIN=2
IF (SOC.LT.S8CH10) MIN=1
X=MIN-1
ID5=100
GO TO 700
C
2010 IF (MIN.EQ.1) GO TO 1010
2020 GO TO (2050,2080),MIN
2050 S80=S8C
S82=S8CH10
2060 G=0.5
GO TO 2100
2080 S80=S8CH10
S82=S8C
G=0.5
2100 LI=MIN-1
DO 2120 N=1,NPARA
I=IDPAR(N)
2110 PARA1(I)=PARA1(I)-0.5*DRP(N)
CALL XSECT(NLO,NLD,PARA1,SOC,1,2)
NDO=N
MIN=2
IF (SOC.LE.S80) MIN=1
X=FLOAT(MIN-1)*G
ID5=3
GO TO 700
2130 IF (MIN.EQ.1) GO TO 1010
C
2160 S81=S8C
A=0.5*(S82+S80)+S81
B=0.5*(S82-S80)
H=0
IF (B.LT.A.*A) H=(1-D,5*B/A)*G
DO 2200 N=1,NPARA
I=IDPAR(N)
2200 PARA1(I)=PARA1(I)-(G*H)*DRP(N)
CALL XSECT(NLO,NLD,PARA1,SOC,1,2)
IF (ITEST.EQ.0) GO TO 2270
LINE=LINE+1
IF (LINE.LE.58) GO TO 2250
WRITE (6,600) TDN
LINE=7
2250 WRITE (6,2260) A,B,H
2260 FORMAT (9X,2HA*,1PE11.3,5X,2HD*,E11.3,5X,2HM*,E11.3)
2270 NDO=N
ID5=2
IF (A.LE.0.) ID5=1
MIN=2
IF (SOC.LE.AMINI(S80,S81)) MIN=1
X=FLOAT(MIN-1)*(X-G*H)
GO TO 700
2290 IF (MIN.EQ.2) GO TO 2350
GO TO 1010
C
2350 S8C=AMINI(S80,S81)
2360 DO 2370 N=1,NPARA
I=IDPAR(N)
2370 PARA1(I)=PARA1(I)-X*DRP(N)
C
4000 IF ((S8C.LE.CHIMIN)-OR.(S8C.HI0-S8C.LE.CHIMIN+0.1)) RETURN
CALL CLOCK(T1)
IF (T1.GE.TEND) RETURN
4010 CONTINUE
C
4100 IF (NEND.EQ.1) RETURN
NEND=1
GO TO (4210,4310),IGO
4210 IGO=2
MNS= NSRCH
GO TO 1150
4310 IGO=1
MNS=ABS(NSRCH)
GO TO 1150
C
END
EIGENVALUE
C
SUBROUTINE EIGEN(M,A,D,U,F,NDIM)
C
A..... MATRIX TO BE DIAGONALIZED (SYMMETRIC)
C
D..... ACTUAL DIMENSION OF THE MATRIX
C
U(N,1): DIAGONALIZED MATRIX (DIAGONAL ELEMENTS ARE EIGENVALUES)
C
U(N,1): ORTHOGONAL MATRIX OF TRANSFORMATION (U*U=I)
C
F..... UPPER LIMIT FOR THE OFF-DIAGONAL ELEMENTS
C
NDIM..... MAXIMUM DIMENSION TAKEN IN THE MAIN PROGRAM
C
DIMENSION A(NDIM,NDIM),D(NDIM,NDIM),U(NDIM,NDIM)
C
G=0.
M1=M-1
DO 1010 I=1,M
DO 1020 J=1,M
D(I,J)=A(I,J)
U(I,J)=0.
IF (I.EQ.J) U(I,J)=1.
1010 CONTINUE
IF (M.EQ.1) GO TO 3000
2000 G=0.
DO 2010 I=1,M1
I1=I+1
DO 2010 J=1,M
G1=ABS(D(I1,J))
IF (G.GE.G1) GO TO 2010
K=1
L=J
G=0
2010 CONTINUE
IF (G.LE.F) GO TO 3000
DD=D(L,1)-D(K,K)
IF (2.*G.GT.ABS(DD)) GO TO 2110
T=2.*D(K,L)/DD
C=-T/(1.+SQRT(1.+T**2))
GO TO 2120
2110 T=DD/(2.+D(K,L))
C=T*SQRT(1.+T**2)*SIGN(1.,T)
2120 A=L./SQRT(1.+C**2)
G=C*G
C
DO 2210 I=1,M
UK=U(K,I)+U(L,I)
UL=U(L,I)+U(K,I)
U(I,I)=UK
2210 U(L,I)=UL
C
DO 2310 I=1,M
IF (I.EQ.K-OR.I.EQ.L) GO TO 2310
DK1=U(K,I)+U(L,I)
DL1=U(L,I)+U(K,I)
DK(I)=DK1
DL(I)=DL1
D(I,I)=DK(I)+DL(I)
2310 CONTINUE
DKK=D(K,K)+D(L,L)+C
DLL=D(L,L)+D(K,K)+C
DKK=DKK-D(K,L)+C
DLL=D(L,L)-D(K,L)+C
DK(K)=DKK
DL(L)=DLL
D(K,L)=D.
