

PROF GROUCH-G  
A processing code for group  
constants for a fast reactor

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July 1970

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## PROF GROUCH-G

### A processing code for group constants for a fast reactor

#### Abstract

The code PROF GROUCH-G has been developed for producing fast reactor group constants by processing directly the nuclear data given by ENDF/A formats. Types of group constants obtained by the code are similar to the Russian group-constants set, but the processing and calculation methods have been improved in comparison with those of the conventional group-constants sets published so far. The characteristic features of the code are as follows :

- (1) All data of DCC 1, 2, 6, 9 and 10 formats can be treated.
- (2) Calculations of different nuclides, reaction and data types can be continuously carried out, and the respective results are obtained.
- (3) For the angular distribution data of elastic and inelastic scattering, the energy transfer probabilities are in detail calculated by using the conservation law of momentum and energy.
- (4) Smooth spectra  $\phi_s(E)$  used for averaging cross sections are fission spectrum and  $1/E^n$ . Composition dependence of cross sections is taken into account by using a spectrum  $\phi_s(E)(\sigma_i(E) + \sigma_0)^{-1}$ .

UKNDL and ENDF/B data also can be processed, because codes converting them into ENDF/A formats have been developed.

The PROF GROUCH-G is written for IBM-360/75 and FACOM-230/60.

November 1969

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## 高速炉用群定数作成コード：PROF GROUCH-G

### 要 旨

PROF GROUCH-G は、ENDF/A の形式で与えられている核データを直接処理することによって、高速炉用群定数を作成する為に開発された計算コードである。このコードによってソ連の群定数セットと類似の群定数が得られるが、核データの処理および計算法に改良が施されている。このコードの特徴は次の点にある。

- (1) DCC 1, 2, 6, 9 および 10 の形式で与えられている全てのデータを処理する。
- (2) 核種、反応および形式の異ったデータを連続的に処理してそれぞれに対応した結果が得られる。
- (3) 弾性および非弾性散乱が角度分布によって与えられている場合、そのエネルギー分布への変換は、運動量とエネルギーの保存則を用いて詳細に計算される。
- (4) 断面積の平均に用いるスペクトルは核分裂スペクトルと  $1/E^a$  を採用し、群定数の組成依存性はこれらのスペクトルに  $(\sigma_i + \sigma_0)^{-1}$  の補正を行うことによって考慮している。

現在 UKNDL および ENDF/B の形式のデータを ENDF/A の形式に変換する計算コードが用意されているので、これらのデータも PROF GROUCH-G によって処理することが出来る。なおこのコードは IBM-360/75 と FACOM-230/60 用に作成されている。

1969 年 11 月

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## 1. Introduction

Several number of nuclear data files for use in calculations of reactor constants have been developed hitherto. The ENDF/A<sup>1)</sup>, ENDF/B<sup>2)</sup> and UKNDL<sup>3)</sup> are typical files for the storage and retrieval of nuclear data. In the files, a huge amount of data for different reaction types of very many nuclides is stored in accordance with their characteristic format. It is desirable that the format is designed to be flexible for the storage of new data and the retrieval of old data, and for general utilization. The ENDF/A format was decided so as to satisfy these requirements.

Two cases might be considered for producing reactor constants by utilizing a nuclear data file. One is the case where the selected data from the file is once converted to an another particular format for a data-library for a spectrum code to calculate reactor constants such as MC<sup>2 4)</sup>, in which the library tape is produced from the ENDF/B data. The other is the case where the file is directly processed and reactor constants are produced by using an approximate weighting spectrum in a processing code. A set of the resultant reactor constants is similar to a usual group-constants set such as YOM<sup>5)</sup> or the Russian set<sup>6)</sup> and to the MUFT or GAM-type library.

In the latter case, the set obtained by producing independently group constants of individual materials may not quite exactly predict the nuclear characteristics of a real reactor system. However, such a set is very convenient for nuclear design calculations, preceding analyses of experiments and so on, if the accuracy of the set is tolerable or reasonable for those purposes. It is also easy to modify or improve worse data of the set by using new nuclear data of higher accuracy or a more exact weighting spectrum.

The code PROF GROUCH-G (the signature "G" means "GAM-type") attempts to produce easily group constants for a fast reactor directly from the ENDF/A file. The original form of PROF GROUCH-G is the PROF GROUCH-M<sup>7)</sup> which produces MUFT-type group constants. But the code has been largely modified and improved. The PROF GROUCH-G can treat all data of DCC 1, DCC 2, DCC 6, DCC 9 and DCC 10 formats. The weighting spectrum used is similar to that of the Russian set which takes approximately composition dependence of group constants into account.

## 2. Description of the Processing Code

### 2.1 Weighting Spectrum

The standard weighting spectrum is given by  $\phi_s(E) = 1/E^\alpha$  below the energy ESPEC and  $\phi_s(E) = A_0 e^{-E/A_1} \sinh \sqrt{A_2 E}$  (fission spectrum) above ESPEC, where ESPEC,  $\alpha$ ,  $A_0$ ,  $A_1$  and  $A_2$  are input data. The  $\phi_s$  is reasonable as a weighting spectrum in an energy range where cross sections are constant or smooth. However, in a case where cross sections vary rapidly especially in resonance energy regions, the  $\phi_s$  should bring errors in resultant group constants. In the

code, influence of variations of cross sections on a spectrum is taken into account by the function  $(\sigma_t^m(E) + \sigma_0^m)^{-1}$  like the Russian set, where  $\sigma_t^m(E)$  is a microscopic total cross section and  $\sigma_0^m$  the sum of macroscopic total cross sections of other materials per element  $m$ .

An effective cross section of the  $i$ -th group extending from  $E_{il}$  to  $E_{iu}$  is given by

$$\sigma_{ix}^m(\sigma_0^m) = \frac{\int_{E_{il}}^{E_{iu}} \sigma_x^m(E) \frac{\phi_s(E)}{\sigma_t^m(E) + \sigma_0^m} dE}{\int_{E_{il}}^{E_{iu}} \frac{\phi_s(E)}{\sigma_t^m(E) + \sigma_0^m} dE} \quad (1)$$

where the subscript  $x$  represents all reactions except for total cross section. An effective total cross section, which is used for the calculation of the transport cross section for diffusion coefficient, is given by

$$\sigma_{it}^m(\sigma_0^m) = \frac{\int_{E_{il}}^{E_{iu}} \frac{\phi_s(E)}{\sigma_t^m(E) + \sigma_0^m} dE}{\int_{E_{il}}^{E_{iu}} \frac{\phi_s(E)}{(\sigma_t^m(E) + \sigma_0^m)^2} dE} - \sigma_0^m \quad (2)$$

Generally, the  $\sigma_0^m$  is defined by

$$\sigma_0^m = \sum_{n \neq m} N^n \sigma_t^n / N^m \quad (3)$$

where  $N^n$  and  $N^m$  are atomic densities of materials  $n$  and  $m$ , and in the PROF GROUCH-G the  $\sigma_0^m$  is given as a parameteric input.

When only the  $\phi_s(E)$  is used as a weighting spectrum, effective cross sections for infinite dilution are obtained, i. e.,

$$\sigma_{ix\infty}^m = \frac{\int_{E_{il}}^{E_{iu}} \sigma_x^m(E) \phi_s(E) dE}{\int_{E_{il}}^{E_{iu}} \phi_s(E) dE} \quad (4)$$

where  $x$  represents all reactions together with total cross section and the  $\sigma_{ix\infty}^m$  is the same as the  $\sigma_{ix}^m(\infty)$ .

The choice between  $\phi_s(E)$  and  $\phi_s(E)(\sigma_t + \sigma_0)^{-1}$  is made by an input option. The factor for representing the composition dependence of cross sections is defined by

$$f_{ix}^m(\sigma_0^m) = \frac{\sigma_{ix}^m(\sigma_0^m)}{\sigma_{ix\infty}^m} \quad (5)$$

The factor  $f$  is sometimes called a self-shielding factor.

The equations (1), (2) and (4) are used in cases for processing one-dimensional data, but in cases for calculating transfer matrices also, the same weighting spectra can be chosen.

## 2. 2 Equations of Elastic and Inelastic Scattering Matrices

Elastic and inelastic scattering matrices are obtained by processing the set of cross section (DCC 1 data) and angular (DCC 2) or energy distribution data (DCC 9, DCC 10) or data on Legendre moments (DCC 6), which are read in immediately after DCC 1 data.

The process of the set of DCC 1 and DCC 9 or DCC 10 data and the numerical method are the same as that of PROF GROUCH-M. The scattering matrices are simply obtained from

$$\sigma^{i \rightarrow j} = \frac{\int_{\Delta E_j} dF_F \int_{\Delta E_i} dE_I \sigma(E_I) P(E_I \rightarrow E_F) \phi(E_I)}{\int_{\Delta E_i} \phi(E_I) dE_I} \quad (6)$$

where  $P(E_I \rightarrow E_F)$  is given according to DCC 9 or DCC 10 data formats, and  $\Delta E_i$  and  $\Delta E_j$  are

the energy intervals of the  $i$ - and  $j$ -th groups.

On the other hand, the numerical procedure for the set of DCC 1 and DCC 2 is rather complex, because the angular distribution  $P(E_I, \mu_c)$  or  $P(E_I, \mu_l)$  must be converted to the energy distribution, where  $\mu_c$  is the cosine of scattering angle in the center of mass system and  $\mu_l$  in the laboratory system. The conversion of  $P(E_I, \mu_c)$  or  $P(E_I, \mu_l)$  into  $P(E_I \rightarrow E_F)$  is made by use of the conservation law of energy and momentum. Both elastic and inelastic scattering matrices can be generally treated by the same formula except for consideration of the  $Q$ -value (or excited level energy) in inelastic scattering.

