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MEDUSA-PIJ : A CODE FOR ONE-DIMENSIONAL  
LASER FUSION ANALYSIS TAKING ACCOUNT  
OF NEUTRON HEATING EFFECT

March 1979

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MEDUSA-PIJ : A Code for One-Dimensional  
Laser Fusion Analysis Taking Account  
of Neutron Heating Effect

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In the one-dimensional Lagrangian hydro-burn code MEDUSA, the neutrons emitted by D-T and D-D reactions are assumed to escape freely from the system without further interaction. The assumption is not practicable, however, in plasmas highly compressed by laser irradiation.

In the modified code MEDUSA-PIJ, the energy deposition due to the emitted neutrons is calculated by collision probability method. The total energy deposition is considered in the energy balance equation of the system at each time step, and a hydrodynamic equation is solved. Neutron spectra escaping from fuel sphere are calculated by a neutron slowing-down equation derived under the two-collision model.

The MEDUSA-PIJ code is designed for FACOM-230/75 computer with a FORTRAN IV H-level compiler.

Keywords : Laser Fusion, Neutron Heating, Collision Probability, MEDUSA Code,  
Neutron Slowing Down, Energy Deposition

中性子加熱を考慮した1次元レーザー核融合  
解析コード (MEDUSA-PIJ)

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1次元のラグランジアン熱核融合燃焼コードMEDUSAにおいては、D-TとD-D反応によって生れた中性子は、衝突なしに系から逃れると仮定されている。この仮定はレーザー照射によって高密度に圧縮されたプラズマ状態では実際的ではない。

新しく改良したMEDUSA-PIJコードでは、この放出中性子によるエネルギーデポジションを衝突確率法を用いて計算することができる。このエネルギーデポジションは各タイムステップ毎に系のエネルギーバランス方程式に考慮され、流体方程式が解かれる。燃料球から逃れる中性子のスペクトラムが2回衝突モデルから導出した中性子の減速方程式を用いて計算される。

MEDUSA-PIJコードはFACOM-30/75のFORTRAN-IV-Hコンパイラーを使用するようになっている。

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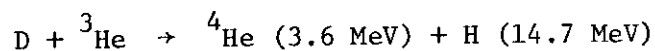
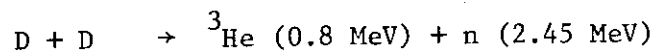
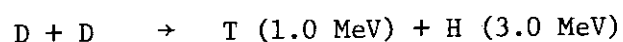
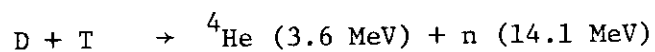
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## 1. Introduction

Theoretical studies for inertially confined fusion reactor indicate that symmetric illumination of a hydrogen fuel pellet will lead to high-pellet-compression ratios of  $\sim 10^4$  times liquid density with an associated implosion heating process. A description based on the classical theory can be given for this process as follows: Under laser irradiation, electrons near the critical surface of the pellet are heated by inverse bremsstrahlung, resonance absorption and various anomalous mechanisms.<sup>(1)</sup> This heat is transported inward by classical thermal conductivity, and transferred to the ions by classical electron-ion collisions. The pellet surface is ablated by the electrons heated and the recoil force due to the expansion of ablating ions compresses the fuel sphere to high dense plasma. The thermal wavefront with the ablation surface acts the role like a piston. Once ignition of the central region of the imploded pellet occurs, a thermonuclear burn front propagates radially outward and ignites additional fuel. The MEDUSA code<sup>(2)</sup> can calculate one-dimensional hydrodynamic and thermodynamic behaviour such as the implosion and burning process of plasma described above.

In the MEDUSA code, the 4 thermonuclear reactions are considered as follows:



Here, the charged reaction products H, T,  ${}^3\text{He}$  and  ${}^4\text{He}$  are assumed to deposit their energy locally to the ions and electrons. The neutrons are

assumed to escape freely from the system without further interaction. The assumption made for the neutron escape is not reliable in the state of high dense plasma compressed. This can be easily understood from the following simple reasoning : The neutron escape probability, which one 14.1 MeV neutron emitted uniformly at the center of a sphere with uniform density escapes without collision, is  $e^{-\Sigma x}$ , where  $\Sigma$  is the macroscopic cross section of the sphere. In the initial state of D-T fuel, the distance  $x$  is  $\sim 10^{-1}$  cm and the density  $\rho_0$  is  $\sim 0.1$  g/cm<sup>3</sup>, and in a over-dense plasma,  $x$  is  $\sim 10^{-3}$  cm and  $\rho$  is  $\sim 1000$  g/cm<sup>3</sup>; Thus, the neutron escape probabilities are 96 % for the initial state and 67 % for the over-dense state, respectively. In the over-dense state, therefore, 30 % of the neutron energy is deposited before the neutron escapes, and the energy deposited becomes larger than that of  $\alpha$ -particle. Hence, the effect of neutron heating by neutron-ion interaction should not be neglected in laser fusion analysis.

The MEDUSA-PIJ code can calculate the effect of the neutron heating and the spectrum of neutron escaping from the fuel sphere to blanket region by using the collision probability method.<sup>(3)</sup> In the Chapter 2, the calculational methods for the neutron spectrum and energy deposition are described. The MEDUSA-PIJ usage is described in the Chapter 3.



2. Computational Method of Neutron Heating Effect and Neutron Spectrum

In the newly modified MEDUSA-PIJ code, the energy deposition by neutrons is calculated with the use of the collision probability method. (3) The D-T fuel sphere is divided into M shells which are expressed at the points of a moving Lagrangian mesh as shown in Fig.1. All physical quantities are calculated at the Lagrangian mesh points at each time step. Then, the collision probability  $P_{IJ}$  is defined as the probability that a neutron emitted isotropically in a shell I has its next collision in a shell J. The probability  $P_{IJ}$  is given by

$$P_{IJ} = \frac{1}{4\pi V_I} \int_{V_I} dx_I \int_{4\pi} d\Omega \int dR \Sigma_J e^{-\tilde{\Sigma}R}$$

where

$V_I$  = volume of sphere I

$R = |x_I - x_J|$  = distance between point  $x_I$  and  $x_J$  in the shell I and J, respectively

$\Omega = (x_I - x_J)/R$

$\tilde{\Sigma}R = \int_0^R \Sigma(x_I - \Omega R') dR' =$  optical distance between  $x_I$  and  $x_J$ .

The neutron slowing down and energy deposition equations for the first collision are given by

$$Q_I^{(1)}(n) = \sum_J N_J P_{JI}^{(1)} \sum_A \frac{\Sigma_{A,I}^{(1)}}{\Sigma_I^{(1)}} T_A(1,n) \quad (2)$$

and

$$\langle DE(n) \rangle_I^{(1)} = \sum_J N_J P_{JI}^{(1)} \sum_A \frac{\Sigma_{A,I}^{(1)}}{\Sigma_I^{(1)}} T_A(1,n) E_A(1,n) \quad , \quad (3)$$

where  $N_J$  is the neutrons emitted isotropically at the J-th shell region,

$P_{JI}(1)$  the collision probability that a neutron of the energy group 1 emitted at the region J will have the next collision at the region I,  $\Sigma_{A,I}(1)$  the macroscopic neutron elastic scattering cross section of material A for the energy group 1,  $\Sigma_I(1)$  the total cross section summed over material A,  $T_A(1,n)$  and  $E_A(1,n)$  mean the energy transfer and energy loss matrices from group 1 to n, respectively.  $Q_I^{(1)}(n)$  expresses the number of neutrons slowing down to group n by the first collision and  $\langle DE(n) \rangle_I^{(1)}$  the energy deposition by the slowing down of neutrons in the region I. The energy group 1 is 14.1 MeV and the other group boundaries are determined by dividing the energy space with equal lethargy width  $\Delta u=0.22$ . The energy group structure and elastic scattering cross sections of D, T, H,  $^3\text{He}$  and  $^4\text{He}$  are shown in Table 1.

In Eq.(2), the energy transfer matrix,  $T_A(m,n)$  is defined as follows:

①  $m=1$

$$T_A(1,n) = \begin{cases} 0.0 & \dots\dots n=1 \\ \Delta E_n / (1-\alpha_A) E_0 & \dots\dots n \neq 1, L_A \\ (E_{L_A-1} - \alpha_A E_0) / (1-\alpha_A) E_0 & \dots\dots n=L_A \\ 0.0 & \dots\dots n > L_A \end{cases} \quad (4)$$

②  $m \neq 1, n=m$

$$T_A(m,n) = \frac{\int_{E_m}^{E_{m-1}} \int_{E_m}^{E'} \frac{1}{(1-\alpha_A) E'} dE dE'}{\int_{E_m}^{E_{m-1}} dE'} = (\Delta E_m - E_u \Delta u) / (1-\alpha_A) \Delta E_m \quad (5)$$

③  $m \neq 1, m < n < m+L_A-1$

$$T_A(m,n) = \frac{\int_{E_m}^{E_{m-1}} \int_{E_n}^{E_{n-1}} \frac{1}{(1-\alpha_A) E'} dE dE'}{\int_{E_m}^{E_{m-1}} dE'} = \Delta E_n \Delta u / (1-\alpha_A) \Delta E_m \quad (6)$$

④  $m \neq 1, n = m + L_A - 1$

$$T_A(m, n) = \left\{ \int_{E_n/\alpha_A}^{E_{m-1}} \int_{\alpha_A E'}^{E_{n-1}} \frac{1}{(1-\alpha_A)E'} dE dE' + \int_{E_m}^{E_n/\alpha_A} \int_{E_m}^{E_{n-1}} \frac{1}{(1-\alpha_A)E'} dE dE' \right\} / \int_{E_m}^{E_{m-1}} dE'$$

$$= \frac{1}{(1-\alpha_A)\Delta E_m} \{ \Delta E_n \Delta u - (\alpha_A E_{m-1} - E_n) + E_n (L_A \Delta u - \xi_A) \} \quad (7)$$

⑤  $m \neq 1, n = m + L_A$

$$T_A(m, n) = \int_{E_m}^{E_{n-1}/\alpha_A} \int_{\alpha_A E'}^{E_{n-1}} \frac{1}{(1-\alpha_A)E'} dE dE' / \int_{E_m}^{E_{m-1}} dE'$$

$$= \frac{1}{(1-\alpha_A)\Delta E_m} [E_{n-1}(\xi_A - (L_A - 1)\Delta u) - (E_{n-1} - \alpha_A E_m)] \quad (8)$$

where  $L_A$  is the maximum group number of slowing down and  $\alpha_A = (M-1)^2 / (M+1)^2$  is a well-known parameter concerning the atomic mass  $M$  of material  $A$ .