D(L,K)=D.
C
GO TO 2000
C
3000 F=G
RETURN
END
C
CROSS SECTION
C
SUBROUTINE XSECT(NLOUT,NLD,PARA1,S8C,IX,IFC)
C
DIMENSION BL(25,11),V(25),W(25),V50(25),RR(25),RS(25),A(25),
1 B(25),PAR(7),PL(11,19),PL1(11,19),PL2(11,19),S3,
3 TRMAT(25),TRPW(25)
C
COMPLEX WR,PHAIR(300),DPHAI,DL,YETA(11,2),YETAN,HI,HR
C
COMMON /CLEV/NLEV,EX(25),SPIN(25),NPAR(25),NLMAX(25),DJ,NLMAXA

```

```

SR 1750 1 /CSTA/T(11,2,25)+TH(11,2,25)+THTOT(21,2)+0(21,2)+10CONX X 10
SR 1760 2 @CONST+FNP+BC+REB+FCBI+COMP X 11
SR 1770 3 /CIN/EN(25)+FK(25)+HM+PLAM+PLAM0+PLAM1+FNMI3+COMPAI X 12
SR 1780 4 /CVA/REL/RE-DR+ORREL(19)+S1EX(19)+ERIA(19) X 13
SR 1790 5 /CPOT/PARA(7)+VE(5)+WE(5)+VSOE(5)+RRE(5)+RSE(5)+AE(5) X 14
SR 1800 6 BE(5) X 15
SR 1810 7 /CCRS/NANGL+STOT+SABS+SEL+SSEL+SG(25)+SELA(19)+SSELA(19) X 16
SR 1820 8 SGA(19,23)+PLANEL X 17
SR 1830 9 /CWEL/VR+VSDR+RV+RW+AR+VR+ER+DRI,Mesh,RPWR+RPWR1 X 18
SR 1840 COMMON /CFXP/IFEXP+WEIGHT+WGTOT+WGL+WGI X 19
SR 1850 1 WELA+WOLA+STOTEX+ERTOT+SELEX+EREL+SISTEX+ERIST X 20
SR 1860 2 /CWA/REL/RE-DR+ORREL(19)+S1EX(19)+ERIA(19) X 21
SR 1870 COMMON /CYET/YETA/COMJ/NJMIN+JMAX+NLMAX(3)+COSCM(19,25)+IFFL X 22
SR 1880 3 /CSGT/STOTT+SABST+SEL+VSELT+SG2T+SELA(19)+SSELA(19) X 22
SR 1890 4 SG2AT(19) X 23
SR 1900 5 /CCGW/CGL(11,11)+WL(11,2,11) X 23
C
CALL /NOVFL/NOVFL
IF (NOVFL.EQ.1) WRITE (6,98765)
98765 FORMAT (////5X,15H*****XSECT 1////)
NLMAXA=0 X 24
SG(2)=0 X 24
NLEV=NLEV X 25
IF (IFC.EQ.1) NLEV=1 X 25
DO 2230 N=1,NLEV X 25
ER=EN(N)*HM X 26
V(N)=PAR(1) X 26
W(N)=PAR(2) X 29
V50(N)=PAR(3) X 30
RR(N)=PAR(4) X 31
RS(N)=PAR(5) X 32
A(N)=PAR(6) X 33
B(N)=PAR(7) X 34
FE=1. X 35
DO 2010 NF=1,5 X 36
FE=FE*EN(N) X 37
V(N)=V(N)+FE*VE(N) X 38
W(N)=W(N)+FE*WE(N) X 39
V50(N)=V50(N)+FE*V50(N) X 40
RR(N)=RR(N)+FE*RR(N) X 41
RS(N)=RS(N)+FE*RS(N) X 42
A(N)=A(N)+FE*A(N) X 43
2010 B(N)=B(N)+FF*BE(N) X 44
VR=V(N)*HM X 45
WR=W(N)*HM*(0.,4.) X 46
V50R=V50(N)*HM*COMPAI/A(N) X 47
RW=RR(N)*FNMI3 X 48
RWS=RS(N)*FNMI3 X 49
AR=A(N) X 50
BR=B(N) X 51
IF (IX.EQ.0) GO TO 2030 X 51
VRX=MAX1(CABS(V(N)*A(N)*FK(N))+ABS(V50(N)*10.*COMPAI/RV)) X 52
IF (VRX.EQ.0.) GO TO 2020 X 52
R1=VRX/((FK(N)*PHIPRE)+A(N)*RV X 53
GO TO 2021 X 53
2020 R1=0. X 53
2021 IF (W(N)+E0.-OR.B(N)+E0.-N.) GO TO 2022 X 54
R2=ALOG(ABS(W(N)+B(N)+FK(N)+A(N)+/CFN(N)*PHIPRE)+A(N)*RV X 54
GO TO 2023 X 54
2022 R2=0. X 54
2023 RMATCH=MAX1(R1,R2) X 54
RPWR=1.2/SBRT(ER+VR) X 56
TRMAT(N)=RMATCH X 56
TRPW(N)=RPWR X 56
GO TO 2040 X 56
2030 RMATCH=TRMAT(N) X 56
RPWR=TRPW(N) X 56
2040 MESH=(RMATCH+RPWR)/PR-2.0999999 X 56
IF (MESH.LE.300) GO TO 2050 X 58
MESH=300 X 59
DRI=(RMATCH+RPWR)/298. X 60
GO TO 2110 X 61
2050 RMATCH=DRFLOAT(MESH-2)*RPWR X 62
DRI=DR X 63
2110 ROH2=RMATCH*FK(N) X 64
ROH1=ROH2-DRI*FK(N) X 65
RPWR1=RPWR-DRI X 65
NLC=ROH2+1.5 X 65
RL=MAX1(RV+0.6*AR,RW+0.6*RR) X 65
ZL=RL*FK(N) X 65