The general relations between the incident energy  $E_I$ , the final energy  $E_F$  and the cosine of scattering angle  $\mu_c$  or  $\mu_l$  are as follows.

$$\frac{E_F}{E_I} = \frac{\left( \sqrt{A^2 - 1 + \mu_l^2} - A(A+1) \frac{w_n}{E_I} + \mu_l \right)^2}{(A+1)^2} \quad (7)$$

$$\mu_l = \frac{W + \mu_c}{\sqrt{W^2 + 2W\mu_c + 1}} \quad (8)$$

where

$$W^2 = \frac{1}{A^2 - A(A+1) \frac{w_n}{E_I}} \quad (9)$$

and

$$\mu_l = \frac{1}{2}(A+1)\sqrt{\frac{E_F}{E_I}} - \frac{1}{2}(A-1)\sqrt{\frac{E_I}{E_F}} + \frac{A}{2} \frac{w_n}{E_I} \sqrt{\frac{E_I}{E_F}} \quad (10)$$

where  $A$  is atomic mass, and  $w_n$  is the excited energy of the  $n$ -th level for inelastic scattering and is set equal to zero for elastic scattering.

The angular distribution of the elastic and inelastic scattering cross sections ( $\sigma_s$  and  $\sigma_{in}$ ) can be represented as a series of Legendre functions:  $P_0(\mu_l) = 1$ ;  $P_1(\mu_l) = \mu_l$ ;  $P_2(\mu_l) = \frac{1}{2}(3\mu_l^2 - 1) \dots$ , and we shall write

$$Y_{x_1}(E_I, \mu_{x_2}) = \sum_{l=0}^{\infty} S_{x_1, l}(E_I \rightarrow E_F) P_l(\mu_l) \quad (11)$$

where the suffix  $x_1$  represents "s" or "in" and the  $x_2$  "l" or "c".

$Y_{x_1}(E_I, \mu_{x_2})$  is given by the product of  $\sigma_{x_1}$  from DCC 1 data and  $P(E_I, \mu_{x_2})$  from DCC 2.

In case that the angular distribution is given in the laboratory system,  $S_{x_1, l}$  is obtained by using Eqs. (7)~(10) and

$$\begin{aligned} S_{x_1, l}(E_I \rightarrow E_F) &= \frac{2l+1}{4\pi} \int_{-1}^{+1} Y_{x_1}(E_I, \mu_l) P_l(\mu_l) \\ &\quad \times \delta\left(E_F - \frac{E_I}{(A+1)^2} \left[ \sqrt{A^2 - 1 + \mu_l^2} - A(A+1) \frac{w_n}{E_I} + \mu_l \right]^2\right) d\mu_l \\ &= \frac{2l+1}{4\pi} Y_{x_1}(E_I, \mu_l(E_I, E_F)) P_l(\mu_l(E_I, E_F)) \\ &\quad \times \frac{(A+1)\sqrt{\frac{E_F}{E_I}} + (A-1)\sqrt{\frac{E_I}{E_F}} - A \frac{w_n}{E_I} \sqrt{\frac{E_I}{E_F}}}{4E_F} \end{aligned} \quad (12)$$

In the center of mass system,

$$\begin{aligned} S_{x_1, l}(E_I \rightarrow E_F) &= \frac{2l+1}{4\pi} \int_{-1}^{+1} Y_{x_1}(E_I, \mu_c) P_l(\mu_l) \frac{d\mu_c}{d\mu_l} \\ &\quad \times \delta\left(E_F - \frac{E_I}{(A+1)^2} \left[ \sqrt{A^2 - 1 + \mu_l^2} - A(A+1) \frac{w_n}{E_I} + \mu_l \right]^2\right) d\mu_l \\ &= \frac{2l+1}{4\pi} Y_{x_1}(E_I, \mu_c(E_I, E_F)) P_l(\mu_l(E_I, E_F)) \frac{[\sqrt{1 - W^2(1 - \mu_l^2)} + W\mu_l]^2}{\sqrt{1 - W^2(1 - \mu_l^2)}} \end{aligned}$$



$$\times \frac{(A+1)\sqrt{\frac{E_F}{E_I}} + (A-1)\sqrt{\frac{E_I}{E_F}} - A\frac{w_n}{E_I}\sqrt{\frac{E_I}{E_F}}}{4E_F} \quad (13)$$

Eq. (13) is reduced to a rather simple form in case of elastic scattering, i. e.,

$$S_{s,l}(E_I \rightarrow E_F) = \frac{2l+1}{4\pi} Y_s(E_I, \mu_c(E_I, E_F)) P_l(\mu_l(E_I, E_F)) \frac{(A+1)^2}{2AE_I} \quad (14)$$

By use of  $S_{x_1,0}(E_I \rightarrow E_F)$  and  $S_{x_1,1}(E_I \rightarrow E_F)$  are obtained  $\sigma_{x_1}^{i \rightarrow j}$ ,  $\mu_{x_1}^{i \rightarrow j}$  and  $\mu_{x_1}^i$ , i.e.,

$$\sigma_{x_1}^{i \rightarrow j} = 4\pi \int_{\Delta E_j} dE_F \int_{\Delta E_i} dE_I S_{x_1,0}(E_I \rightarrow E_F) \phi(E_I) / \int_{\Delta E_i} dE_I \phi(E_I) = 4\pi S_{x_1,0}(i \rightarrow j), \quad (15)$$

$$\mu_{x_1}^{i \rightarrow j} = \frac{\int_{\Delta E_j} dE_F \int_{\Delta E_i} dE_I \phi(E_I) \int_{-1}^{+1} d\mu_l \mu_l Y_{x_1}(E_I, \mu_{x_2})}{\int_{\Delta E_j} dE_F \int_{\Delta E_i} dE_I \phi(E_I) \int_{-1}^{+1} d\mu_l Y_{x_1}(E_I, \mu_{x_2})} = \frac{S_{x_1,1}(i \rightarrow j)}{3S_{x_1,0}(i \rightarrow j)} \quad (16)$$

and

$$\mu_{x_1}^i = \frac{\int_{E_F \leq E_I} dE_F \int_{\Delta E_i} dE_I \phi(E_I) \int_{-1}^{+1} d\mu_l \mu_l Y_{x_1}(E_I, \mu_{x_2})}{\int_{E_F \leq E_I} dE_F \int_{\Delta E_i} dE_I \phi(E_I) \int_{-1}^{+1} d\mu_l Y_{x_1}(E_I, \mu_{x_2})} = \frac{\sum_{j \geq i} S_{x_1,1}(i \rightarrow j)}{3 \sum_{j \geq i} S_{x_1,0}(i \rightarrow j)} \quad (17)$$

Eq. (17) is rewritten as

$$\mu_{x_1}^i = \frac{\sum_{j \geq i} \mu_{x_1}^{i \rightarrow j} \sigma_{x_1}^{i \rightarrow j}}{\sigma_{x_1}^i} \quad (18)$$

where

$$\sigma_{x_1}^i = \sum_{j \geq i} \sigma_{x_1}^{i \rightarrow j} \quad (19)$$

The  $\sigma_{x_1}^i$  should coincide with the  $\sigma_{ix}$  from Eq. (1), but because of an error introduced from numerical integrations, the normalization of above transfer matrices is carried out in the code, to make the sum of the matrices coincide with  $\sigma_{ix}$ .

For the angular distribution, Legendre coefficients of the probability distribution or coefficients of differential cross sections are sometimes given. In this case we must use DCC 6 type data. The  $Y_{x_1}(E_I, \mu_{x_2})$  is produced by the Legendre coefficients and cross sections. Then Eq. (12)~(17) are applied to this  $Y_{x_1}(E_I, \mu_{x_2})$ . However DCC 6 is designed to use only for elastic scattering.

When angular distribution data are processed, another transfer matrices  $\sigma_{x_1}^{i \rightarrow j}$  also is obtained, where the diagonal component  $\sigma_{x_1}^{i \rightarrow i}$  is given by

$$\sigma_{x_1}^{i \rightarrow i} = \sigma_{x_1}^{i \rightarrow i} - \sigma_{x_1}^i \mu_{x_1}^i \quad (20)$$

Equation (20) is proposed in the code PIXSE<sup>8)</sup>. The  $\sigma_{x_1}^{i \rightarrow j}$  is obtained only in card form.

## 2.3 Method of Processing

The processing is performed according to specifications of data, i. e., by distinguishing varieties of nuclides, reaction types, data types (here all types of DCC 1, DCC 2, DCC 6, DCC 9, and DCC 10) and so on. In TABLE 1 are listed types of reactions and data which are processed by the code.

Calculations of different nuclides, reaction types and data types can be continuously carried out and the corresponding results are obtained. Then it should be noticed that angular or energy distribution data (DCC 2, DCC 6, DCC 9 and DCC 10) must be always read subsequently to cross sections (DCC 1). Consequently, fission cross section is also necessary for the

calculation of fission spectrum  $\chi$  before energy distribution data, though fission cross section has no influence on the results. Total cross section also must be read before all cross sections, when  $\phi_s(\sigma_t + \sigma_0)^{-1}$  spectrum is used. In case where both angular and energy distribution data in the same energy interval are prepared in ENDF data, either the former or the latter should be chosen as input data.