In Eq.(3), the energy loss matrix  $E_A(m, n)$  is defined as follows:

①  $m=1$

$$E_A(1, n) = \begin{cases} 0 & \dots n=1 \\ E_0 - \bar{E}_n & \dots n \neq 1, L_A \\ E_0 - \frac{1}{2}(E_{L_A-1} + \alpha_A E_0) & \dots n=L_A \\ 0 & \dots n > L_A \end{cases} \quad (9)$$

②  $m \neq 1, n=m$

$$E_A(m, n) = \int_{E_m}^{E_{m-1}} \int_{E_n}^{E'} \frac{E' - E}{(1-\alpha_A)E'} dE dE' / \int_{E_m}^{E_{m-1}} \int_{E_n}^{E'} \frac{1}{(1-\alpha_A)E'} dE dE'$$

$$= \frac{1}{2} \{ \Delta E_m (E_n - 2E_n) + E_n^2 \Delta u \} / T_A(m, n) (1-\alpha_A) \Delta E_m \quad (10)$$

③  $m \neq 1, m < n < m+L_A-1$

$$E_A(m,n) = \int_{E_m}^{E_{m-1}} \int_{E_n}^{E_{n-1}} \frac{E'-E}{(1-\alpha_A)E'} dEdE' / \int_{E_m}^{E_{m-1}} \int_{E_n}^{E_{n-1}} \frac{1}{(1-\alpha_A)E'} dEdE'$$

$$= \Delta E_m (\Delta E_m - \bar{E}_n \Delta u) / T_A(m,n) \Delta E_m (1-\alpha_A) \quad (11)$$

④  $m \neq 1, n = m+L_A-1$

$$E_A(m,n) = \left\{ \int_{E_n/\alpha_A}^{E_{m-1}} \int_{\alpha_A E'}^{E_{n-1}} \frac{E'-E}{(1-\alpha_A)E'} dEdE' + \int_{E_m}^{E_n/\alpha_A} \int_{E_n}^{E_{n-1}} \frac{E'-E}{(1-\alpha_A)E'} dEdE' \right\} /$$

$$\left\{ \int_{E_n/\alpha_A}^{E_{m-1}} \int_{\alpha_A E'}^{E_{n-1}} \frac{1}{(1-\alpha_A)E'} dEdE' + \int_{E_m}^{E_n/\alpha_A} \int_{E_n}^{E_{n-1}} \frac{1}{(1-\alpha_A)E'} dEdE' \right\}$$

$$= \left\{ \frac{1}{2} [(E_{m-1} - E_n/\alpha_A) (2E_{n-1} - \alpha_A (1 - \frac{\alpha_A}{2}) (E_{m-1} + E_n/\alpha_A)) - E_{n-1}^2 (L_A \Delta u - \xi_A)] + \Delta E_n [(E_n/\alpha_A - E_m) - \bar{E}_n (\xi_A - (L_A - 1) \Delta u)] \right\} /$$

$$T_A(m,n) (1-\alpha_A) \Delta E_m \quad (12)$$

⑤  $m \neq 1, n = m+L_A$

$$E_A(m,n) = \int_{E_m}^{E_{n-1}/\alpha_A} \int_{\alpha_A E'}^{E_{n-1}} \frac{E'-E}{(1-\alpha_A)E'} dEdE' / \int_{E_m}^{E_{n-1}/\alpha_A} \int_{\alpha_A E'}^{E_{n-1}} \frac{1}{(1-\alpha_A)E'} dEdE'$$

$$= \frac{1}{2} [(E_{n-1}/\alpha_A - E_m) (2E_{n-1} - \alpha_A (1 - \frac{\alpha_A}{2}) (E_{n-1}/\alpha_A + E_n)) - E_{n-1}^2 (\xi_A - (L_A - 1) \Delta u)] / T_A(m,n) (1-\alpha_A) \Delta E_m \quad (13)$$

By considering the second collision of neutrons, the slowing down and energy deposition equations are respectively expressed by

$$Q_I^{(2)}(n) = \sum_J \sum_m Q_J^{(1)}(m) P_{JI}(m) \sum_A \frac{\Sigma_{A,I}(m)}{\Sigma_I(m)} T_A(m,n), \quad (14)$$

and

$$\langle DE(n) \rangle_I^{(2)} = \sum_J \sum_m Q_J^{(1)}(m) P_{JI}(m) \sum_A \frac{\Sigma_{A,I}(m)}{\Sigma_I(m)} T_{A(m,n)} E_{A(m,n)} . \quad (15)$$

Thus the total energy deposition by the neutrons with considering two collisions is given by,  $\langle DE \rangle_I = \sum_n [\langle DE(n) \rangle_I^{(1)} + \langle DE(n) \rangle_I^{(2)}]$ . The spectrum of neutrons escaping from the spherical shell to blanket region is given by

$$S(n) = \begin{cases} \sum_J N_J P_J(n) , & n = 1 \end{cases} \quad (16)$$

$$\begin{cases} \sum_J (Q_J^{(1)}(n) + Q_J^{(2)}(n)) P(n) , & 1 < n \end{cases} \quad (17)$$

where  $P_J(n) = 1 - \sum_I P_{JI}(n)$  is the collision escape probability.

In the MEDUSA-PIJ code, the total energy deposition of the emitted neutrons is considered in the energy balance equation of the system at each time step, and a hydrodynamic equation is solved. The internal energy per unit mass,  $U$ , is assumed to be separable into electric and ionic components,  $U_e$  and  $U_i$ , which are directly coupled only through electron-ion collisional energy exchange. The electron and ion pressures are assumed only to do  $PdV$  work on the electrons and ions, respectively. The internal energy balance equations for electrons and ions are written as

$$\frac{dU_e}{dt} = - P_e \frac{dv}{dt} + H_e + K + Y_e + J + X \quad (18)$$

$$\frac{dU_i}{dt} = - P_i \frac{dv}{dt} + H_i - K + Y_i + Q , \quad (19)$$

where  $H$  is the flow heat due to thermal conduction,  $K$  the rate of energy exchange between electrons and ions,  $Y$  the rate of thermonuclear energy release,  $J$  the rate of bremsstrahlung emission,  $X$  the absorption rate of laser light and  $Q$  the rate of viscous shock heating. The energy deposition by the neutron is added to the thermonuclear energy rate  $Y_i$  in the MEDUSA-PIJ code.

3. Input Data Description

In the MEDUSA-PIJ code, the input of data are mainly made with "namelist" function. The input variables are defined by "default" values. Therefore, only the variables to be changed in the calculation need to be read with the use of "namelist" function.

3.1 Input Data

# 1 (12A4) Tittle is described by four cards

LABEL 1(12), LABEL 2(12), LABEL 3(12), LABEL 4(12)

# 2 Namelist NEWRUN

Variable	Descriptions	Default value
AK 0	$a_0$ , Control variable of time centering in Eq.(A-39)	5.0
AK 1	$a_1$ , Control variable of time step $\Delta t$ by sound speed in Eq.(A-35)	0.25
AK 2	$a_2$ , Control variable of time step $\Delta t$ by density variation in Eq.(A-36)	0.25
AK 3	$a_3$ , Control variable of time step $\Delta t$ by ion temperature variation in Eq.(A-37)	0.25
AK 4	$a_4$ , Control variable of time step $\Delta t$ by electron temperature variation in Eq.(A-38)	0.25
BNEUM	b, Factor for the rate of shock heating in Eq.(A-33)	1.0
DELTAT	Time step interval $\Delta t_0$ , $\Delta t \leq \Delta t_0$ (sec)	$1 \times 10^{-12}$
DTEMAX	$\delta T_e$ , Convergence condition for electron temperature in Eq.(A-42)	0.1
DTIMAX	$\delta T_i$ , Convergence condition for ion temperature in Eq.(A-41)	0.1
DUMAX	$\delta u$ , Convergence condition for plasma velocity in Eq.(A-40)	0.1
DEUTER	$f_D$ , Initial fraction of D	0.5

TRITIU	$f_T$ , Initial fraction of T	0.5
HELIU 3	$f_{3He}$ , Initial fraction of $^3He$	0.0
NETRAL	$f_N$ , Initial fraction of neutrals	0.0
HELIU 4	$f_{4He}$ , Initial fraction of $^4He$	0.0
XTRA	$f_X$ , Initial fraction of arbitrary ion X	0.0
HYDROG	$f_H$ , Initial fraction of H	0.0
GAMMAE	$\gamma_e$	$\frac{5}{3}$
GAMMAI	$\gamma_i$	$\frac{5}{3}$
LAMDA 1	$\lambda$ , Wavelength of laser light (m)	$1 \times 10^{-5}$
NTRLMS	$M_N$ , Atomic mass number of neutrals (kg)	0.0
RHOINI	$\rho_0$ , Initial density of fuel ( $kg/m^3$ )	124.0
RINI	$R_0$ , Initial size of system (m)	$4.8 \times 10^{-4}$
TEINI	$T_e$ , Initial electron temperature ( $^{\circ}K$ )	500
TIINI	$T_i$ , Initial ion temperature ( $^{\circ}K$ )	500
TINUCL	Ignition temperature for nuclear fusion reaction ( $^{\circ}K$ )	$1 \times 10^7$
TSTOP	Maximum calculational time (sec)	$1 \times 10^{-6}$
XMASS	$M_X$ , Atomic mass number of arbitrary ion X (kg)	0.0
XZ	$Z_X$ , Charge number of X	0.0
SCP	Scaling factor of pressure for output format	1.0
SCR	Scaling factor of coordinate r for output format	1.0
SCRHO	Scaling factor of density $\rho$ for output format	1.0
SCTE	Scaling factor of electron temperature $T_e$	1.0
SCTI	Scaling factor of ion temperature $T_i$	1.0
SCTIME	Scaling factor of time	1.0
MESH	Number of Lagrangian mesh cells ( $\leq 150$ )	40
NCASE	Boundary condition =1 $P(R_0) = 0$ , zero thermal flux =2 $u(R_0) = 0$ , $T_i(R_0) = T_i(t)$ , $T_e(R_0) = T_e(t)$	