BL=ZL*(ZL)/ZL X 65
NLO=COS(71)/ZL X 65
BMIN=0./((ZL-664*RL**2) X 65
DO 2112 NL=2,11 X 66
BS=FLOAT(2*NL-3)*BL/ZL-BLO X 66
IF (BS**2>FLOAT(2*NL-1).LE.BJMIN+AND.NL.GT.NLC) GO TO 2113 X 66
ALD=BL1 X 66
NL=12 X 66
2113 NL=NL-1 X 67
CALL /NOVFL/NOVFL X 67
98766 FORMAT (////5X,15H*****XSECT 2////) X 67
IF (NOVFL.EQ.1) WRITE (6,98766) X 67
NLMAX(N)=MNL X 67
T(1,1)=0. X 67
DO 2170 NL=1,MNL X 67
L=NL-1 X 68
CALL BESSEL(NL,BJ2*RN2/ROH2) X 69
2115 CALL BESSEL(NL,BJ1*RN1/ROH1) X 70
LS=1 X 71
IF (L.EQ.0) LS=2 X 72
2120 CALL DIFE0(L,LS,PHAIR) X 73
DPHAI=PHAIR(MESH)*ROH1/(PHAIR(MESH-1)*ROH2) X 74
HR=BJ2*DPHAI*BJ1 X 75
HI=(0.-1.)*(BN2+DPHAI*BN1) X 80
YETAN=(HI+HR)/(HI+HR) X 81
T(NL,LS)=1-(CABS(YETAN))**2 X 82
IF (N.EQ.1) YETA(N,LS)=YETAN X 83
IF (LS.FQ.2) GO TO 2170 X 84
LS=2 X 85
GO TO 2120 X 85
2170 CONTINUE X 92
2200 NLMAXA=MAX0(NLMAXA,NLMAX(N)) X 96
IF (NLMAX(N).FM.11) GO TO 2220 X 97
NL=NLMAX(N)+1 X 98
DO 2210 NL=NL+1 X 99
T(NL,1)=0. X 100
2210 T(NL,2)=0. X 101
2220 RMATCH=RMATCH X 102
DRN(N)=DRI X 104
2230 CONTINUE X 105
IF (NOVFL/NOVFL) X 105
IF (NOVFL.EQ.1) WRITE (6,98767) X 105
98767 FORMAT (////5X,15H*****XSECT 3////) X 105
C
IF (IFC.EQ.1) GO TO 3010 X 105
FNL=NLE=1 X 106
DO 2310 L=1,NL X 106
FL=2*L-1 X 106
2310 CGL(NL,L)=CGL(FNL,FNL*FL)*(2.*FNL*1.) X 106
DO 2330 LS=1,2 X 106
L1=NL-LS X 106
IF (L1.EQ.0) GO TO 2330 X 106
FLS=FNL+FLOAT(LS)-1.5 X 106
DO 2320 L1=1,L1 X 107
FL=2*L1-1 X 107
2320 WL(NL,LS,L)=WRAC(FNL,FLS,FNL,FLS,0.,FL) X 107
2330 CONTINUE X 107
2340 CONTINUE X 107
IF (NOVFL/NOVFL) X 107
98768 FORMAT (////5X,15H*****XSECT 4////) X 107
98768 FORMAT (////5X,15H*****XSECT 4////) X 107
C
3010 MAX1=NLMAX(1) X 107
SARS=0. X 109
FNL=0. X 110
DO 3020 NL=1,MAX1 X 110
FNL=FNL X 112
FNL=FNL+1. X 113
3020 SARS=SARS+FNL*1*(NL+1)+FNL*(NL+2) X 114
SARS=SARS*PLAM X 115
C
IF (COMP.FQ.0.) GO TO 3700 X 116
IF (IFC.EQ.1) GO TO 6010 X 116
DO 3510 NF=1,NLOUT X 117
DO 3510 NI=1,MAX1 X 118
3510 BL(NF,NL)=0. X 119
NJMAX=MAX1+INT(SPIN(1)-.4*DJ) X 120
F=FLOAT(MAX1)-SPIN(1)-1. X 121
NJMIN=0.-A*F X 122
IF (NJMIN.LE.0) NJMIN=1 X 123
DO 3610 N=1,NJMIN X 124
CALL /NOVFL/NOVFL X 124
98705 FORMAT (////5X,15H*****XSECT 5////) X 124
IF (NOVFL.EQ.1) WRITE (6,98705) X 124
CALL TRANS(IMAGY+NU+NNNNNN) X 126

```

```

CALL OVERFL(NOVFL)
IF (NOVFL.EQ.1) WRITE (6,98706)
98706 FORMAT (//////X,15H*****XSECT 6////)
3610 CALL NLOUT.NLUD.NJ.NL)
C
DO 3670 N=1,NLOUT
  X=MINO(N,33)
  SG(N)=PLAM1*NL(N,1)+COMP
  IF (N.GT.NLUD) GO TO 3670
  NL2=MINO(NLMAX(N)+MAXL1)
  NL1=1
  IF (N.LT.EQ.33) GO TO 3620
  IF (NL2.LE.NLMAX(N)) GO TO 3660
  NL1=NLMAX(N)+1
  IF (N.LT.EQ.23) NLMAX(N)=NL2
3620 CALL LEGND2(NL1,NL2,NANGL,PL2,COSCH,N,IFFL)
  CALL OVERFL(NOVFL)
  IF (NOVFL.EQ.1) WRITE (6,98707)
98707 FORMAT (//////X,15H*****XSECT 7////)
3640 DO 3660 NA=1,NANGL
  SGA(NA,N)=0.