The procedure for reading and examination of data and the rearrangement is almost the same as in PROF GROUCH-M<sup>7)</sup>, but the whole structure of the code has been quite changed, and some subroutines for numerical calculations have been considerably modified, removed and added. The block diagrams of processing data are shown in Figs. 1~4.

PROF GROUCH-G can read data of ENDF/A format either in the form of cards or directly from the ENDF/A library tape whose number for reading is specified by an input. The cards from the ENDF/A tape is prepared by use of DFSR 2<sup>17)</sup> code. PROF GROUCH-G can process UKNDL and ENDF/B data also, because codes converting them to ENDF/A format have been developed<sup>9)</sup>.

In order to calculate the products (cross section)  $\times$  (probability (here angular or energy distribution))  $\times$  (spectrum) or (cross section)  $\times$  (spectrum), it is necessary to rearrange energy points of data by using interpolation rules. For instance, the procedure of producing probabilities of DCC 2 data at the same energy points as those of cross sections (DCC 1) was given in PROF GROUCH-M (see Tables 2.2.3 and 2.2.4 in the previous report<sup>7)</sup>). Further description will not be repeated, because the procedure in PROF GROUCH-G is the same.

In PROF GROUCH-G the maximum number (M) of energy points admissible initially in each group (cf. PROF GROUCH-M) has been extended from 200 to 500 (note that the M does not contain the mesh points increased by interpolation). And the maximum number (LDB\*) of incident energy points at which probability is given, has been also increased to 150, and the maximum number of initial data points is 2500. The length of data block (LDB\*) in DCC 2 is 100. Therefore it is not necessary to divide the ENDF/A data of some nuclides listed in the previous report<sup>7)</sup> into a few subblocks. The UK data converted into ENDF/A format also can be processed without modifying the data block.

## 2.4 Method of Numerical Integrations

### 2.4.1 Calculations of Cross Sections

Data are rearranged in the same manner as those for which NDBA=1, i. e.,  $[E_1, \sigma_1, E_2, \sigma_2, \dots]$  (see ENDF format). Integrations of the product (cross section:  $\sigma$ )  $\times$  (spectrum:  $\phi$ ) and of  $\phi$  are performed over individual group intervals. The same energy mesh points as determined for interpolation of  $\sigma$  are used for integrations of  $\sigma\phi$  and  $\phi$ , in order to avoid an error brought by the difference between both mesh structures, especially for  $\phi = \phi_s(\sigma_t + \sigma_0)^{-1}$ .

The number of interpolated points between the initial energy points of cross section is determined as follows. Let  $E_l$  be the  $l$ -th initial energy point and  $\sigma_l$  be the value of cross section at  $E_l$ . If the quantity  $\gamma = |\sigma_{l+1} - \sigma_l| / (\sigma_{l+1} + \sigma_l)$  is greater than  $\epsilon$ , the number MESH of energy points interpolated between  $E_l$  and  $E_{l+1}$  is determined as the integer part of  $\{(\gamma/\epsilon) + 1.0\}$ . Final energy points  $E_{l'}$  which are used for integrations are determined by dividing equally

\* See ENDF/A format in BNL-8381

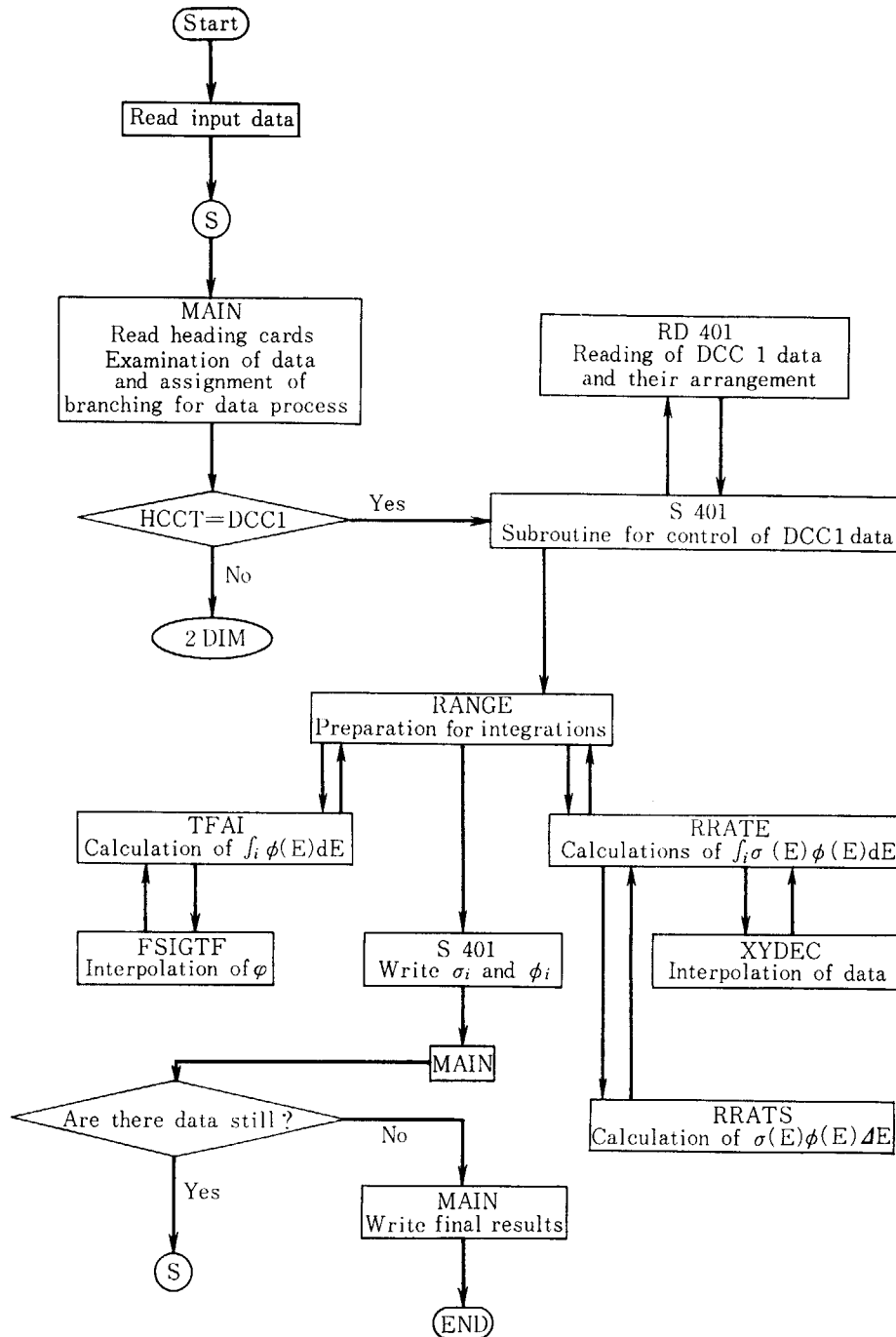


Fig. 1 Block diagram I for processing DCC1 data

in lethargy the interval between  $E_l$  and  $E_{l+1}$  into sub-intervals by MESH. The  $\epsilon$  is set equal to 0.1 in the code, by considering the computation time and the accuracy of integrations.

The integration over a sub-interval between  $E_l$  and  $E_{l+1}$  is given by an average of the values at  $E_l$  and  $E_{l+1}$ , i. e., by the trapezoid formula.

When reactions of the same nuclide are continuously processed, the sum of the reactions is obtained. For instance, when capture and fission cross sections are continuously processed, absorption cross section also is obtained as final results.