	=3 $P(R_0) = P(t)$ , zero thermal flux =4 $u(R_0) = u(t)$ , zero thermal flux	
GEOM	Geometry =1 Slab =2 Cylinder =3 Sphere	
NITMAX	Maximum number of iteration for convergence criterion	5
NP1	Physical quantities $T_i$ , $T_e$ , $\rho$ etc are printed from NP1 mesh point to NP2 points	1
NP2	Last mesh number for printing the physical quantities	MESH
NP3	The physical quantities are printed for mesh interval of NP3 increment	<u>MESH</u> 20
NPRNT	Print at each NPRNT time step	100
NLPRNT	Print option	T
NPRINT	Logical unit number of print	6
NIN	Logical unit number of read	5
NPUNCH	Logical unit number of punch	7
NRUN	Maximum number of time steps	0
NLABS	Absorption by inverse-bremsstrahlung	T
NLBRMS	Bremsstrahlung	T
NLBURN	Burnup of fuel	T
NLCRI1	Dump laser power at critical density $\rho_c$	T
NLDEPO	Deposition of thermonuclear energies	T
NLECON	Electron thermal conduction	T
NLICON	Ion thermal conduction	T
NLFUSE	Thermonuclear reactions	T
NLITE	Solve equations by iterations	T
NLMOVE	Hydrodynamic motion	T
NLPFE	Perfect electron gas laws	T
NLPFI	Perfect ion gas laws	T



NLX	Ion-electron energy change	T
LASER 1(1)	Starting time for laser irradiation (sec)	0.0
LASER 1(2)	Initial laser power $P_0$ (w)	0.0
LASER 1(3)	Laser pulse duration $\tau$ (sec)	0.0
LASER 1(4)	Laser pulse shape factor q	0.0
LASER 1(5)	Time to switch off the laser (sec)	0.0
LASER 1(6)	Maximum permissible power level (w)	0.0
LASER 1(7)	Total maximum energy by laser irradiation (joule)	0.0
PIQ (10)	Modify free streaming limit coefficient, a in Eq.(A-6), $a' = a \times \text{PIQ}(10)$	0.0
PIQ (11)	=1, initial mesh width is divided by equal volume	0.0
PIQ (12)	=1, change initial density distribution, data of density are read in "namelist" INRHO	0.0
PIQ (13)	=1, change initial composition of pellet, composition data are read in "namelist" INFRAC	0.0
PIQ (20)	Modify coulomb logarithm coefficient $\Lambda$ in Eq.(A-4), $\Lambda' = \Lambda \times \text{PIQ}(20)$	0.0
PIQ (21)	Modify absorption coefficient $\alpha$ in Eq.(A-16), $\alpha' = \alpha \times \text{PIQ}(21)$	0.0
PIQ (24)	Modify bremsstrahlung coefficient, J in Eq.(A-11), $J' = J \times \text{PIQ}(24)$	0.0
PIQ (25)	Modify rate of energy exchange, K in Eq.(A-9), $K' = K \times \text{PIQ}(25)$	0.0
PIQ (26)	Modify ion conduction coefficient, $K_i$ in Eq.(A-3), $K_i' = K_i \times \text{PIQ}(26)$	0.0
PIQ (27)	Modify electron conduction coefficient, $K_e$ in Eq.(A-2), $K_e' = K_e \times \text{PIQ}(27)$	0.0
PIQ (28)	$P_{DT}^n = T_e^n / [T_e^n + 3.71 \times 10^8 (1 + \text{PIQ}(28))] in Eq.(A-27)$	0.0
PIQ (29)	$P_{DD}^n = T_e^n / [T_e^n + 1.2 \times 10^9 (1 + \text{PIQ}(29))] in Eq.(A-28)$	0.0
PIQ (30)	$P_{D^3He}^n = T_e^n / [T_e^n + 1.2 \times 10^9 (1 + \text{PIQ}(30))] in Eq.(A-29)$	0.0

PIQ (31)	Print option for coordinate	} if PIQ < 0.0, none print = 0.0, only array > 0.0, array and scale factor	0.0
PIQ (32)	Print option for velocity		0.0
PIQ (33)	Print option for density		0.0
PIQ (34)	Print option for pressure		0.0
PIQ (35)	Print option for Ti		0.0
PIQ (36)	Print option for Te		0.0
PIQ (38)	Modify reaction rate in Eq.(A-24) $(\overline{\sigma v})'_{DD} = (\overline{\sigma v})_{DD} (1 + \text{PIQ}(38))$		0.0
PIQ (39)	$(\overline{\sigma v})'_{DT} = (\overline{\sigma v})_{DT} (1 + \text{PIQ}(39))$ in Eq.(A-22)		0.0
PIQ (40)	$(\overline{\sigma v})'_{DT} = 7.5 \times 10^{-22} (7 + \text{PIQ}(40))$ in Eq.(A-23)		0.0
PIQ (41)	$(\overline{\sigma v})'_{D^3\text{He}} = (\overline{\sigma v})_{D^3\text{He}} (1 + \text{PIQ}(41))$ in Eq.(A-26)		0.0
NL PLOT	Switch option of plotter or COM production		F
N PLOT	Frequency of dumping physical variables for plotting production. The variables are written on disk or tape of unit 20. The graphical representations of the physical variables can be performed with the MEDU-PLOT <sup>(4)</sup> code.		10

# 3 Namelist INRAD if PIQ(11)  $\neq$  0.0

(BUF 1(J), J = 1, MESH + 1) : coordinates r

if PIQ(11) > 0.0

$$R(J) = R_0 \times (J-1)^{1/3} / (\text{MESH})^{1/3}$$

if PIQ(11) < 0.0

$$R(J) = \text{BUF } 1(J)$$

$$R_0 = \text{BUF } 1(\text{MESH} + 1)$$

# 4 Namelist INRHO if PIQ(12)  $\neq$  0.0

(BUF 3(J), J = 1, MESH + 1) : densities  $\rho$

if PIQ(12) > 0.0

$$\rho_J = \rho_0 \times \text{BUF}3(J)$$

if PIQ(12) < 0.0

$$\rho_J = \text{BUF}3(J)$$

# 5 Namelist INFRAC if PIQ(13) ≠ 0.0

$$(\text{BUF } 1(J), J = 1, \text{ MESH} + 1) : f_{D,J} = f_D \times \text{BUF } 1(J)$$

$$(\text{BUF } 2(J), J = 1, \text{ MESH} + 1) : f_{H,J} = f_H \times \text{BUF } 2(J)$$

$$(\text{BUF } 4(J), J = 1, \text{ MESH} + 1) : f_{N,J} = f_N \times \text{BUF } 4(J)$$

$$(\text{BUF } 5(J), J = 1, \text{ MESH} + 1) : f_{T,J} = f_T \times \text{BUF } 5(J)$$

$$(\text{BUF } 6(J), J = 1, \text{ MESH} + 1) : f_{X,J} = f_X \times \text{BUF } 6(J)$$

# 6 Namelist INTEM if PIQ(14) < 0.0

$$(\text{BUF } 5(J), J = 1, \text{ MESH} + 1) : T_{i,J} = \text{BUF } 5(J)$$

$$(\text{BUF } 6(J), J = 1, \text{ MESH} + 1) : T_{e,J} = \text{BUF } 6(J)$$

# 7 Namelist PIJRUN

Variable	Descriptions	Default value
NGR	Order of division of Gauss integration for radial direction	8
NGA	Order of division of Gauss integration for angular direction	2
NDA	Order of division for angular integration	1
NDR	Number of coarse mesh regions for collision probability calculation, $\leq 10$	10
NNR	Lagrangian mesh number corresponding to coarse mesh number	
SIGD	Neutron elastic scattering cross sections of D	Values in Table 1
SIGH	Neutron elastic scattering cross sections of H	Values in Table 1
SIGHE 3	" of $^3\text{He}$	"
SIGHE 4	" of $^4\text{He}$	"
SIGT	" of T	"
MUD	$\mu = \overline{\cos\theta}$ for D, $\theta$ is the scattering angle in the center of mass system	-1
MUT	$\mu = \overline{\cos\theta}$ for T, $\theta$ is the same described above.	-1

YIELDS	When the yield calculated become larger than YIELDS, calculations of neutron heating are performed	$1 \times 10^{-3}$
NLPIJ	Switch option for neutron heating calculation	F

### 3.2 Usage of MEDUSA-PIJ in the computer of FACOM 230/75

The jobcontrol cards of MEDUSA-PIJ in FACOM-230/75 are arranged as follows:

```

1   ¥ NO RMEDU,0350 ,      T.5C.4W.4P.0 ,      LRGC35
2   ¥ USER      J2031 ,
3 VM ¥ GJOB      3102031,HI,TAKANO,431,12,MEDUSA.P,SMF=CLS
4   ¥
5   ¥           IF COM PRODUCTION REQUIRED COMLIB=CALL
6   ¥           IF PLOT PRODUCTION REQUIRED PLTLIB=CALL
7 VM ¥ HLIED     RFNAME=J2031,MDS1HRB,GRFD=ON,COMLIB=CALL,
              SIMPL=OVLY,A=NOLIST,B=NOMAP,SYSOUS=CLS
8 VM ¥ HRUN     SYSOUT=CLS,SIZE=20
9   ¥           LOGICAL UNIT NUMBER F12 IS SCRATCH UNIT FOR PIJ
              CALCULATION
10 VM ¥ TPDISK   F12,PIJWRK,RSIZE=216,BSIZE=12960
11  ¥           LOGICAL UNIT NUMBER F20 IS SCRATCH UNIT FOR
              COM AND PLOT
12 VM ¥ TPDISK   F20,PLTWRK,RSIZE=910,BSIZE=6370
13  ¥           IF PLOT PRODUCTION REQUIRED REPLACE NEXT CARD TO
              ¥ PLOT
14  ¥           IF NOT ¥ DISK   F77
15 VM ¥ DISK     F77
16  ¥           IF COM PRODUCTION REQUIRED REPLACE NEXT CARD TO
              ¥ GCOM35
17  ¥           IF NOT ¥ DISK   F78

```

18 VM ¥ GCOM35  
 19 VM ¥ PRTFD F51, SYSOUT=CLS  
 20 VM ¥ DATA

Title cards (4)

\$ NEWRUN  
 Input variable = x ,  
 " = z

\$ PIJRUN  
 Input variable = y

\$ END

21 ¥ JEND

### 3.3 Output Data

# 1 The list of the input data

All the input data in Section 3.2 are printed.