  DO 3650 NL=1,NL2
  SGA(NA,NL)=SGA(NA,N)+NL(N,1)*PL2(NL,NA)*K
3650 SGA(NA,NL)=SGA(NA,N)+NL(N,1)*PL2(NL,NA)*K
3660 SGA(NA,N)=SGA(NA,N)+PLAMPCOMP
3670 CONTINUE
C
3700 IF (NLMAX(1).LE.NLMAX(2)) GO TO 3701
  CALL LEGND1(NLMAX(1)+1,NLMAX(1)+NANGL,PL,PL1,COSCH)
  CALL OVERFL(NOVFL)
  IF (NOVFL.EQ.1) WRITE (6,98708)
98708 FORMAT (//////X,15H*****XSECT 8////)
  NLMAX(1)=NLMAX(2)
3701 CALL SHAPE(FK(1),NANGL,MAXL1,YETA,SSELA,SSEL,PL,PL1,PLAMEL,PLAMOX)
  CALL OVERFL(NOVFL)
  IF (NOVFL.EQ.1) WRITE (6,98709)
98709 FORMAT (//////X,15H*****XSECT 9////)
  DO 3710 NA=1,NANGL
  IF (COMP.EQ.0) SGA(NA,1)=0.
  SELA(NA)=SSELA(NA)+SGA(NA,1)
  IF (COMP.EQ.0) SG(1)=0.
  SEL=SELA+SG(1)
  STOT=SARS+SELA
C
IF (IFEXP.EQ.0) RETURN
SWE=0.
IF (WGTOT.EQ.0) GO TO 3810
SSE=SSE+WGTOT*(STOT-STOTF)/ERTOT)**2
3810 IF (WGLE.EQ.0) GO TO 3820
SSE=SSE+WGLE*(SEL-SELE)/EREL)**2
3820 IF (WGLA.EQ.0) GO TO 3830
SSE=SSE+WGLA*(SG(2)-SISTEX)/ERIST)**2
3830 IF (WGLEA.EQ.0) GO TO 3840
SSE=SSE+WGLEA*(SELA(NA)-SELAFX(NA))/ERELA(NA)**2
  DO 3840 NA=1,NANGL
  IF (SELAEX(NA).LE.0.) GO TO 3840
  SSE=SSE+(SELA(NA)-SELAFX(NA))/ERELA(NA)**2
3840 CONTINUE
SSE=SSE+WGLFA*SSA
3850 IF (WGLA.EQ.0) GO TO 3870
  SGA=0.
  DO 3860 NA=1,NANGL
  IF (SIAEX(NA).LE.0.) GO TO 3860
  SGA=SSE*(SGA(NA)-SIAEX(NA))/ERIA(NA)**2
3860 CONTINUE
SSE=SSE+WGLA*SSA
3870 SSK=SSE/WIGHT
  CALL OVERFL(NOVFL)
  IF (NOVFL.EQ.1) WRITE (6,98710)
98710 FORMAT (//////X,15H*****XSECT 10////)
  I=0
  RETURN
C
6010 RATS=SARS/SARST
  SG(1)=SCFLT*RATS
  SG(2)=SG2T*RATS
  IF (WGLEA.EQ.0) GO TO 6030
  DO 6020 NA=1,NANGL
  SGA(NA,1)=SCELAT(NA)*RATS
  6020 SGA(NA,1)=SCELAT(NA)*RATS
  6030 IF (WGLA.EQ.0) GO TO 3700
  DO 6040 SGA(NA,2)=SGPAT(NA)*RATS
  GO TO 3700
C
END
C
SOLUTION OF SCHROEDINGER EQUATION
C
SUBROUTINE DIFEQ(L,LS,PHAIR)
C
COMPLEX PHAIR(300),WR,P(8),UO(8),FO(300),F(300),C(8)
C
COMMON /CDEL/WR,VR,VDR,VRH,RS,AR,B,EN,DR,MN1,RPWR,RPWR1
C
COMMON /COUT/P
DIMENSION FSO(300),FR(10),G(8),FR(9),H(8)
DOUBLE PRECISION FR,G,FB,FEV,EW,EA,ED,DEL,DE2
C
CALL OVERFL(NOVFL)
IF (NOVFL.EQ.1) WRITE (6,98765)
98765 FORMAT (//////X,15H*****XSECT 11////)
GO TO (10,1020),LS
1010 FACL=L-1
  RST1=RST1*RPWR1
  RST2=RST2*RPWR
  FZL=FZL*Z,
  GO TO 1030
1020 FACL=L
  FLL=FACL*(FACL+1.)