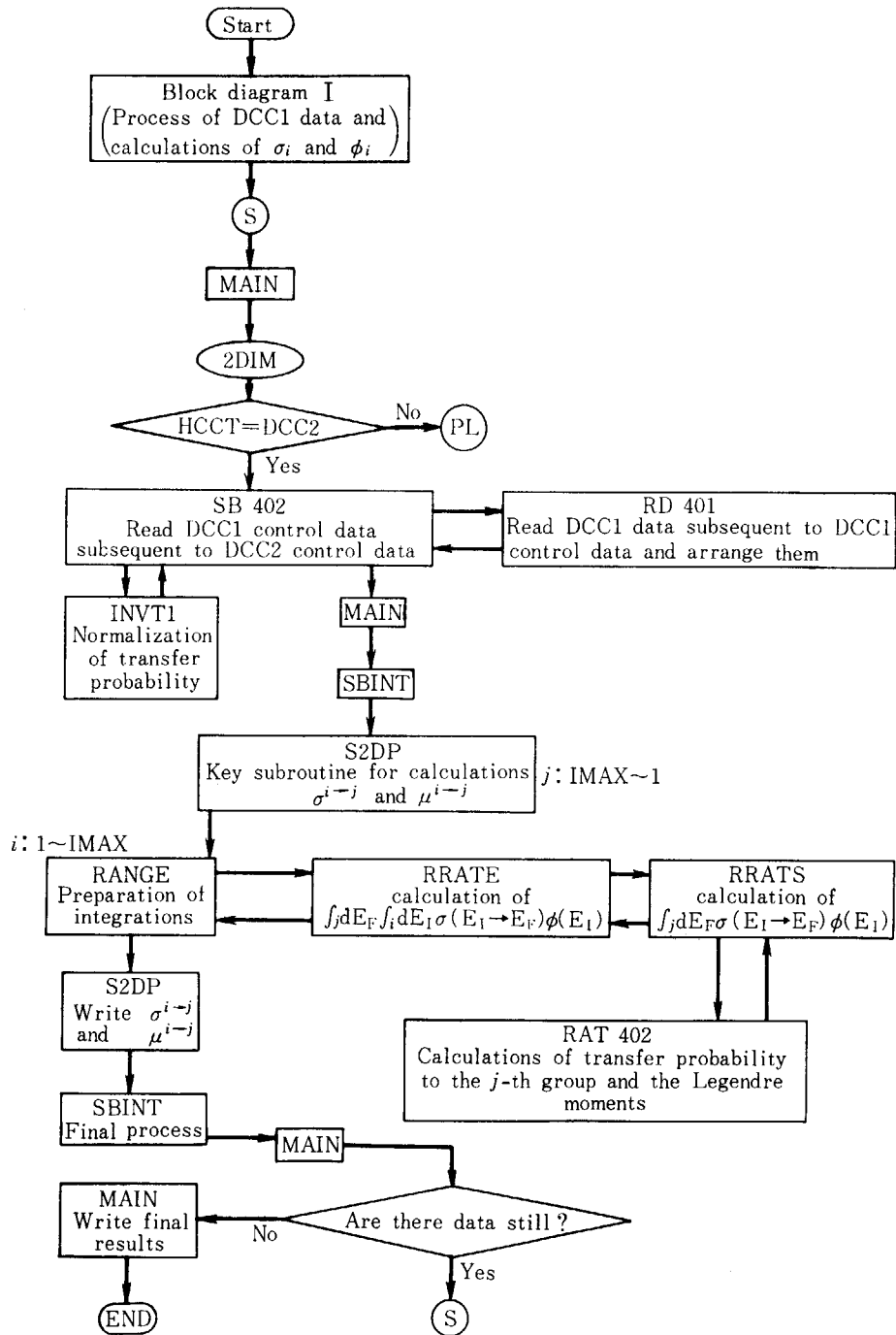


Fig. 2 Block diagram II for processing DCC 1 and DCC 2 data

2. 4. 2 Calculation of Transfer Matrices

The calculation of  $\int dE_F \int dE_I \sigma \cdot P \cdot \phi$  for obtaining the transfer matrix  $\sigma^{i \rightarrow j}$  in the code is performed by integrating it over all initial group  $i$  (the group containing incident energy) as fixing final group  $j$  (the group containing the energy of scattered neutron). The maximum number of the difference between the initial and final group is limited to 30 including the initial group itself, i. e., the maximum of  $j$  is  $i+29$ . This number is large enough for all cases except for hydrogen and the continuum level of inelastic scattering. However for the continuum level of inelastic scattering, errors brought by this truncation have been proved to be admissible in

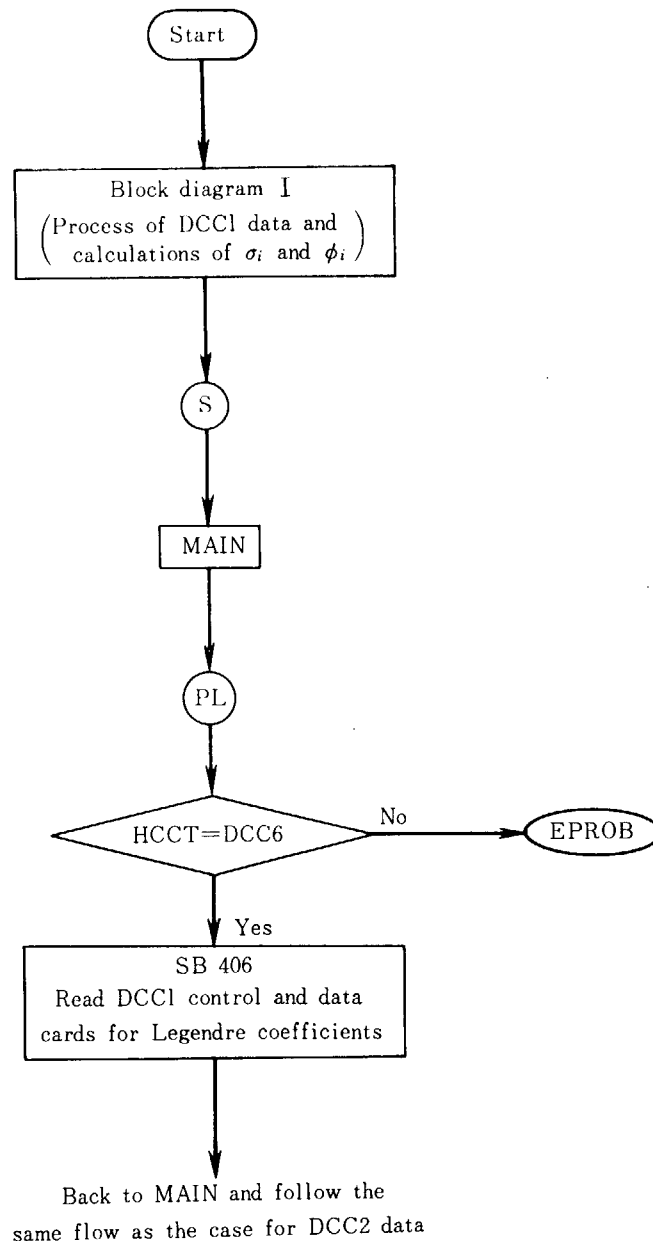


Fig. 3 Block diagram III for processing DCC 6 data

practice. Moreover it should be noticed that when  $j$  is larger than IMAX (maximum group number given by an input), the calculation of  $\sigma^{i \rightarrow j}$  is not carried out.

Before calculations of integrations probability is renormalized in the code, but because of an error brought by truncation and numerical procedure,  $\sum_j \sigma^{i \rightarrow j}$  does not precisely coincide with  $\sigma_i$  obtained from cross section data only. Hence the  $\sigma^{i \rightarrow j}$  is finally again renormalized to make  $\sum_j \sigma^{i \rightarrow j}$  equal to  $\sigma_i$ .

In the present code  $P_0$  and  $P_1$  components of Legendre moments can be obtained from angular distribution data. For higher moments only treatments of the final stage are not ready yet. For energy distribution data only  $P_0$  component is obtained.

The number MESH D of interpolated points between the initial incident energy points for the integration  $\int dE_F \int dE_I \sigma \cdot P \cdot \phi$  is determined as follows. In this case an another quantity besides the  $\gamma$  is necessary for reducing the error of the integration. Especially for the calculation

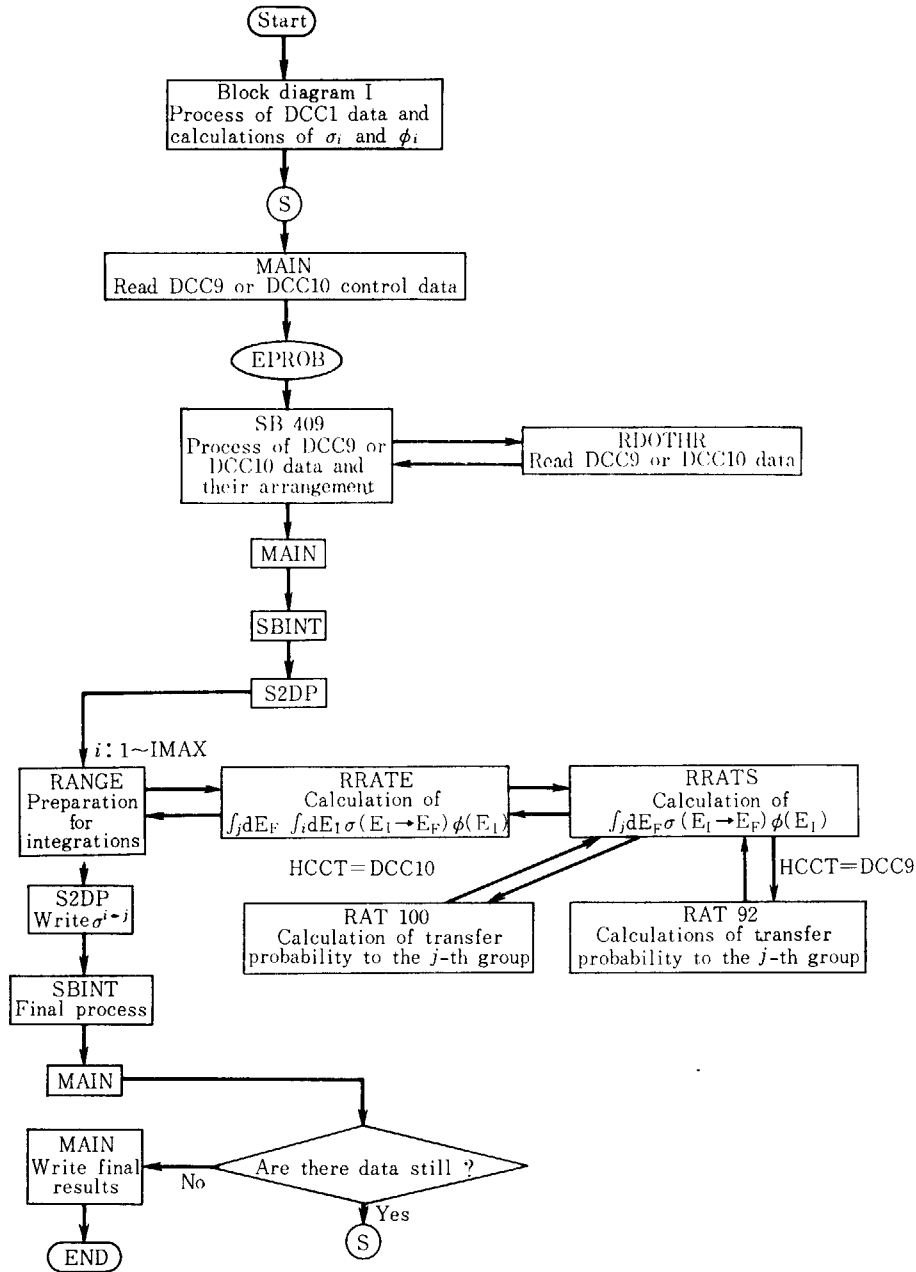


Fig. 4 Block diagram IV for processing DCC 1 and DCC 9 or DCC 10

of an elastic scattering matrix  $\sigma_s^{i \rightarrow J}$  ( $J < i$ ), we must carefully determined the number MESH in the energy interval from the lower energy boundary of the  $i$ -th group  $E_{Li}$  to  $E_{Li}/\alpha$ , from which scattered neutrons are slowed down into group  $J$ .  $\alpha$  is given by  $\{(A-1)/(A+1)\}^2$ . Even if scattering cross section is constant, it is necessary to determine the MESH depending on the value of  $\alpha$ . In the code the method determining MESH of the  $i$ -th group differs between above and below  $E_{Li}/\alpha$ .