# 2 At each time step, the following quantities are printed: time t(sec),

$\Delta t$ ;  $R_o$ , u, P,  $T_i$ ,  $T_e$  and  $C_s$  at outer boundary ;  $E_{th}(=PV/r-1)$ ;

$E_k(= \frac{1}{2} mv^2)$ ,  $E_f(=Y_i + Y_e)$ ;  $\Delta E(=E_k + E_{th} - E_f - E_{in} + E_b)$ , where  $E_{in}$  is  
 the input laser power and  $E_b$  the radiation loss ;  $\langle \rho R \rangle = \int_0^{R_o} \rho R dr /$

$\int_0^{R_o} dr$  ; Energy gain =  $(E_f + E_n)/E_{in}$  ; Produced neutrons ;

$E_n(=R_{DT}E_{DT} + \frac{1}{2} R_{DD}E_{DD})$  ; Laser power  $P_t(J)$

# 3 Lagrangian radial coordinates from center to outer boundary (m)

# 4 Hydrodynamic velocities (m/sec)

# 5 Densities  $\rho(kg/m)$  : (=log  $\rho$ )

# 6 Hydrodynamic pressures  $P(J/m^3)$  : (log P)

# 7 Ion temperatures  $T_i(^{\circ}K)$  : log  $T_i$

# 8 Electron temperatures  $T_e(^{\circ}K)$  : log  $T_e$

- # 9 Thermonuclear energy
- #10 Neutrons released by D-D reactions
- #11 Neutrons released by D-T reactions
- #12 Neutron energy produced by D-D reactions
- #13 Neutron energy produced by D-T reactions
- #14 Neutrons with the first collision at the cell j
- #15 Neutrons with the second collision at the cell j
- #16 Energy deposited by the first collision
- #17 Energy deposited by the second collision
- #18 Energy loss of neutrons by the first collision
- #19 Energy loss of neutrons by second collision
- #20 Escape neutron spectrum
- #21 Escape neutron probabilities  $P_{gj}$ , which the neutron produced at the region j and energy group g escape from the system without the next collision

## Acknowledgements

The authors are indebted to Mr. K. Tsuchihashi for his offering to use collision probability calculation code. They also wish to thank Mr. K. Nakano of CRC for programming MEDUSA-PIJ code.

## References

- 1) BRUECKNER, K.A., JORNA, S. : Rev. Mod. Phys., 46, No.2, 325 (1974)
- 2) ROBERTS, K.V. : Comp. Phys. Comm., 7, 237 (1974)
- 3) TSUCHIHASHI, K., GOTOH, Y. : Nucl. Sci. Eng., 58, 213 (1975)
- 4) TAKANO, H., SUWA, to be published
- 5) SPITZER, L. : "Physics of fully ionized gases", 2nd Ed., Interscience, New York (1961)
- 6) ROSE, D.J., CLARKE, M. : "Plasma and Controlled Fusion", Jhon Wilay and Sons Inc., New York, London (1961)
- 7) KIDDER, R.E. : "Interaction of Intense Photon and Electron Beams with Plasmas", Proc. Int. School Phys., Academic Press (1971)
- 8) KIDDER, R.E., BARNES, W.S. : "WAZER, A one-dimensional two temperature hydrodynamic code", UCRL-50583 (1969).
- 9) CHRISTIANSEN, J.P., et al. : CLM-RBO (1973)
- 10) RICHTMEYER, R.D., MORTON, K.W. : "Difference methods for initial value problems", 2nd Ed. Inter. Publishers (1967)

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- 3) TSUCHIHASHI, K., GOTOH, Y. : Nucl. Sci. Eng., 58, 213 (1975)
- 4) TAKANO, H., SUWA, to be published
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- 6) ROSE, D.J., CLARKE, M. : "Plasma and Controlled Fusion", Jhon Wilay and Sons Inc., New York, London (1961)
- 7) KIDDER, R.E. : "Interaction of Intense Photon and Electron Beams with Plasmas", Proc. Int. School Phys., Academic Press (1971)
- 8) KIDDER, R.E., BARNES, W.S. : "WAZER, A one-dimensional two temperature hydrodynamic code", UCRL-50583 (1969).
- 9) CHRISTIANSEN, J.P., et al. : CLM-RBO (1973)
- 10) RICHTMEYER, R.D., MORTON, K.W. : "Difference methods for initial value problems", 2nd Ed. Inter. Publishers (1967)



Table 1. Neutron elastic scattering cross sections (barns)

Group No.	E <sub>g</sub> (MeV)	D	T	α	<sup>3</sup> He <sub>2</sub>	H
1	14.1 ~11.31	0.7	1.1	1.25	1.1	0.75
2	11.31~9.07	0.9	1.4	1.45	1.3	0.90
3	9.07 ~7.27	1.1	1.6	1.60	1.5	1.1
4	7.27 ~5.83	1.3	1.9	1.80	1.7	1.25
5	5.83 ~4.68	1.55	2.2	2.15	1.9	1.45
6	4.68 ~3.75	1.80	2.4	2.35	2.0	1.75
7	3.75 ~3.01	2.0	2.5	2.60	2.1	2.10
8	3.01 ~2.41	2.2	2.4	3.00	2.15	2.40
9	2.41 ~1.94	2.4	2.2	3.65	2.15	2.75
10	1.94 ~1.56	2.6	2.0	4.90	2.1	3.2
-----						
11	1.56 ~1.25	2.7	1.9	6.40	2.1	3.6
12	1.25 ~1.00	2.8	1.8	7.30	2.0	4.0
13	1.00 ~0.803	2.9	1.75	5.70	1.9	4.4
14	0.803~0.644	3.0	1.70	3.30	1.9	4.9
15	0.644~0.517	3.1	1.6	2.0	1.9	5.3
16	0.517~0.414	3.1	1.5	1.4	1.9	5.8
17	0.414~0.332	3.15	1.45	1.1	1.9	6.2
18	0.332~0.267	3.15	1.42	0.91	1.9	6.7
19	0.267~0.214	3.15	1.40	0.83	1.9	8.5
20	0.214~0.172	3.20	1.35	0.80	1.9	10.0
-----						
21	0.172~0.138	3.20	1.35	0.78	1.8	11.5
22	0.138~0.110	3.20	1.35	0.77	1.7	13.0
23	0.110~0.886	3.20	1.35	0.76	1.6	14.0

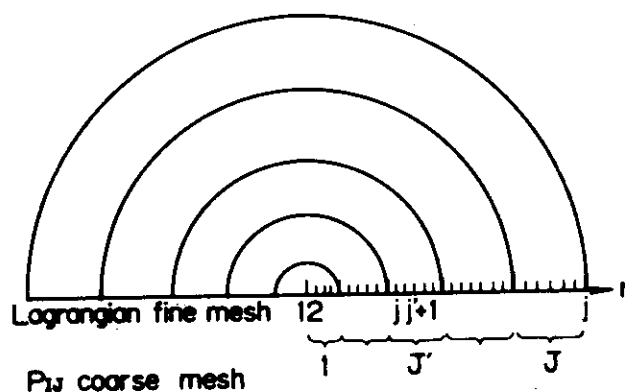


Fig. 1 Lagrangian mesh and coarse mesh for calculation of collision probability

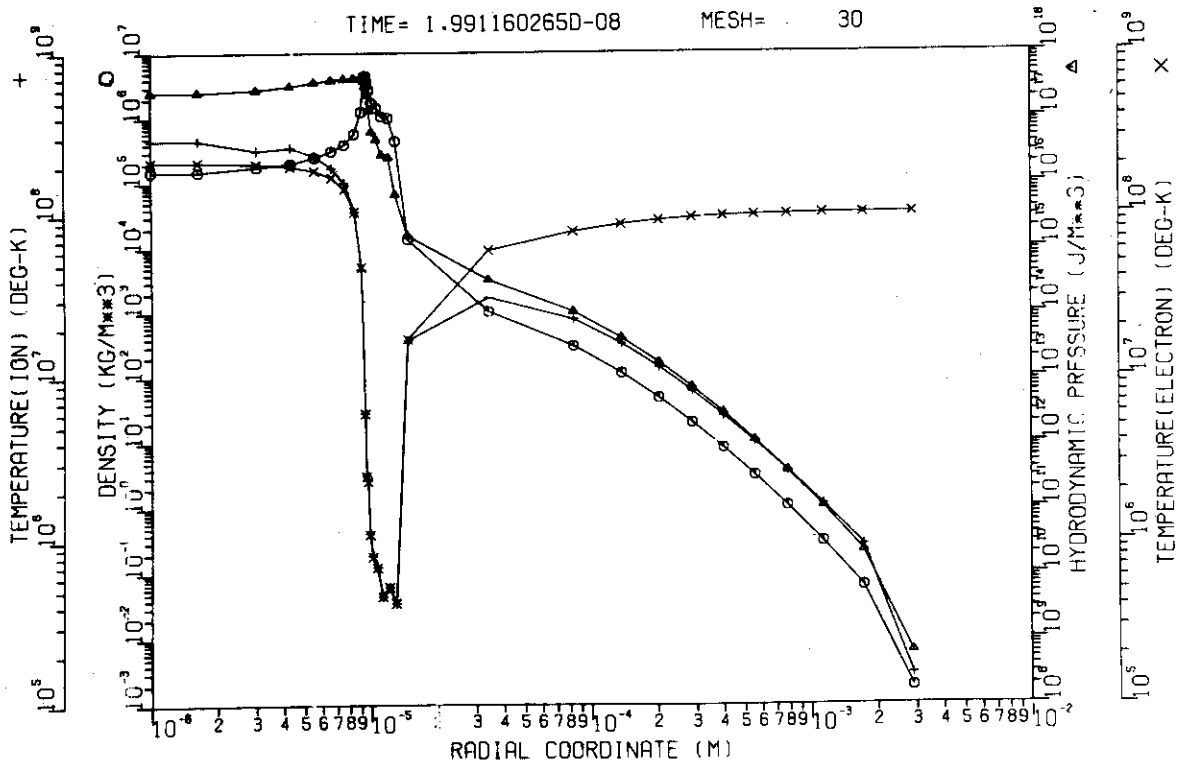


Fig. 2 Density, hydrodynamic pressure, electron and ion temperatures versus radius during the implosion of D-T pellet