  DR1=DR+2*FLL
  FACL=FACL+VSOR
C
STARTING VALUES
C
IF (L.NE.0) GO TO 1200
  FV=DFXP(DRLE(-HR/AR))
  FV=DFXP(DRLE(-HS/AS))
  RST1=RPWR1
  RST2=RPWR
  FZL=Z,
  FR(1)=1./L*(1.+EV)
  FR(2)=1./L*(1.+E*)
  FN=0.
  DO 1120 N=2,9
  FN=FN+1.
  FN1=N-1
  FR(N)=FR(N-1)
  GN1=FR(N1)
  DO 1110 M=1,N1
  MN=M
  FR(N)=FR(N)+FR(M)*FR(NM)
  1110 GN1=GN1+FR(M)*FR(NM)
  FR(N)=FR(N)+FR(M)
  1120 FB(N)=GN1/(N*FN)
  FR(10)=FR(9)
  DO 1130 M=1,9
  FN1=M
  FR(10)=FR(10)+FR(M)*FR(M10)
  FR(10)=FR(10)/FR(M)
  FN1=M
  DO 1140 N=1,8
  FN=FN+1.
  FN1=N-1
  C(N)=VR+SNGL(FR(N))+WR*SNGL(G(N))
  H(N)=WR+VR*SNGL(FR(N)+2)
  CALL OVERFL(NOVFL)
  IF (NOVFL.EQ.1) WRITE (6,98766)
98766 FORMAT (//////X,15H*****XSECT 12////)
  C(1)=C(1)+ER
C
1200 RESID=-FACL*FR(2)*AR
  DO 1210 N=1,8
  UO(N)=C(N)+H(N)*FACL
  M(1)=1.
  M(2)=-RESID/FZL2
  FN1=1
  DO 1320 N=3,8
  M(N)=-RESID*M(N-1)
  FN=FN+1.
  N2=N-2
  DO 1310 M=1,N2
  MN=M-1
  1310 M(N)=M(N)+M(N)*UO(N)
  1320 M(N)=M(N)/CFNA(FZL2+FN-1.)
C
PHAIR(1)=M(1)*RST1
PHAIR(2)=M(2)*RST2
ST1=RST1
ST2=RST2
DO 1330 N=2,8
  ST1=ST1*RPWR1
  ST2=ST2*RPWR1
  DIF 740
  DIF 750
  DIF 760
  DIF 770
  DIF 780
  DIF 790
  DIF 800
  DIF 810
  DIF 820
  DIF 830
  DIF 840
  DIF 850
  DIF 860
  DIF 870
  DIF 880
  DIF 890
  DIF 900
  DIF 910
  DIF 920
  DIF 930
  DIF 940
  DIF 950
  DIF 960
  DIF 970
  DIF 980
  DIF 990
  DIF 1000
  DIF 1010
  DIF 1020
  DIF 1030
  DIF 1040
  DIF 1050
  DIF 1060
  RES 20
  RES 30
  RES 40
  RES 50
  RES 60
  RES 70
  RES 80
  RES 90
  RES 100
  RES 110
  RES 120
  RES 130
  RES 140
  RES 150
  RES 160
  RES 170
  RES 180
  RES 190
  RES 200
  RES 210
  RES 220
  RES 230
  RES 240
  RES 250
  RES 260
  RES 270
  RES 280
  RES 290
  RES 300
  RES 310
  RES 320
  RES 330
  RES 340
  RES 350
  RES 360
  RES 370
  RES 380
  RES 390
  RES 400
  RES 410
  RES 420
  RES 430
  RES 440
  RES 450
  RES 460
  RES 470
  RES 480
  RES 490
  RES 500
  RES 510
  RES 520
  RES 530
  RES 540
  RES 550
  RES 560
  RES 570
  RES 580
  RES 590
  RES 600
  RES 610
  RES 620
  RES 630
  RES 640
  RES 650
  RES 660
  RES 670
  RES 680
  RES 690
  RES 700
  RES 710
  RES 720
  RES 730
  RES 740
  RES 750
  RES 760
  RES 770
  RES 780
  RES 790
  RES 800
  RES 810
  RES 820
  RES 830
  RES 840
  RES 850
  RES 860
  RES 870
  RES 880
  RES 890
  RES 900
  RES 910
  RES 920
  RES 930
  RES 940
  RES 950
  RES 960
  RES 970
  RES 980
  RES 990
  RES 1000
  RES 1010
  RES 1020
  RES 1030
  RES 1040
  RES 1050
  RES 1060
  RES 1070
  RES 1080
  RES 1090
  RES 1100
  RES 1110
  RES 1120
  RES 1130
  RES 1140
  RES 1150
  RES 1160
  RES 1170
  RES 1180
  RES 1190
  RES 1200
  RES 1210
  RES 1220
  RES 1230
  RES 1240
  RES 1250
  RES 1260
  RES 1270
  RES 1280
  RES 1290
  RES 1300
  RES 1310
  RES 1320
  RES 1330
  RES 1340
  RES 1350
  RES 1360
  RES 1370
  RES 1380
  RES 1390
  RES 1400
  RES 1410
  RES 1420
  RES 1430
  RES 1440
  RES 1450
  RES 1460
  RES 1470
  RES 1480
  RES 1490
  RES 1500
  RES 1510
  RES 1520
  RES 1530
  RES 1540
  RES 1550
  RES 1560
  RES 1570
  RES 1580
  RES 1590
  RES 1600
  RES 1610
  RES 1620
  RES 1630