If the energy point  $E_i$  is above  $E_{Li}/\alpha$ , MESH is the integer part of  $\{(E_{Li+1}-E_i)/(AMESH 2 \cdot E_{Li})+1\}$ , and if the  $E_i$  is below  $E_{Li}/\alpha$ , MESH takes a larger value of the integer part of  $\{(E_{Li+1}-E_i)/(AMESH 1 \cdot (1/\alpha-1) \cdot E_{Li})+1\}$  and the integer part of  $\{(Y/0.1)+1.0\}$ . AMESH 1 and AMESH 2 are given by inputs and when those are less than or equal to zero, values of AMESH 1=0.1 and AMESH 2=0.025 are set respectively in the code. Next the MESH below  $E_{Li}/\alpha$  is compared with MESH(I) of an input specified by a user and the larger value of both is adopted as final MESH.

In case of calculating inelastic scattering matrices, MESH D is the same as that obtained in 2.4.1.

It should be noticed that when both elastic and inelastic scattering for the same nuclide are processed, the sum of  $\sigma_s^{i \rightarrow j}$  and  $\sigma_{in}^{i \rightarrow j}$  is finally obtained in form of card and list. When we want  $\sigma_s^{i \rightarrow j}$  and  $\sigma_{in}^{i \rightarrow j}$  individually as final results, it is necessary to process both data separately.

### 3. Input Data and the Restrictions

Control input data before the ENDF/A data (or data converted to the ENDF/A format) are as follows.

# 1: Format (I1, 12 A 6); NFC, CCRD

NFC:  $\begin{cases} =0 & \text{indicates continuation of the card \# 1.} \\ \neq 0 & \text{indicates last information of the card \# 1.} \end{cases}$

CCRD: Title, comment etc.

# 2: Format (4 I 5); IMAX, LMAXI, ITAPE, MSIGS

IMAX: Number of groups ( $\leq 70$ ).

LMAXI: Number of Legendre moments required ( $\leq 2$ ).

ITAPE: If ITAPE  $\leq 0$ , the ENDF/A data are read in card form, and if ITAPE  $> 0$ , the data are read from a tape with the number specified by ITAPE.

MSIGS:  $\begin{cases} =0 & \text{indicates that the result for DCC 6 is cross section itself.} \\ \neq 0 & \text{indicates that the result for DCC 6 is normalized to unity.} \end{cases}$

# 3: Format (14 I 5); MESH (I), I=1, IMAX

MESH (I) is used for calculations of elastic scattering matrices and gives an admissible minimum number of interpolation between initial energy points of the input data contained below  $EL(I+1)/\alpha$  in the I-th group. The final interpolated number may increased more than MESH (I), according to the rule set in the code. If MESH (I)  $\leq 0$ , the interpolation is not performed. For MESH (I)  $> 30$ , the MESH (I) is set to be equal to 30 in the code.

# 4: Format (4 E 15. 8); EMAX, EMIN, ALPHA, EPS4

EMAX: Upper energy limit of the highest group.

EMIN: Lower energy limit of the lowest group.

ALPHA: Power  $\alpha$  to be used for standard weighting spectrum  $\phi_s = 1/E^\alpha$  below the energy ESPEC.

EPS4: Constant determining the integration interval of  $\mu$  for DCC6. If EPS4  $\leq 0$ , EPS4 is set equal to 1/16 in the code.

# 5: Format (6 E 12. 5); EL(I), I= 2, IMAX+1

EL (I+1) is the lower energy boundary of the I-th group. Note EL(1)=EMAX.

# 6: Format (4 E 15. 8); A0, A1, A2, ESPEC

A0: } Constants for fission spectrum,  $A_0 e^{-E/A_1} \sinh \sqrt{A_2 E}$ .

A1: } An arbitrary value does well for  $A_0$ , because  $A_0$  is modified after reading so that

A2: } at ESPEC the value of fission spectrum may be connected with that of  $1/E^\alpha$ .

ESPEC:  $\phi_s(E)$  is fission spectrum above the energy ESPEC and  $\phi_s(E) = 1/E^\alpha$  below ESPEC.

# 7: Format (A 3); ENUNI

- Energy unit used in the code, i. e., "EV", "MEV" etc.
- # 8: Format (3 I 5); NDIV, IPRT1, IPRT2
- NDIV:  $\begin{cases} =0 & \text{gives transfer matrix } \sigma^{i \rightarrow j}. \\ \neq 0 & \text{gives transfer probability } \sigma^{i \rightarrow j} / \Delta E_j. \end{cases}$
- IPRT1:  $\begin{cases} \neq 0 & \text{gives intermediate results of the coefficients of Legendre expansion.} \\ =0 & \text{does not give above results.} \end{cases}$
- IPRT2:  $\begin{cases} \neq 0 & \text{gives intermediate results of } \sigma^{i \rightarrow j} \text{ and } \mu^{i \rightarrow j}. \\ =0 & \text{does not give above results.} \end{cases}$
- # 9: Format (I 5, 3 E 15. 5); KCALF, SIGMA0, AMESH1, AMESH2
- KCALF:  $\begin{cases} \leq 0 & \text{for } \phi = \phi_s. \\ > 0 & \text{for } \phi = \phi_s (\sigma_t + \sigma_0)^{-1}. \end{cases}$
- SIGMA0: Value of  $\sigma_0$  in case of  $\text{KCALF} > 0$ .
- AMESH1: Constant used for determining the number of interpolation below  $E_{Li}/\alpha$  for the calculation of  $\sigma_s^{i \rightarrow j} (j > i)$ , where  $E_{Li}$  is the lower energy boundary of the  $i$ -th group. If  $\text{AMESH1} \leq 0$ , AMESH1 is set equal to 0.1 in the code.
- AMESH2: Constant used for determining the number of interpolation above  $E_{Li}/\alpha$  for the calculation of  $\sigma_s^{i \rightarrow j}$ .  
If  $\text{AMESH2} \leq 0$ , AMESH2 is set equal to 0.025 in the code.
- # 10: Format (2 A 5); ICALP, LCACP
- ICALP: ID number of 1st block of ENDF/A data processed.
- LCACP: ID number of 1st block of ENDF/A data processed.  
The ID number is specified in columns from 73 to 77.
- ENDF/A format data follow above input data.

## 4. Output Data

### 4. 1 Description of Output Data

The following results are obtained from the code.

1st page: List of input data

- 1) Title
- 2) Group boundaries in an energy unit: EL(I).
- 3) Maximum numbers of interpolated points for each group: MESH(I).
- 4) AMESH1 and AMESH2 used for determining the number of interpolation for the calculation of the elastic matrix  $\sigma_s^{i \rightarrow j}$ .
- 5) Coefficients of fission spectrum and ESPEC.
- 6)  $\sigma_0$  (in case of using  $\phi_s (\sigma_t + \sigma_0)^{-1}$  spectrum).  
2nd page: List of ENDF/A data
- 7) Heading and control data of ENDF/A formats.
- 8) First and last values of energy and its dependent variable of DCC 1 data.  
3rd page: calculation results of cross section (DCC 1 data)
- 9) Effective cross section.
- 10) Flux used for averaging.

In case of processing cross section data of various reactions continuously, the outputs of 7)



~10) are edited iteratively. When scattering probability data (DCC 2, DCC 6, DCC 9 and DCC 10) follow cross section data (DCC 1), the code prints :

- 11) Heading and control data of ENDF/A formats, and informations on DCC 2, DCC 6, DCC 9 or DCC 10 data.

If the data are DCC 2 type :

- 12) An arranged two-dimensional array of scattering probability and its renormalized array are listed.

Using the options IPRT1 and IPRT2, the following intermediate results are obtained :

- 13) The intermediate results of  $\sigma^{i \rightarrow j}$  and  $\mu^{i \rightarrow j}$ , and of their corresponding Legendre components. For DCC 9 and DCC 10,  $\mu^{i \rightarrow j}$  and the  $P_1$  component are not obtained.

Finally, the following results are obtained for every isotopes.

- 14) Final results.

The summation of all cross sections, except for total cross section, processed for the same isotope is obtained. When scattering probability data also are processed, the summation  $\sum_j \sigma^{i \rightarrow j}$ ,  $\mu$  (for DCC 2 and DCC 6) and  $\sigma^{i \rightarrow j}$  are listed.