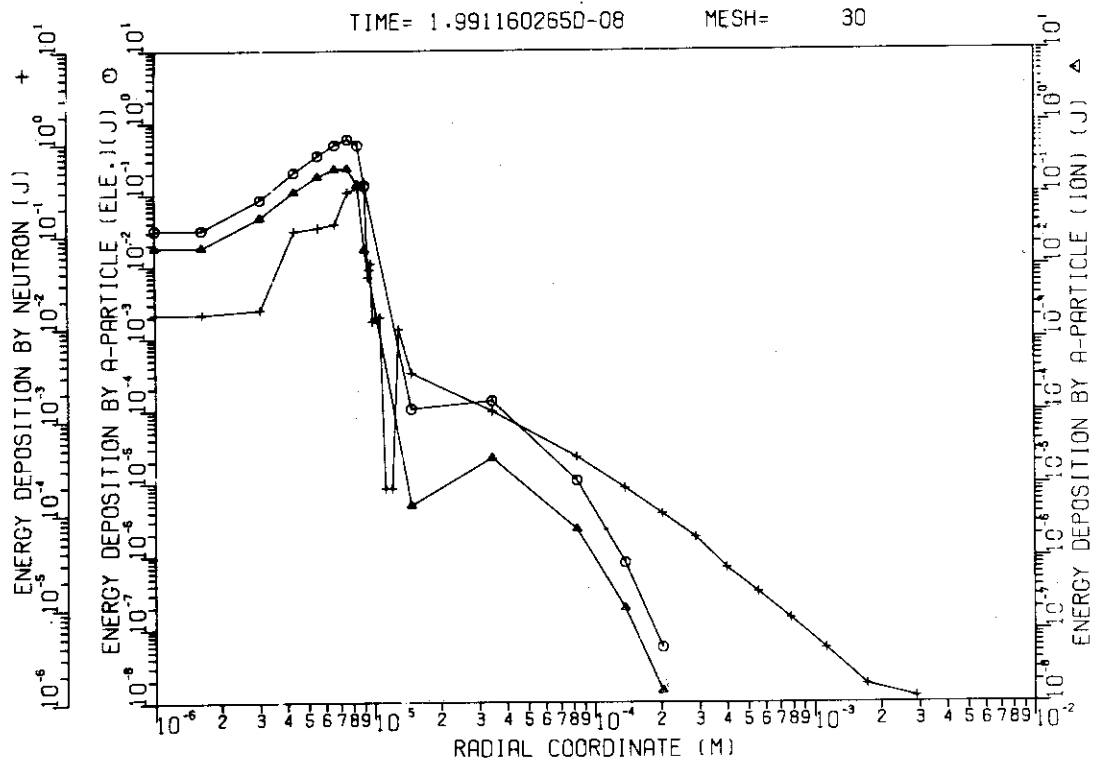


Fig. 3 Comparison of energy deposition by neutrons and alpha-particles

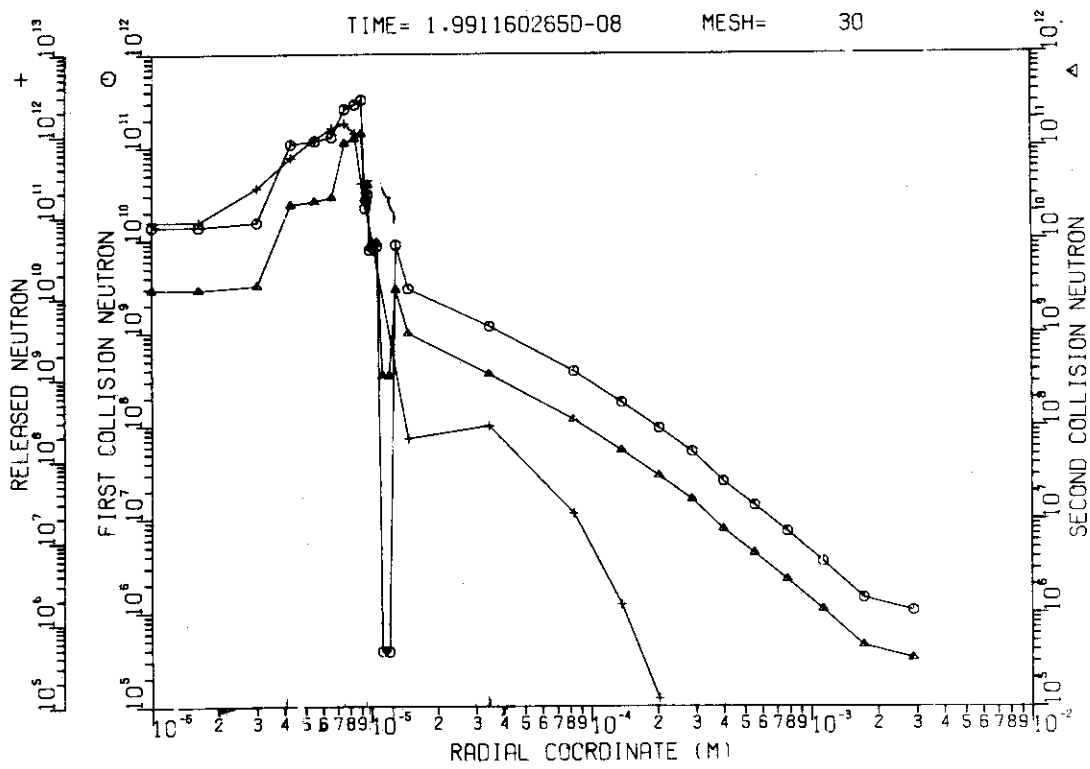


Fig. 4 Fusion neutrons and the first and second collision neutrons

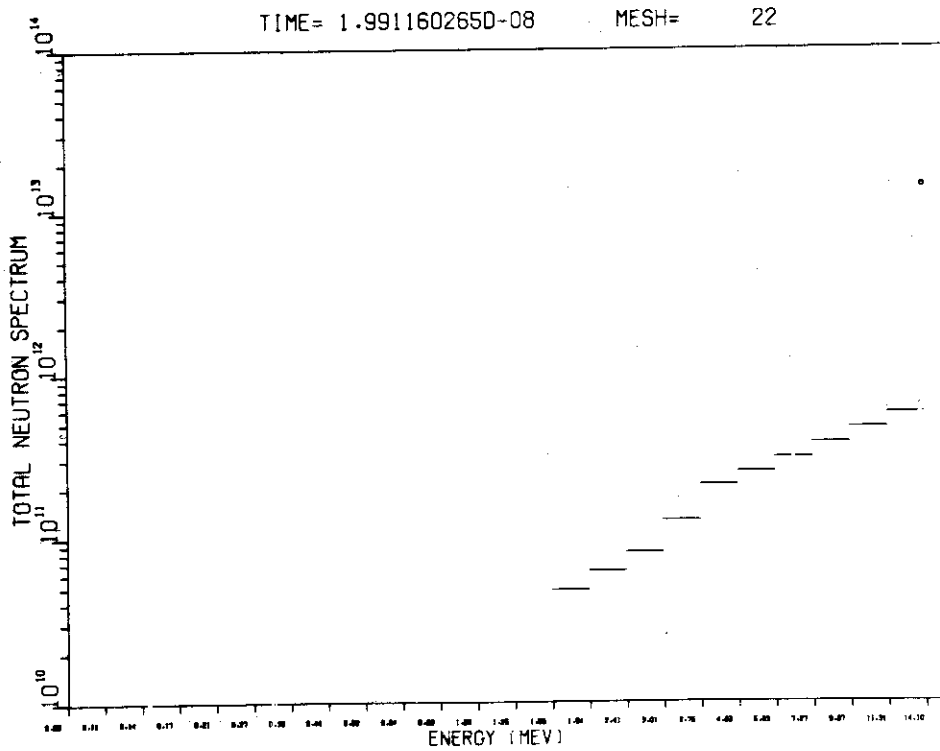


Fig. 5 Neutron energy spectrum degraded by elastic scattering

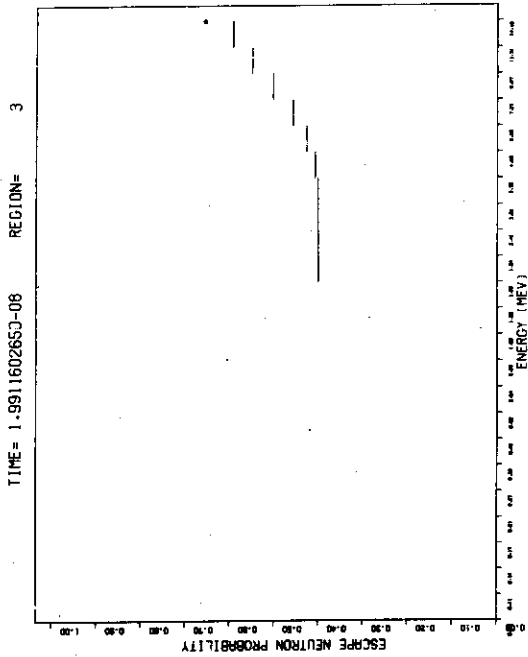


Fig. 8 Escape neutron probability depending on each energy group

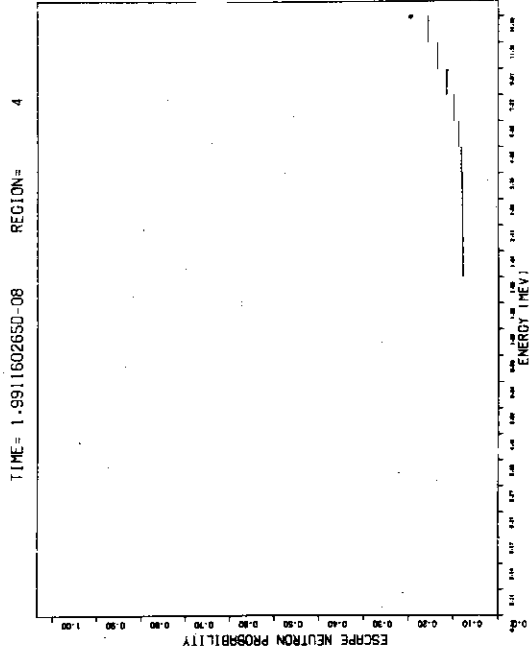


Fig. 9 Escape neutron probability depending on each energy group

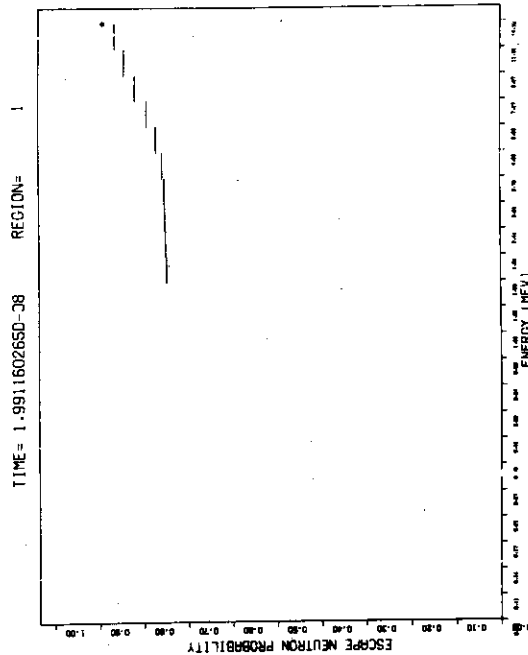


Fig. 6 Escape neutron probability depending on each energy group

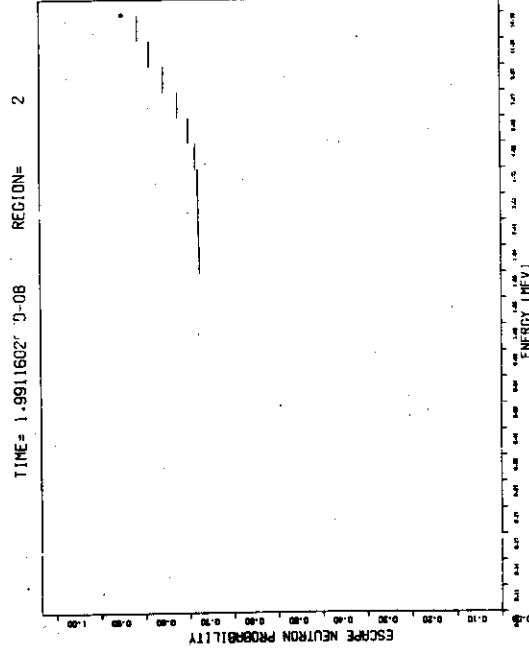


Fig. 7 Escape neutron probability depending on each energy group

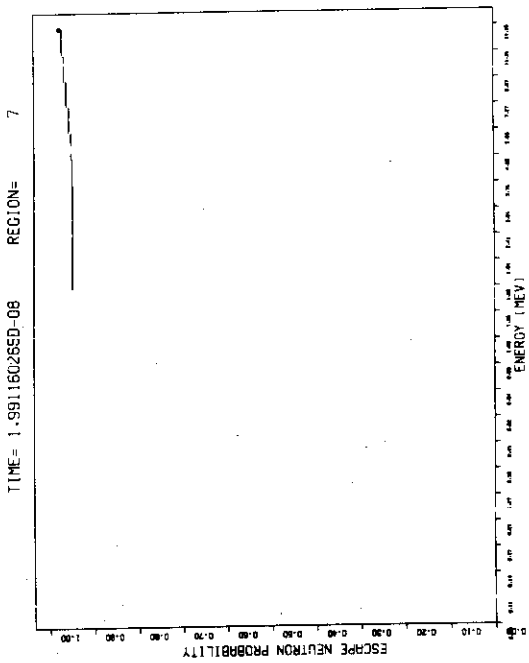


Fig. 12 Escape neutron probability depending on each energy group

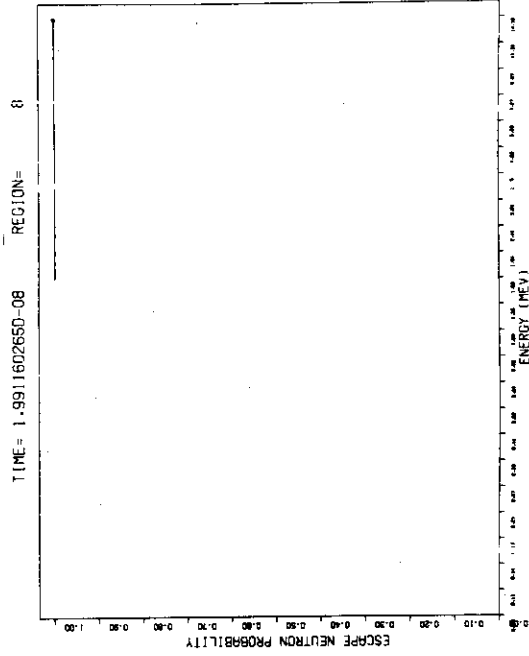


Fig. 13 Escape neutron probability depending on each energy group