  RES 1640
  RES 1650
  RES 1660
  RES 1670
  RES 1680
  RES 1690
  RES 1700
  RES 1710
  RES 1720
  RES 1730
  RES 1740
  RES 1750
  RES 1760
  RES 1770
  RES 1780
  RES 1790
  RES 1800
  RES 1810
  RES 1820
  RES 1830
  RES 1840
  RES 1850
  RES 1860
  RES 1870
  RES 1880
  RES 1890
  RES 1900
  RES 1910
  RES 1920
  RES 1930
  RES 1940
  RES 1950
  RES 1960
  RES 1970
  RES 1980
  RES 1990
  RES 2000
  RES 2010
  RES 2020
  RES 2030
  RES 2040
  RES 2050
  RES 2060
  RES 2070
  RES 2080
  RES 2090
  RES 2100
  RES 2110
  RES 2120
  RES 2130
  RES 2140
  RES 2150
  RES 2160
  RES 2170
  RES 2180
  RES 2190
  RES 2200
  RES 2210
  RES 2220
  RES 2230
  RES 2240
  RES 2250
  RES 2260
  RES 2270
  RES 2280
  RES 2290
  RES 2300
  RES 2310
  RES 2320
  RES 2330
  RES 2340
  RES 2350
  RES 2360
  RES 2370
  RES 2380
  RES 2390
  RES 2400
  RES 2410
  RES 2420
  RES 2430
  RES 2440
  RES 2450
  RES 2460
  RES 2470
  RES 2480
  RES 2490
  RES 2500
  RES 2510
  RES 2520
  RES 2530
  RES 2540
  RES 2550
  RES 2560
  RES 2570
  RES 2580
  RES 2590
  RES 2600
  RES 2610
  RES 2620
  RES 2630
  RES 2640
  RES 2650
  RES 2660
  RES 2670
  RES 2680
  RES 2690
  RES 2700
  RES 2710
  RES 2720
  RES 2730
  RES 2740
  RES 2750
  RES 2760
  RES 2770
  RES 2780
  RES 2790
  RES 2800
  RES 2810
  RES 2820
  RES 2830
  RES 2840
  RES 2850
  RES 2860
  RES 2870
  RES 2880
  RES 2890
  RES 2900
  RES 2910
  RES 2920
  RES 2930
  RES 2940
  RES 2950
  RES 2960
  RES 2970
  RES 2980
  RES 2990
  RES 3000
  RES 3010
  RES 3020
  RES 3030
  RES 3040
  RES 3050
  RES 3060
  RES 3070
  RES 3080
  RES 3090
  RES 3100
  RES 3110
  RES 3120
  RES 3130
  RES 3140
  RES 3150
  RES 3160
  RES 3170
  RES 3180
  RES 3190
  RES 3200
  RES 3210
  RES 3220
  RES 3230
  RES 3240
  RES 3250
  RES 3260
  RES 3270
  RES 3280
  RES 3290
  RES 3300
  RES 3310
  RES 3320
  RES 3330
  RES 3340
  RES 3350
  RES 3360
  RES 3370
  RES 3380
  RES 3390
  RES 3400
  RES 3410
  RES 3420
  RES 3430
  RES 3440
  RES 3450
  RES 3460
  RES 3470
  RES 3480
  RES 3490
  RES 3500
  RES 3510
  RES 3520
  RES 3530
  RES 3540
  RES 3550
  RES 3560
  RES 3570
  RES 3580
  RES 3590
  RES 3600
  RES 3610
  RES 3620
  RES 3630
  RES 3640
  RES 3650
  RES 3660
  RES 3670
  RES 3680
  RES 3690
  RES 3700
  RES 3710
  RES 3720
  RES 3730
  RES 3740
  RES 3750
  RES 3760
  RES 3770
  RES 3780
  RES 3790
  RES 3800
  RES 3810
  RES 3820
  RES 3830
  RES 3840
  RES 3850
  RES 3860
  RES 3870
  RES 3880
  RES 3890
  RES 3900
  RES 3910
  RES 3920
  RES 3930
  RES 3940
  RES 3950
  RES 3960
  RES 3970
  RES 3980
  RES 3990
  RES 4000
  RES 4010
  RES 4020
  RES 4030
  RES 4040
  RES 4050
  RES 4060
  RES 4070
  RES 4080
  RES 4090
  RES 4100
  RES 4110
  RES 4120
  RES 4130
  RES 4140
  RES 4150
  RES 4160
  RES 4170
  RES 4180
  RES 4190
  RES 4200
  RES 4210
  RES 4220
  RES 4230
  RES 4240
  RES 4250
  RES 4260
  RES 4270
  RES 4280
  RES 4290
  RES 4300
  RES 4310
  RES 4320
  RES 4330
  RES 4340
  RES 4350
  RES 4360
  RES 4370
  RES 4380
  RES 4390
  RES 4400
  RES 4410
  RES 4420