Every cross sections, fluxes, the summation of cross sections for the same isotope,  $\mu$  (for DCC 2 and DCC 6),  $\sigma^{i \rightarrow j}$  and  $\sigma^{i \rightarrow j}$  (see Eq. (20)) are obtained also in card form.

## 4. 2 Examples of Output Data

- (1) Case of processing only cross section data (DCC 1)

TABLE 2 shows the results for a case where the data of all discrete and continuum levels of inelastic scattering of  $^{238}\text{U}$  have been continuously processed with using the standard weighting spectrum  $\phi_s(E)$ . The group structure of the example is similar to that of the Russian set. The "LOWER BOUNDARY" in the first page means ESPEC. The "SUMMATION OF ONE DIMENSIONAL DATA" in the final page gives the total inelastic scattering cross section of  $^{238}\text{U}$ . Cross sections of individual levels are listed in the pages following immediately after lists of heading cards of the corresponding levels.

In TABLE 3 are presented the results of total and elastic scattering cross sections of Fe in case of using the spectrum  $\phi_s(E)(\sigma_t(E) + \sigma_0)^{-1}$  for  $\sigma_0 = 10^2$  barns. The number of groups is 70 obtained by dividing each group of the Russian set into two or three groups. The "SUMMATION OF ONE DIMENSIONAL DATA" in the final page gives only the result of elastic scattering cross section.

- (2) Case of processing a set of cross section (DCC 1) and scattering probability data (DCC 2, DCC 6, DCC 9 and DCC 10)

In TABLE 4 are shown the results of elastic scattering of natural boron in case of processing cross section (DCC 1) and angular distribution data (DCC 2) with use of the spectrum  $\phi_s(E)$ . In the lists subsequent to the result of cross section, informations on DCC 2 data and the intermediate results of calculations are given. The intermediate results of the example have been obtained by setting IPRT1=0 and IPRT2=0. The symbols LGF, SIG and MU represent respectively the final group into which neutrons are scattered,  $\sigma_s^{i \rightarrow LGF}$  and  $\mu^{i \rightarrow LGF}$ . And the symbol LGI means the range of the initial group  $i$ .

Next to the above list the final results are shown by the titles subsequent to "SUMMATION OF ONE DIMENSIONAL DATA". The values of  $\sum \sigma^{i \rightarrow j}$ ,  $\mu^i$  and  $\sigma^{i \rightarrow j}$  of the first three

groups are set to be zero, because the scattering probability data of that energy region is not given in the input data converted from the UKNDL<sup>2)</sup>. Values in "SUMMATION OF TRANSFER MATRICES" indicate those before the normalization in order to know an error brought by numerical integrations.

For the other sets of cross section data (DCC 1) and DCC 6 or DCC 9 or DCC 10, almost the same results as for the set of DCC 1 and DCC 2 are also obtained.

### Acknowledgment

The authors wish to express their gratitude to users of the code for their comments on it.

### References

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TABLE 1 Data Processed by PROF GROUCH-G

Data form	Reaction type	Data type
C  DCC 1	Total	Cross section
	Elastic scattering	"
	Inelastic scattering	"
	Pair production	"
	Triplet production	"
	Fission	"
	Parasitic absorption	"
	(n, $\gamma$ )	"
	(n, p)	"
	(n, t)	"
	(n, $^3\text{He}$ )	"
	(n, $\alpha$ )	"
	(n, $2\alpha$ )	"
P (probability) DCC 2, DCC 6 DCC 9 and DCC 10	Elastic scattering	Angular and energy distribution
	Inelastic scattering	
	Pair production	
	Triplet production	
	Fission	

TABLE 2 Output for processing only DCC 1 data of inelastic scattering cross section of  $^{238}\text{U}$  with use of  $\phi_s(E)$

```

PRODUCTIVE CALCULATION OF *U238-INELASTIC*
GROUP STRUCTURE IS THE SAME AS ESELEM * 1/E WEIGHT *
DATA ARE TAKEN FROM A/W JUL./30/1968 T.TONE

EL= 0.10500000E 08 0.65000000E 07 0.40000000E 07 0.19951000E 07 0.13969000E 07 0.79882000E 06
EL= 0.40127000E 06 0.19994000E 06 0.99628000E 05 0.46152000E 05 0.21380000E 05 0.99843999E 04
EL= 0.46253000E 04 0.21500000E 04 0.10000000E 04 0.46500000E 03 0.21500000E 03 0.10000000E 03
EL= 0.46500000E 02 0.21500000E 02 0.10000000E 02 0.46500000E 01 0.21500000E 01 0.10000000E 01
EL= 0.46500000E 00 0.21500000E 00

MESH= 1 1 1 1 1 1
MESH= 1 1 1 1 1 1
MESH= 1 1 1 1 1 1
MESH= 1 1 1 1 1 1
MESH= 1
      AMESH1 = 0.10000
      AMESH2 = 0.02500

FISSION SPECTRUM
A0 = 0.48400E-06 A1 = 0.10000E 07 A2 = 0.20000E-05 LOWER BOUNDARY = 0.10000E 07 EV

U 238
U 238 N INELAS 1 N CROS C E1 BN A/W 41 0 63 1.000E-03 1.500E 07EV A0041 0
 50 1 92 238 1 4 1 1 1 1 1 2 A0041 1
 1 2 1 0 0 0 0 128 41 0 63 0 A0041 2
2.3807E 02 -4.4000E 04 0. 1.0000E 00 1.0000E-03 1.5000E 07 A0041 3
ALDERMASTON/WINFRITH DATA FILE. 7/1/64 NIN= 5 GCN= 1 PCN= 5 A0041 4
AEEW R351. 2/64. BARRINGTON ETAL. AWRE 0-79/63. 1/64. K.PARKER A0041 5
DCC1 401 1 5 0 124 0 0 -0. 0. A0041 6
0.40000000E 05 0.10000000E-02 0.20000000E 07 0.10000000E-02 (FIRST AND LAST VALUE)
NOTYP=1
    
```

\*\*\* U 238 (N-INFLAS) 1,N CROSS C(EI, , ) \*\*\*

SEC= 1 (LGI=1...25)  
 0. 0. 0.57476375E-04 0.66619598E-01 0.73184316E 00 0.15274226E 01  
 0.13310057E 01 0.89416164E 00 0.21200201E 00 0.48062913E-04 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

(TOTAL SEC) \*\*\* (LGI=1...25)  
 0. 0. 0.57476375E-04 0.66619598E-01 0.73184316E 00 0.15274226E 01  
 0.13310057E 01 0.89416164E 00 0.21200201E 00 0.48062913E-04 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

\*\*\* FLUX \*\*\*

0. 0. 8.4291242 E-01 4.7372668 E-01 5.8141658 E-01 6.8869995 E-01  
 6.9738338 E-01 6.9951868 E-01 7.6963085 E-01 7.7010402 E-01 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

U 238  
 U 238 N INELAS 2 N CROSS C EI BN A/W 44 0 63 1.000E-03 1.500E U7EV A0044 0  
 46 1 92 238 1 4 2 1 1 1 2 A0044 1  
 1 2 1 0 0 0 128 44 U 63 0 A0044 2  
 2.3807E 02 -1.4600E 05 0. 1.0000E 00 1.0000E-03 1.5000E 07 A0044 3  
 ALDERMASTON/WINFRITH DATA FILE: 7/1/64 NIN= 5 GCN= 1 PCN= 6 A0044 4  
 AEEW R351: 2/64, BARRINGTON ETAL. AWRE 0-79/63, 1/64, K-PARKER A0044 5  
 DCC1 401 1 5 0 113 0 0-0. 0. A0044 6  
 0.15000000E 06 0.10000000E-02 0.20000000E 07 0.10000000E-02 (FIRST AND LAST VALUE)  
 NDYTP=1

\*\*\* U 238 (N-INFLAS) 2,N CROSS C(EI, , ) \*\*\*

SEC= 1 (LGI=1...25)  
 0. 0. 0.57653844E-04 0.29017699E 00 0.58414833E 00 0.28631179E 00  
 0.58651783E-01 0.35761356E-02 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

(TOTAL SEC) \*\*\* (LGI=1...25)  
 0. 0. 0.57653844E-04 0.29017699E 00 0.58414833E 00 0.28631179E 00  
 0.58651783E-01 0.35761356E-02 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

\*\*\* FLUX \*\*\*

0. 0. 8.4291242 E-01 4.7372574 E-01 5.8141658 E-01 6.8869995 E-01  
 6.9701371 E-01 6.9668743 E-01 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

U 238  
 U 238 N INELAS 3 N CROSS C EI BN A/W 47 0 63 1.000E-03 1.500E U7EV A0047 0  
 42 1 92 238 1 4 3 1 1 1 2 A0047 1  
 1 2 1 0 0 0 128 47 U 63 0 A0047 2  
 2.3807E 02 -3.0000E 05 0. 1.0000E 00 1.0000E-03 1.5000E 07 A0047 3  
 ALDERMASTON/WINFRITH DATA FILE: 7/1/64 NIN= 5 GCN= 1 PCN= 7 A0047 4  
 AEEW R351: 2/64, BARRINGTON ETAL. AWRE 0-79/63, 1/64, K-PARKER A0047 5  
 DCC1 401 1 5 0 102 0 0-0. 0. A0047 6  
 0.31000000E 06 0.10000000E-02 0.20000000E 07 0.10000000E-02 (FIRST AND LAST VALUE)  
 NDYTP=1

\*\*\* U 238 (N-INFLAS) 3,N CROSS C(EI, , ) \*\*\*

SEC= 1 (LGI=1...25)  
 0. 0. 0.57476375E-04 0.40409434E-01 0.90498418E-01 0.53319902E-01  
 0.34631272E-02 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

(TOTAL SEC) \*\*\* (LGI=1...25)  
 0. 0. 0.57476375E-04 0.40409434E-01 0.90498418E-01 0.53319902E-01  
 0.34631272E-02 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

\*\*\* FLUX \*\*\*

0. 0. 8.4291242 E-01 4.7372641 E-01 5.8141658 E-01 6.8869995 E-01  
 6.9674813 E-01 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.  
 0. 0. 0. 0. 0. 0.