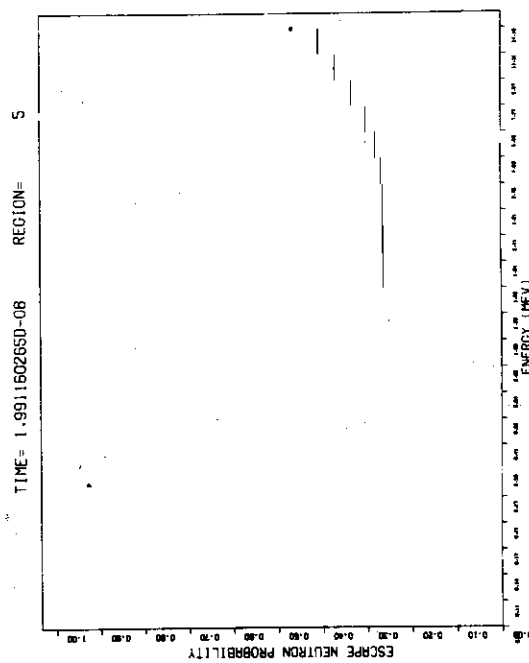


Fig. 10 Escape neutron probability depending on each energy group

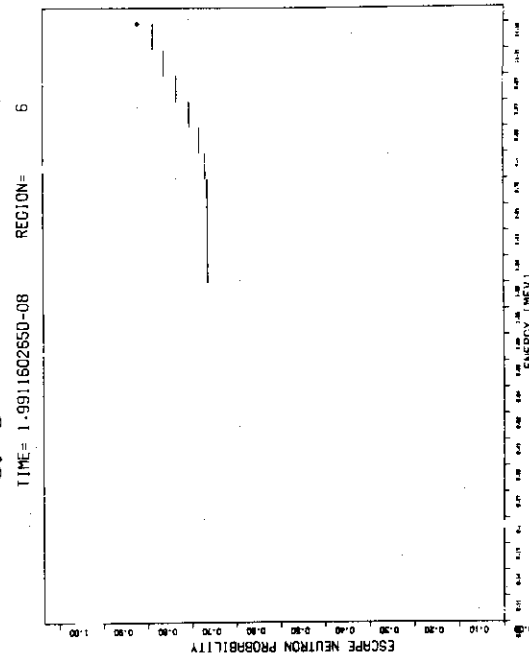


Fig. 11 Escape neutron probability depending on each energy group

Appendix A. Expressions for Various Terms in Energy Balance Equations

a) Thermal conductivity (H)

$$H = \frac{1}{\rho} \nabla \cdot k \nabla T \quad (A-1)$$

where  $\rho$  is the density of plasma,  $T$  the temperature of plasma and  $k$  the thermal conductivity. The expressions for  $k$  by Spitzer<sup>(5)</sup> are given by

$$k_e = 1.83 \times 10^{-10} T_e^{5/2} / (Z \log \Lambda) \quad (A-2)$$

$$k_i = 4.30 \times 10^{-12} T_i^{5/2} / (M^{1/2} Z^4 \log \Lambda) \quad (A-3)$$

$$\Lambda = 1.24 \times 10^7 T_e^{3/2} n_e^{-1/2} / Z \quad (A-4)$$

where the subscripts  $e$  and  $i$  stand for the electron and ion, respectively,  $M$  the atomic mass number,  $Z$  the charge number and  $\log \Lambda$  the Coulomb logarithm.