  RES 4430
  RES 4440
  RES 4450
  RES 4460
  RES 4470
  RES 4480
  RES 4490
  RES 4500
  RES 4510
  RES 4520
  RES 4530
  RES 4540
  RES 4550
  RES 4560
  RES 4570
  RES 4580
  RES 4590
  RES 4600
  RES 4610
  RES 4620
  RES 4630
  RES 4640
  RES 4650
  RES 4660
  RES 4670
  RES 4680
  RES 4690
  RES 4700
  RES 4710
  RES 4720
  RES 4730
  RES 4740
  RES 4750
  RES 4760
  RES 4770
  RES 4780
  RES 4790
  RES 4800
  RES 4810
  RES 4820
  RES 4830
  RES 4840
  RES 4850
  RES 4860
  RES 4870
  RES 4880
  RES 4890
  RES 4900
  RES 4910
  RES 4920
  RES 4930
  RES 4940
  RES 4950
  RES 4960
  RES 4970
  RES 4980
  RES 4990
  RES 5000
  RES 5010
  RES 5020
  RES 5030
  RES 5040
  RES 5050
  RES 5060
  RES 5070
  RES 5080
  RES 5090
  RES 5100
  RES 5110
  RES 5120
  RES 5130
  RES 5140
  RES 5150
  RES 5160
  RES 5170
  RES 5180
  RES 5190
  RES 5200
  RES 5210
  RES 5220
  RES 5230
  RES 5240
  RES 5250
  RES 5260
  RES 5270
  RES 5280
  RES 5290
  RES 5300
  RES 5310
  RES 5320
  RES 5330
  RES 5340
  RES 5350
  RES 5360
  RES 5370
  RES 5380
  RES 5390
  RES 5400
  RES 5410
  RES 5420
  RES 5430
  RES 5440
  RES 5450
  RES 5460
  RES 5470
  RES 5480
  RES 5490
  RES 5500
  RES 5510
  RES 5520
  RES 5530
  RES 5540
  RES 5550
  RES 5560
  RES 5570
  RES 5580
  RES 5590
  RES 5600
  RES 5610
  RES 5620
  RES 5630
  RES 5640
  RES 5650
  RES 5660
  RES 5670
  RES 5680
  RES 5690
  RES 5700
  RES 5710
  RES 5720
  RES 5730
  RES 5740
  RES 5750
  RES 5760
  RES 5770
  RES 5780
  RES 5790
  RES 5800
  RES 5810
  RES 5820
  RES 5830
  RES 5840
  RES 5850
  RES 5860
  RES 5870
  RES 5880
  RES 5890
  RES 5900
  RES 5910
  RES 5920
  RES 5930
  RES 5940
  RES 5950
  RES 5960
  RES 5970
  RES 5980
  RES 5990
  RES 6000
  RES 6010
  RES 6020
  RES 6030
  RES 6040
  RES 6050
  RES 6060
  RES 6070
  RES 6080
  RES 6090
  RES 6100
  RES 6110
  RES 6120
  RES 6130
  RES 6140
  RES 6150
  RES 6160
  RES 6170
  RES 6180
  RES 6190
  RES 6200
  RES 6210
  RES 6220
  RES 6230
  RES 6240
  RES 6250
  RES 6260
  RES 6270
  RES 6280
  RES 6290
  RES 6300
  RES 6310
  RES 6320
  RES 6330
  RES 6340
  RES 6350
  RES 6360
  RES 6370
  RES 6380
  RES 6390
  RES 6400
  RES 6410
  RES 6420
  RES 6430
  RES 6440
  RES 6450
  RES 6460
  RES 6470
  RES 6480
  RES 6490
  RES 6500
  RES 6510
  RES 6520
  RES 6530
  RES 6540
  RES 6550
  RES 6560
  RES 6570
  RES 6580
  RES 6590
  RES 6600
  RES 6610
  RES 6620
  RES 6630
  RES 6640
  RES 6650
  RES 6660
  RES 6670
  RES 6680
  RES 6690
  RES 6700
  RES 6710
  RES 6720
  RES 6730
  RES 6740
  RES 6750
  RES 6760
  RES 6770
  RES 6780
  RES 6790
  RES 6800
  RES 6810
  RES 6820
  RES 6830
  RES 6840
  RES 6850
  RES 6860
  RES 6870
  RES 6880
  RES 6890
  RES 6900
  RES 6910
  RES 6920
  RES 6930
  RES 6940
  RES 6950
  RES 6960
  RES 6970
  RES 6980
  RES 6990
  RES 7000
  RES 7010
  RES 7020
  RES 7030
  RES 7040
  RES 7050
  RES 7060
  RES 7070
  RES 7080
  RES 7090
  RES 7100
  RES 7110
  RES 7120
  RES 7130
  RES 7140
  RES 7150
  RES 7160
  RES 7170
  RES 7180
  RES 7190
  RES 7200
  RES 7210