U 238  
 U 238 N INELAS 4 N CROSS C EI BN A/W 50 0 63 1.000E-03 1.500E U7EV A0050 0  
 36 1 92 238 1 4 4 1 1 1 2 A0050 1  
 1 2 1 0 0 0 128 50 0 63 0 A0050 2  
 2.3807E 02 -7.0000E 05 0. 1.0000E 00 1.0000E-03 1.5000E 07 A0050 3  
 ALDERMASTON/WINFRITH DATA FILE: 7/1/64 NIN= 5 GCN= 1 PCN= 8 A0050 4  
 AEEW R351: 2/64, BARRINGTON ETAL. AWRE 0-79/63, 1/64, K-PARKER A0050 5  
 DCC1 401 1 5 0 82 0 0-0. 0. A0050 6  
 0.71000000E 06 0.10000000E-02 0.20000000E 07 0.10000000E-02 (FIRST AND LAST VALUE)  
 NDYTP=1

\*\*\* U 238 (N-INFLAS) 4-N CRUS C(EI, , ) \*\*\*

```

SEC= 1 (LGI=1...25)
0.      0.      0.57704802E-04  0.17959515E 00  0.46518985E 00  0.10958495E-01
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

```

(TOTAL SEC) *** (LGI=1...25)
0.      0.      0.57704802E-04  0.17959515E 00  0.46518985E 00  0.10958495E-01
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

\*\*\* FLUX \*\*\*

```

0.      0.      8.4291242 E-01  4.7372597 E-01  5.8141658 E-01  6.8875705 E-01
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

```

U 238
U 238 N INELAS 5 N CRUS C EI      BN A/W 53 0 63 1.000E-03 1.500E 07EV A0053 0
23 1 92 238 1 4 5 1 1 1 2 A0053 1
1 2 1 0 0 0 0 128 53 0 63 0 A0053 2
2.3807E 02 -9.8000E 05 0. 1.0000E 00 1.0000E-03 1.5000E 07 A0053 3
ALDERMASTON/WINFRITH DATA FILE, 7/1/64 NIN= 5 GCN= 1 PCN= 9 A0053 4
AEEW R351, 2/64, BARRINGTON ETAL, AWRE 0-79/63, 1/64, K-PARKER A0053 5
DCC1 401 1 5 0 45 0 0 -0. 0. A0053 6
0.99000000E 06 0.10000000E-02 0.20100000E 07 0.10000000E-02 (FIRST AND LAST VALUE)
NDTYP=1

```

\*\*\* U 238 (N-INFLAS) 5-N CRUS C(EI, , ) \*\*\*

```

SEC= 1 (LGI=1...25)
0.      0.      0.10664659E-02  0.37472566E 00  0.34777647E 00  0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

```

(TOTAL SEC) *** (LGI=1...25)
0.      0.      0.10664659E-02  0.37472566E 00  0.34777647E 00  0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

\*\*\* FLUX \*\*\*

```

0.      0.      8.4290313 E-01  4.7372717 E-01  5.8141728 E-01  0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

```

U 238
U 238 N INELAS 6 N CRUS C EI      BN A/W 56 0 63 1.000E-03 1.500E 07EV A0056 0
22 1 92 238 1 4 6 1 1 1 2 A0056 1
1 2 1 0 0 0 0 128 56 0 63 0 A0056 2
2.3807E 02 -1.0600E 06 0. 1.0000E 00 1.0000E-03 1.5000E 07 A0056 3
ALDERMASTON/WINFRITH DATA FILE, 7/1/64 NIN= 5 GCN= 1 PCN= 10 A0056 4
AEEW R351, 2/64, BARRINGTON ETAL, AWRE 0-79/63, 1/64, K-PARKER A0056 5
DCC1 401 1 5 0 41 0 0 -0. 0. A0056 6
0.10700000E 07 0.10000000E-02 0.20100000E 07 0.10000000E-02 (FIRST AND LAST VALUE)
NDTYP=1

```

\*\*\* U 238 (N-INFLAS) 6-N CRUS C(EI, , ) \*\*\*

```

SEC= 1 (LGI=1...25)
0.      0.      0.14380618E-02  0.41957524E 00  0.19576222E 00  0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

```

(TOTAL SEC) *** (LGI=1...25)
0.      0.      0.14380618E-02  0.41957524E 00  0.19576222E 00  0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

\*\*\* FLUX \*\*\*

```

0.      0.      8.4290313 E-01  4.7372895 E-01  5.8141506 E-01  0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.
0.      0.      0.      0.      0.      0.

```

```

U 238
U 238 N INELAS 98 N CRUS C EI      BN A/W 59 0 63 1.000E-03 1.500E 07EV A0059 0
32 1 92 238 1 4 98 1 1 1 2 A0059 1
1 2 1 0 0 0 0 128 59 0 63 0 A0059 2
2.3807E 02 -4.4000E 05 0. 1.0000E 00 1.0000E-03 1.5000E 07 A0059 3
ALDERMASTON/WINFRITH DATA FILE, 7/1/64 NIN= 5 GCN= 1 PCN= 15 A0059 4
AEEW R351, 2/64, BARRINGTON ETAL, AWRE 0-79/63, 1/64, K-PARKER A0059 5
DCC1 401 1 5 0 72 0 0 -0. 0. A0059 6
0.12700000E 07 0.10000000E-02 0.15000000E 08 0.14700000E 00 (FIRST AND LAST VALUE)
NDTYP=1

```



```

FE 0 N TOTAL 0 CRUS C(EI) BN UK- 36 0 0 1.000E-04 1.500E 07EV E0291 0
270 0 1 20 0 1 1 0 0 1 1 1 2 E0291 1
1 0 1 0 0 0 0 0 0 0 0 0 0 E0291 2
5.5045E 01 0.0 0.0 1.0000E 00 1.0000E-04 1.5000E 07 E0291 3
ENR 200 1 2 0 0 0 0 0 1.0000E-04 2.7820E 06 E0291 4
DCCI 401 1 5 0 690 0 0 0.0 0.0 E0291 5
9.9999902E-05 4.00000000E 01 2.78200000E 06 3.18399906E 00 (FIRST AND LAST VALUE)
0.99999990E-04 0.14000000E 03 0.27820000E 07 0.10318399E 03 (FIRST AND LAST VALUE)=SIG.TOT( 1, 1379)
ENR 200 2 2 0 0 0 0 0 2.7820E 06 1.5000E 07 E0291236
DCCI 401 1 5 0 92 0 0 0.0 0.0 E0291237
2.7820000E 06 3.18399906E 00 1.5000000E 07 3.20349979E 00 (FIRST AND LAST VALUE)
0.2782000E 07 0.10318399E 03 0.1500000E 08 0.10320349E 03 (FIRST AND LAST VALUE)=SIG.TOT( 1381, 1563)

```

```

FE 0 N ELAST 0 N CRUS C(EI) BN UK- 36 0 0 1.000E-04 1.500E 07EV E0292 0
270 1 1 20 0 1 2 0 0 1 1 1 2 E0292 1
1 0 1 0 0 0 0 0 0 0 0 0 0 E0292 2
5.5045E 01 0.0 0.0 1.0000E 00 1.0000E-04 1.5000E 07 E0292 3
ENR 200 1 2 0 0 0 0 0 1.0000E-04 2.7820E 06 E0292 4
DCCI 401 1 5 0 690 0 0 0.0 0.0 E0292 5
9.9999902E-05 3.79999924E 00 2.78200000E 06 2.28869915E 00 (FIRST AND LAST VALUE)
ENR 200 2 2 0 0 0 0 0 2.7820E 06 1.5000E 07 E0292236
DCCI 401 1 5 0 92 0 0 0.0 0.0 E0292237
2.7820000E 06 2.28869915E 00 1.5000000E 07 1.12999916E 00 (FIRST AND LAST VALUE)

```

\*\*\* FE 0 (N,TOTAL) 0, CRUS C(EI, , ) \*\*\*

SEC= 1 (LGI=1...70)

0.2149350E 01	0.35306244E 01	0.36883392E 01	0.36867765E 01	0.34729919E 01	0.33255157E 01
0.31506151E 01	0.28081335E 01	0.25450439E 01	0.24707031E 01	0.30222321E 01	0.26510315E 01
0.37592775E 01	0.32530975E 01	0.24235077E 01	0.31864929E 01	0.37312164E 01	0.41543427E 01
0.29002380E 01	0.63298035E 01	0.40079041E 01	0.43726196E 01	0.54647773E 01	0.18713501E 02
0.51777649E 01	0.2221832E 01	0.3747051E 01	0.57673187E 01	0.11027100E 02	0.18883423E 02
0.62766665E 01	0.54957581E 01	0.52636414E 01	0.59508362E 01	0.68818817E 01	0.77271271E 01
0.9212310E 01	0.93731232E 01	0.93598863E 01	0.10396744E 02	0.10735397E 02	0.10999878E 02
0.1187454E 02	0.11314346E 02	0.11393158E 02	0.11428055E 02	0.11443283E 02	0.11454437E 02
0.11472397E 02	0.11475403E 02	0.11479202E 02	0.11484665E 02	0.11493881E 02	0.11505783E 02
0.11519958E 02	0.11536255E 02	0.11554703E 02	0.11575851E 02	0.11599274E 02	0.11626755E 02
0.11659348E 02	0.11694195E 02	0.11734177E 02	0.11779770E 02	0.11829559E 02	0.11888702E 02
0.11955063E 02	0.12030380E 02	0.12117645E 02	0.12215927E 02		