The electron thermal flux  $F_e = k_e \nabla T_e$  is limited by the so-called streaming limit.<sup>(2)</sup>

$$\frac{1}{F_e'} = \frac{1}{F_e} + \frac{1}{F_{e,max}} \quad (A-5)$$

$$F_{e,max} = \frac{a}{4} n_e \bar{v}_e k T_e \quad (A-6)$$

$$k_e' = k_e \left( 1 + a \frac{\lambda_e}{T_e} \frac{dT_e}{dx} \right)^{-1} \quad (A-7)$$

$$\lambda_e = 7.92 \times 10^9 T_e^2 / n_i Z^2 \log \Lambda \quad (m) \quad (A-8)$$

where  $a$  is an adjustable constant for  $(F_e)_{max}$  and  $\lambda_e$  the electron mean free path.

b) Electron-ion collisional exchange<sup>(6)</sup> (K)

$$K = \frac{\omega}{\rho} n_e (T_i - T_e) = 0.59 \times 10^{-8} n_e (T_i - T_e) T_e^{3/2} \frac{M^{-1} Z^{-2}}{m_H^{-1}} \log \Lambda k M^{-1} m_H^{-1} \quad (W/kg) \quad (A-9)$$

$$\omega = \frac{M^{-1}Z^{-2} e^4 n_i \log \Lambda m_e^{1/2}}{32\sqrt{2\pi} \epsilon_0^2 m_H} (k T_e)^{-3/2} \quad (\text{A-10})$$

c) Radiation loss<sup>(6)</sup> (J)

$$J = - \frac{Z^2 e^6 n_e \bar{v}_e}{24\pi \epsilon_0^3 c^3 m_e m_H M h} \quad (\text{A-11})$$

$$= - 8.5 \times 10^{-14} n_e (\delta T)^{1/2} \frac{Z^2 M^{-1}}{Z^2 M^{-1}} \quad (\text{A-12})$$

where

$$\delta T = T_e - (n_e/n_{e^0})^{\gamma_e-1} T_{e^0} \quad (\text{A-13})$$

d) Laser absorption<sup>(7)</sup> (X)

$$X(r,t) = \frac{1}{dM} P_L(r,t) (1 - e^{-\alpha(R_0-r)}) \quad (\text{W/kg}) \quad (\text{A-14})$$

$$P_L(r,t) = P_L(R_0,t) e^{-\alpha(R_0-r)} \quad (\text{A-15})$$

$$\alpha = 13.5 \lambda^{-2} \beta^2 (1-\beta)^{-1/2} T_e^{-3/2} [5.05 + \log(\lambda T_e)] \bar{Z} \quad (\text{m}^{-1}) \quad (\text{A-16})$$

where  $P_L(R_0,t)$  is the laser power incident on the plasma boundary at  $r = R_0$ ,  $\alpha$  the absorption coefficient,<sup>(7)</sup>  $\lambda$  the wavelength of laser light and  $\beta = \rho/\rho_c < 1$ . At the critical surface ( $r=r_c$ ,  $\rho=\rho_c$ ,  $\omega_p=\omega_c$ ), the remaining laser power is all assumed to be absorbed.

e) Thermonuclear processes

$$- \frac{dN_D}{dt} = - \frac{dN_T}{dt} = \frac{dN_{^4\text{He}}}{dt} = \frac{dN_n}{dt} = N_D N_T (\overline{\sigma v})_{DT} \quad (\text{A-17})$$

$$\frac{dN_D}{dt} = - N_D^2 (\overline{\sigma v})_{DD} - N_D N_T (\overline{\sigma v})_{DT} - N_D N_{^3\text{He}} (\overline{\sigma v})_{D^3\text{He}} \quad (\text{A-18})$$

$$\frac{dN_T}{dt} = -N_D N_T (\overline{\sigma v})_{DT} + \frac{1}{4} N_D^2 (\overline{\sigma v})_{DD} \quad (\text{A-19})$$

$$\frac{dN_n}{dt} = \frac{1}{4} N_D^2 (\overline{\sigma v})_{DD} \quad (\text{A-20})$$

$$\frac{dN_{^3\text{He}}}{dt} = \frac{1}{4} N_D^2 (\overline{\sigma v})_{DD} \quad (\text{A-21})$$

where the reaction rates are given as the function of ion temperature as follows:

$$(\overline{\sigma v})_{DT} = \begin{cases} 3.68 \times 10^{-18} T_i^{-3/2} \exp(-19.94 \times T_i^{-3/2}) & (\text{m}^3/\text{sec}) \\ 7.5 \times 10^{-22}, & T_i > 35 \text{ keV} \end{cases} \quad (\text{A-22})$$

$$\quad \quad \quad (\text{A-23})$$

$$(\overline{\sigma v})_{DD} = 2.23 \times 10^{-20} T_i^{-2/3} \exp(-19.42 T_i^{-1/3}) \quad (\text{A-24})$$

$$(\overline{\sigma v})_{D^3\text{He}} = \begin{cases} 0 \\ 10^{-23}, & T_i > 50 \text{ keV} \end{cases} \quad (\text{A-25})$$

$$\quad \quad \quad (\text{A-26})$$

The charged reaction products H, T,  $^3\text{He}$  and  $^4\text{He}$  are assumed to deposit their energy locally to the ions and electrons in ratios as a function of  $T_e$  (8):

$$P_{DT} = \frac{T_e}{T_e + 3.71 \times 10^8} \quad (\text{A-27})$$

$$P_{DD} = \frac{T_e}{T_e + 1.2 \times 10^9} \quad (\text{A-28})$$

$$P_{D^3\text{He}} = \frac{T_e}{T_e + 1.2 \times 10^9} \quad (\text{A-29})$$

Then, the thermonuclear energy release is

$$Y_i = P_{DD} E_{DD} R_{DD} + P_{DT} E_{DT} R_{DT} + P_{D^3\text{He}} E_{D^3\text{He}} R_{D^3\text{He}} + \langle DE \rangle \quad (\text{A-30})$$



$$Y_e = (1-P_{DD})E_{DD}R_{DD} + (1-P_{DT})E_{DT}R_{DT} + (1-P_{D^3He})E_{D^3He}R_{D^3He} \quad (A-31)$$

where  $E_{XX}$  is the energy of the charged reaction product,  $R_{XX}$  the total reaction rate and  $\langle DE \rangle$  the energy deposition of the emitted neutrons.

f) Shock heating<sup>(9)</sup> (Q)

$$Q = -q \frac{dV}{dt} \quad (A-32)$$

where  $q$  is the viscous pressure expressed as

$$q = b^2 \frac{1}{2} \rho u \eta \quad (A-33)$$

$$\eta = (\partial r \nabla u) / (\partial r \nabla \cdot u) = r^{g-1} \partial u / \partial r^{g-1} u \quad (A-34)$$

g) Boundary condition

The MEDUSA code does not specifically allow for hollow shells, and the inner boundary conditions are  $u(r=0) \equiv 0$  and zero thermal flux. At the outer boundary point  $r = R_0$ , the following four cases can be chosen:

- (1)  $\underline{P}(R_0) = 0$ ; zero thermal flux
- (2)  $u(R_0) = 0$ ;  $T_i(R_0) = T_i(t)$ ;  $T_e(R_0) = T_e(t)$
- (3)  $\underline{P}(R_0) = P(t)$ ; zero thermal flux
- (4)  $u(R_0) = u(t)$ ; zero thermal flux

h) Control of convergence by iteration for solving difference equations

The difference equations used in the calculation of MEDUSA are in the Lagrange formulation, i.e. in a reference frame moving with the plasma. For solving the difference equations, the arrangement of the space and time meshes are devised for various basic quantities. Several iterations are performed on one time-step, because a non-linear dependence of the physical terms  $H$ ,  $K$ ,  $Y$ ,  $J$  and  $X$  in Eqs.(18) and (19) on the temperature  $T_i$  and  $T_e$ . The difference equations are numerically solved with the explicit and implicit

methods<sup>(10)</sup> as described in Ref.(2). The time-interval  $\Delta t$  for time step  $n$  are automatically determined by the following conditions:

$$\Delta t^{n+\frac{1}{2}} < a_1 M_{in} \left( \frac{R_{j+1}^n - R_j^n}{C_\ell^n} \right) \quad (A-35)$$

$$\Delta t^{n+\frac{1}{2}} \leq a_2 M_{in} \left( \frac{V^{n+1} - V^n}{V^{n+1} + V^n} \right)_\ell \Delta t^{n-1/2} \quad (A-36)$$

$$\Delta t^{n+\frac{1}{2}} \leq a_3 M_{in} \left( \frac{T_i^{n+1} - T_i^n}{T_i^{n+1} + T_i^n} \right)_\ell \quad (A-37)$$

$$\Delta t^{n+\frac{1}{2}} \leq a_4 M_{in} \left( \frac{T_e^{n+1} - T_e^n}{T_e^{n+1} + T_e^n} \right)_\ell \quad (A-38)$$

$$\Delta t^{n+\frac{1}{2}} \leq a_0 \Delta t^{n-\frac{1}{2}} \quad (A-39)$$

where  $a_1, a_2, a_3$  and  $a_4$  are respectively constants less than 1,  $a_0$  the time-centering constant and  $C_\ell^n$  the sound speed in cell  $\ell$ . The smallest  $\Delta t$  satisfying these conditions is found and the iterations on  $T_i, T_e$  and  $u$  are performed by checking the following convergence conditions:

$$M_{ax} \left( \frac{|u^m - u^{m-1}|}{|u^m + u^{m-1}|} \right)_j^{n+1/2} \leq \delta u \quad (A-40)$$

$$M_{ax} \left( \frac{T_i^m - T_i^{m-1}}{T_i^m + T_i^{m-1}} \right)_j^{n+1/2} \leq \delta T_i \quad (A-41)$$

$$M_{ax} \left( \frac{T_e^m - T_e^{m-1}}{T_e^m + T_e^{m-1}} \right)_j^{n+1/2} \leq \delta T_e \quad (A-42)$$

where  $m$  indicates the iteration numbers.