  RES 7220
  RES 7230
  RES 7240
  RES 7250
  RES 7260
  RES 7270
  RES 7280
  RES 7290
  RES 7300
  RES 7310
  RES 7320
  RES 7330
  RES 7340
  RES 7350
  RES 7360
  RES 7370
  RES 7380
  RES 7390
  RES 7400
  RES 7410
  RES 7420
  RES 7430
  RES 7440
  RES 7450
  RES 7460
  RES 7470
  RES 7480
  RES 7490
  RES 7500
  RES 7510
  RES 7520
  RES 7530
  RES 7540
  RES 7550
  RES 7560
  RES 7570
  RES 7580
  RES 7590
  RES 7600
  RES 7610
  RES 7620
  RES 7630
  RES 7640
  RES 7650
  RES 7660
  RES 7670
  RES 7680
  RES 7690
  RES 7700
  RES 7710
  RES 7720
  RES 7730
  RES 7740
  RES 7750
  RES 7760
  RES 7770
  RES 7780
  RES 7790
  RES 7800
  RES 7810
  RES 7820
  RES 7830
  RES 7840
  RES 7850
  RES 7860
  RES 7870
  RES 7880
  RES 7890
  RES 7900
  RES 7910
  RES 7920
  RES 7930
  RES 7940
  RES 7950
  RES 7960
  RES 7970
  RES 7980
  RES 7990
  RES 8000
  RES 8010
  RES 8020
  RES 8030
  RES 8040
  RES 8050
  RES 8060
  RES 8070
  RES 8080
  RES 8090
  RES 8100
  RES 8110
  RES 8120
  RES 8130
  RES 8140
  RES 8150
  RES 8160
  RES 8170
  RES 8180
  RES 8190
  RES 8200
  RES 8210
  RES 8220
  RES 8230
  RES 8240
  RES 8250
  RES 8260
  RES 8270
  RES 8280
  RES 8290
  RES 8300
  RES 8310
  RES 8320
  RES 8330
  RES 8340
  RES 8350
  RES 8360
  RES 8370
  RES 8380
  RES 8390
  RES 8400
  RES 8410
  RES 8420
  RES 8430
  RES 8440
  RES 8450
  RES 8460
  RES 8470
  RES 8480
  RES 8490
  RES 8500
  RES 8510
  RES 8520
  RES 8530
  RES 8540
  RES 8550
  RES 8560
  RES 8570
  RES 8580
  RES 8590
  RES 8600
  RES 8610
  RES 8620
  RES 8630
  RES 8640
  RES 8650
  RES 8660
  RES 8670
  RES 8680
  RES 8690
  RES 8700
  RES 8710
  RES 8720
  RES 8730
  RES 8740
  RES 8750
  RES 8760
  RES 8770
  RES 8780
  RES 8790
  RES 8800
  RES 8810
  RES 8820
  RES 8830
  RES 8840
  RES 8850
  RES 8860
  RES 8870
  RES 8880
  RES 8890
  RES 8900
  RES 8910
  RES 8920
  RES 8930
  RES 8940
  RES 8950
  RES 8960
  RES 8970
  RES 8980
  RES 8990
  RES 9000
  RES 9010
  RES 9020
  RES 9030
  RES 9040
  RES 9050
  RES 9060
  RES 9070
  RES 9080
  RES 9090
  RES 9100
  RES 9110
  RES 9120
  RES 9130
  RES 9140
  RES 9150
  RES 9160
  RES 9170
  RES 9180
  RES 9190
  RES 9200
  RES 9210
  RES 9220
  RES 9230
  RES 9240
  RES 9250
  RES 9260
  RES 9270
  RES 9280
  RES 9290
  RES 9300
  RES 9310
  RES 9320
  RES 9330
  RES 9340
  RES 9350
  RES 9360
  RES 9370
  RES 9380
  RES 9390
  RES 9400
  RES 9410
  RES 9420
  RES 9430
  RES 9440
  RES 9450
  RES 9460
  RES 9470
  RES 9480
  RES 9490
  RES 9500
  RES 9510
  RES 9520
  RES 9530
  RES 9540
  RES 9550
  RES 9560
  RES 9570
  RES 9580
  RES 9590
  RES 9600
  RES 9610
  RES 9620
  RES 9630
  RES 9640
  RES 9650
  RES 9660
  RES 9670
  RES 9680
  RES 9690
  RES 9700
  RES 9710
  RES 9720
  RES 9730
  RES 9740
  RES 9750
  RES 9760
  RES 9770
  RES 9780
  RES 9790
  RES 9800
  RES 9810
  RES 9820
  RES 9830
  RES 9840
  RES 9850
  RES 9860
  RES 9870
  RES 9880
  RES 9890
  RES 9900
  RES 9910
  RES 9920
  RES 9930
  RES 9940
  RES 9950
  RES 9960
  RES 9970
  RES 9980
  RES 9990
  RES 10000

```