(TOTAL SEC) \*\*\* (LGI=1...70)

0.2149350E 01	0.35306244E 01	0.36883392E 01	0.36867765E 01	0.34729919E 01	0.33255157E 01
0.31506151E 01	0.28081335E 01	0.25450439E 01	0.24707031E 01	0.30222321E 01	0.26510315E 01
0.37592775E 01	0.32530975E 01	0.24235077E 01	0.31864929E 01	0.37312164E 01	0.41543427E 01
0.29002380E 01	0.63298035E 01	0.40079041E 01	0.43726196E 01	0.54647773E 01	0.18713501E 02
0.51777649E 01	0.2221832E 01	0.3747051E 01	0.57673187E 01	0.11027100E 02	0.18883423E 02
0.62766665E 01	0.54957581E 01	0.52636414E 01	0.59508362E 01	0.68818817E 01	0.77271271E 01
0.9212310E 01	0.93731232E 01	0.93598863E 01	0.10396744E 02	0.10735397E 02	0.10999878E 02
0.1187454E 02	0.11314346E 02	0.11393158E 02	0.11428055E 02	0.11443283E 02	0.11454437E 02
0.11472397E 02	0.11475403E 02	0.11479202E 02	0.11484665E 02	0.11493881E 02	0.11505783E 02
0.11519958E 02	0.11536255E 02	0.11554703E 02	0.11575851E 02	0.11599274E 02	0.11626755E 02
0.11659348E 02	0.11694195E 02	0.11734177E 02	0.11779770E 02	0.11829559E 02	0.11888702E 02
0.11955063E 02	0.12030380E 02	0.12117645E 02	0.12215927E 02		

\*\*\* FLUX \*\*\*

3.65900923E-05	1.29267806E-04	3.09529947E-04	5.64257847E-04	8.73690937E-04	9.06016212E-04
1.25365795E-03	1.35200541E-03	9.43141524E-04	1.01611624E-03	2.31877342E-03	2.25147605E-03
2.15048343E-03	2.46798061E-03	2.09974684E-03	2.16220692E-03	2.77131237E-03	2.13958882E-03
1.77176297E-03	2.41457298E-03	2.46388093E-03	2.41007656E-03	2.42719054E-03	2.12415308E-03
2.41532177E-03	2.53046118E-03	2.43100710E-03	2.40818970E-03	2.31680833E-03	2.15365924E-03
2.32165121E-03	2.42597237E-03	2.45573372E-03	2.42605060E-03	2.42027827E-03	2.34123319E-03
2.35139345E-03	2.35435553E-03	2.33463128E-03	2.27930211E-03	2.31192634E-03	2.32901052E-03
2.3117981E-03	2.32436694E-03	2.26492435E-03	2.28766166E-03	2.31062993E-03	2.30353139E-03
2.25750799E-03	2.29756720E-03	2.32015178E-03	2.30676495E-03	2.32062116E-03	2.26319768E-03
2.28954852E-03	2.31463020E-03	2.30282173E-03	2.25682557E-03	2.29502097E-03	2.31801718E-03
2.30702479E-03	2.31823884E-03	2.25858204E-03	2.28422135E-03	2.30344571E-03	2.29584426E-03
2.24917941E-03	2.28618830E-03	2.30786577E-03	2.29178555E-03		

\*\*\* FE 0 (N,ELAST) 0,N CRUS C(EI, , ) \*\*\*

SEC= 1 (LGI=1...70)

0.18152962E 01	0.21443005E 01	0.23297539E 01	0.22693338E 01	0.22813969E 01	0.23630152E 01
0.22970629E 01	0.22302628E 01	0.21284046E 01	0.22710209E 01	0.30244150E 01	0.26465282E 01
0.37562799E 01	0.32649155E 01	0.24301033E 01	0.31890669E 01	0.37950144E 01	0.42793713E 01
0.29006672E 01	0.66294899E 01	0.41735220E 01	0.43718300E 01	0.54455833E 01	0.21678406E 02
0.63901510E 01	0.22163916E 01	0.37416410E 01	0.57617216E 01	0.11123222E 02	0.19170624E 02
0.83416834E 01	0.54925241E 01	0.52554979E 01	0.59390068E 01	0.68537159E 01	0.76412354E 01
0.84268742E 01	0.93192663E 01	0.99409733E 01	0.10380540E 02	0.10710527E 02	0.10968756E 02
0.11160350E 02	0.11285377E 02	0.11355880E 02	0.11388144E 02	0.11399502E 02	0.11400007E 02
0.11400002E 02	0.1139997E 02	0.11400002E 02	0.11399989E 02	0.11399994E 02	0.11399997E 02
0.11399991E 02	0.11399989E 02	0.11399994E 02	0.11400002E 02	0.11399995E 02	0.11399993E 02
0.11399995E 02	0.11400001E 02	0.11399994E 02	0.11399998E 02	0.11399992E 02	0.11399995E 02
0.11399989E 02	0.11399998E 02	0.11399996E 02	0.11400000E 02		

(TOTAL SEC) \*\*\* (LGI=1...70)

0.18152962E 01	0.21443005E 01	0.23297539E 01	0.22693338E 01	0.22813969E 01	0.23630152E 01
0.22970629E 01	0.22302628E 01	0.21284046E 01	0.22710209E 01	0.30244150E 01	0.26465282E 01
0.37562799E 01	0.32649155E 01	0.24301033E 01	0.31890669E 01	0.37950144E 01	0.42793713E 01
0.29006672E 01	0.66294899E 01	0.41735220E 01	0.43718300E 01	0.54455833E 01	0.21678406E 02
0.63901510E 01	0.22163916E 01	0.37416410E 01	0.57617216E 01	0.11123222E 02	0.19170624E 02
0.83416834E 01	0.54925241E 01	0.52554979E 01	0.59390068E 01	0.68537159E 01	0.76412354E 01
0.84268742E 01	0.93192663E 01	0.99409733E 01	0.10380540E 02	0.10710527E 02	0.10968756E 02
0.11160350E 02	0.11285377E 02	0.11355880E 02	0.11388144E 02	0.11399502E 02	0.11400007E 02
0.11400002E 02	0.1139997E 02	0.11400002E 02	0.11399989E 02	0.11399994E 02	0.11399997E 02
0.11399991E 02	0.11399989E 02	0.11399994E 02	0.11400002E 02	0.11399995E 02	0.11399993E 02
0.11399995E 02	0.11400001E 02	0.11399994E 02	0.11399998E 02	0.11399992E 02	0.11399995E 02
0.11399989E 02	0.11399998E 02	0.11399996E 02	0.11400000E 02		





\*\*\* B 0 (N,ELAST) 0,N CROSS (C, E I, , ) \*\*\*

SEC= 1 (LGI=1...69)

Table with 6 columns of numerical data representing cross-sections for various energy groups and materials.

(TOTAL SEC) \*\*\* (LGI=1...69)

Table with 6 columns of numerical data, identical to the one above, representing total cross-sections.

\*\*\* FLUX \*\*\*

Table with 6 columns of numerical data representing flux values for different energy groups.

Table with 6 columns of numerical data, including labels like DCC2, DCC1, and various energy group identifiers.

\*\* ARRANGED DATA \*\*

I= 1 EI= 0.1000000E-03 DATA(X,Y)

Table with 6 columns of numerical data representing arranged data for I=1.

\*\* TOTAL P= 0.9999999E 00

\*\* RE-NORMALIZED DATA \*\*

I= 1 EI= 0.1000000E-03 DATA(X,Y)

Table with 6 columns of numerical data representing re-normalized data for I=1.

\*\* TOTAL P= 0.9999999E 00

\*\* ARRANGED DATA \*\*

I= 2 EI= 0.1000000E 01 DATA(X,Y)

Table with 6 columns of numerical data representing arranged data for I=2.

B 0 N ELAST 0 N ANG D P CC EI UK- 15 -0 -0 1.000E-04 6.000E 06EV E0393 28

NENR= 1(S2DP)

\*LGF= 69 E( 2.78000E-01 ... 3.60000E-01) LGI= 4...70

SIG.

Table with 6 columns of numerical data representing sigma values for various energy groups.



\*\*\*SUMMATION OF ONE DIMENSIONAL DATA

Table with 6 columns of numerical data in scientific notation, representing summation of one-dimensional data.

\*\*\*SUMMATION OF TRANSFER MATRICES

Table with 6 columns of numerical data in scientific notation, representing summation of transfer matrices.

MU(I)

Table with 6 columns of numerical data in scientific notation, representing MU(I) values.

SIG(I+J)

J = 1

Table with 6 columns of numerical data in scientific notation, representing SIG(I+J) for J=1.

J = 2

Table with 6 columns of numerical data in scientific notation, representing SIG(I+J) for J=2.

J = 3

Table with 6 columns of numerical data in scientific notation, representing SIG(I+J) for J=3.