## Appendix B. Example of Input Data

.....1.....2.....3.....4.....5.....6.....7.....8

```

TEST 1 20/9/1978
STEP 2 TEST 1
CO2 LASERLIGHT ON A 60
MGRM 50-50 DT PELLET
#NEWRUN

```

```

DELTAT=1,0E-13,
LAMDA1=1,0E-5,
LASER1(2)=1,5E9,
LASER1(3)=2,0E-8,
LASER1(4)=-2,0,
LASER1(5)=2,0E-8,
LASER1(6)=1,0E16,
LASER1(7)=4,0E3,
MESH=30,
NLDUMP=F,
NLFILM=F,
NLHCPY=F,
NLPFE=F,
NRUN=1300,
NPRNT=100,
TEINI=1,0E3,
TIINI=1,0E3,
PIQ(31)=1,
PIQ(32)=1,
PIQ(33)=1,
PIQ(34)=1,
PIQ(35)=1,
PIQ(36)=1,
PIQ(42)=1,0,
PIQ(43)=1,0,
PIQ(44)=1,0,
TSTOP=2,0-8,
TIMEE(1)=6*2,0E-8,
TIMES(1)=6*1,1E-8,
NPLOT=10,
NTYPE(1,1)=6*1,
NTYPE(1,2)=6*2,
NTYPE(1,3)=6*3,
NTYPE(1,4)=6*4,
NTYPE(1,5)=6*5,
NTYPE(1,6)=6*6,
OPT1=0,
NLPLOT=T

```

```

#END
#PIJRUN

```

```

NGR=8,
NDR=10,
NNR=1,4,7,10,13,16,18,20,25,30,
YIELDS=1,0-3,
NLPIJ=T

```

```

#END

```





TIMESTEP NUMBER 900  
 -----  
 DELTA T DETERMINED BY CONDITION 3 AT MESHPOINT 18  
 -----  
 TIME = 1.9914406D-08  
 DELTA T = 1.6902254D-14  
 -----  
 BOUNDARY : R = 9.0748D-03 U = 0.0 P = 1.0000D-16 T1 = 0.0 TE = 0.0 SOUND SPEED = 7.4845D+05  
 ENERGIES : THERMAL 1.00993D+04 KINETIC 3.39742D+03 NUCLEAR 2.12745D+04 ERROR -3.25997D+02 RHO R 3.35082D+00  
 FUSION : YIELD 5.29587D+00 NEUTRONS 7.48119D+15 RATE 1.33724D+28 ENERGY 1.11257D+04  
 DEPOSIT : ENERGY 5.78365D+03  
 RADIATION LOSS -3.32678D+02

COORDINATES (M) * 1.00000D+00	3.4211D-06	4.4927D-06	5.5108D-06	6.5085D-06	7.4408D-06	8.0909D-06	8.5598D-06
0.0	1.1792D-06	2.3106D-06	1.0305D-02	1.1375D-05	1.1739D-05	1.2286D-05	1.3851D-05
8.9463D-06	9.3135D-06	9.7917D-06	2.9093D-04	4.0255D-04	5.1695D-04	7.1818D-04	1.1327D-03
8.5642D-05	1.4066D-04	2.0688D-04	1.2366D+06	1.2570D+06	1.2447D+06	1.1983D+06	1.0098D+06
HYDRODYNAMIC VELOCITIES (M/SEC) * 1.00000D+00	-2.0197D+05	-2.6989D+05	-4.0784D+05	-5.8886D+05	-8.4630D+05	-1.5088D+06	-2.4272D+06
0.0	-3.1231D+06	-2.0433D+06	-8.2883D+05	2.6062D+05	1.1896D+05	2.3687D+06	1.2240D+06
-3.1231D+06	-2.0433D+06	-8.2883D+05	2.6062D+05	1.1896D+05	2.3687D+06	1.2240D+06	1.2240D+06
8.1843D+05	1.0608D+06	1.1775D+06	1.2366D+06	1.2570D+06	1.2447D+06	1.1983D+06	1.0098D+06
LOG(DENSITY (KG/M**3)) * 1.00000D+00	5.98316	5.51442	5.53403	5.59943	5.62313	5.66860	5.83661
5.98316	5.51442	5.53403	5.59943	5.62313	5.66860	5.83661	6.04683
6.25207	6.17805	6.17581	6.18241	6.19611	6.39584	6.23731	5.76006
2.47313	2.06432	1.68930	1.51848	0.92978	0.51142	0.04023	-0.50932
LOG(HYDRODYNAMIC PRESSURE (JOULE/M**3)) * 1.00000D+00	18.23207	18.22980	18.23604	18.29652	18.34998	18.48171	18.72574
18.23207	18.22980	18.23604	18.29652	18.34998	18.48171	18.72574	18.93620
19.13496	19.03903	19.02836	19.02103	19.00319	19.18118	18.95743	17.90906
13.99945	13.59708	13.22346	12.85025	12.46286	11.57118	11.02033	10.34118
LOG(ION TEMPERATURE (DEGREE K)) * 1.00000D+00	9.05587	9.00122	8.97439	8.95823	9.00222	9.13262	9.21147
9.05587	9.00122	8.97439	8.95823	9.00222	9.13262	9.21147	9.24008
9.23363	9.21374	9.20630	9.18416	9.11896	9.00308	7.95709	7.35075
7.36052	7.21023	7.06568	6.91972	6.76974	6.61035	6.43277	6.22727
LOG(ELECTRON TEMPERATURE (DEGREE K)) * 1.00000D+00	8.76498	8.76735	8.77230	8.77800	8.78649	8.79445	8.80200
8.76498	8.76735	8.77230	8.77800	8.78649	8.79445	8.80200	8.80472
8.79137	8.78032	8.76887	8.75742	8.74515	8.73652	8.71126	8.52346
7.89661	7.93949	7.96342	7.97862	7.98884	7.99583	8.00057	8.00365
RELEASED THERMONUCLEAR ENERGY (ELECTRON) * 1.00000D+00	1.9108D-03	1.4510D-02	4.1193D-02	8.4417D-02	1.4986D-01	2.3430D-01	3.6098D-01
1.9108D-03	1.4510D-02	4.1193D-02	8.4417D-02	1.4986D-01	2.3430D-01	3.6098D-01	7.3773D-01
3.7112D+00	3.6466D+00	4.3335D+00	5.1662D+00	6.2649D+00	1.2743D+01	1.1206D+01	2.1526D-01
2.2850D-06	1.7251D-07	1.2150D-08	0.0	0.0	0.0	0.0	0.0
RELEASED THERMONUCLEAR ENERGY (ION) * 1.00000D+00	2.8983D-03	2.2171D-02	6.3771D-02	1.3290D-01	2.4046D-01	3.8312D-01	5.9952D-01
2.8983D-03	2.2171D-02	6.3771D-02	1.3290D-01	2.4046D-01	3.8312D-01	5.9952D-01	1.2336D+00
6.0250D+00	5.7763D+00	6.7209D+00	7.7733D+00	9.6142D+00	1.8324D+01	1.5184D+01	1.8161D-01
4.1787D-07	4.0237D-08	2.9934D-09	0.0	0.0	0.0	0.0	0.0
RELEASED NEUTRONS D-D REACTION	4.4131D+07	3.0368D+08	8.3164D+08	1.7106D+09	2.9856D+09	5.0680D+09	9.7827D+09
4.4131D+07	3.0368D+08	8.3164D+08	1.7106D+09	2.9856D+09	5.0680D+09	9.7827D+09	2.2659D+10
1.1460D+11	1.0927D+11	1.2728D+11	1.4516D+11	1.7203D+11	3.1658D+11	2.1423D+11	1.4119D+09
1.11276D+04	9.1940D+02	6.9087D+01	0.0	0.0	0.0	0.0	0.0
RELEASED NEUTRONS D-T REACTION	8.1851D+09	6.2512D+10	1.7900D+11	3.7062D+11	5.6590D+11	1.9525D+12	1.6322D+12
8.1851D+09	6.2512D+10	1.7900D+11	3.7062D+11	5.6590D+11	1.9525D+12	1.6322D+12	3.9467D+12
1.6503D+13	1.5979D+13	1.8782D+13	2.1956D+13	2.7478D+13	4.5001D+13	6.8141D+13	1.1891D+08
4.7068D+06	3.6483D+05	2.5959D+04	0.0	0.0	0.0	0.0	0.0

RELEASED ENERGY D-D REACTION (NEUTRON)										
1.7397D-05	1.2050D-04	3.2784D-04	6.7432D-04	1.1769D-03	5.4211D-01	1.9978D-03	3.8564D-03	8.9324D-03	1.8690D-02	3.2492D-02
4.5726D-02	4.3676D-02	5.0176D-02	5.7222D-02	6.7722D-02	0.0	1.2480D-01	8.4451D-02	5.5657D-04	1.1330D-07	5.3015D-08
4.4450D-09	3.6243D-10	2.7234D-11	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
RELEASED ENERGY D-T REACTION (NEUTRON)										
1.8569D-02	1.4182D-01	4.0609D-01	8.8082D-01	1.5107D+02	5.0969D+02	2.3877D+00	3.7030D+00	7.5925D+00	1.5237D+01	2.5920D+01
3.7440D+01	3.6251D+01	4.4812D+01	4.9810D+01	6.2398D+01	0.0	1.1981D+02	1.0209D+02	1.5439D+00	2.6978D-04	1.5291D-04
1.0678D-05	8.2768D-07	5.8892D-08	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
FIRST COLLISION SLOWED DOWN NEUTRON										
2.2484D+10	2.3167D+10	2.3885D+10	1.5655D+11	1.6177D+11	9.2011D+13	1.6743D+11	1.3742D+12	1.4969D+12	1.6407D+12	7.3437D+12
7.6075D+12	8.2567D+12	1.5503D+13	1.6449D+13	1.7146D+13	1.7146D+13	6.3955D+12	6.9896D+12	9.1781D+11	2.7479D+11	3.5373D+10
1.1507D+10	5.4067D+09	2.8934D+09	1.6278D+09	1.6600D+08	0.0	4.2587D+08	2.2449D+08	1.0779D+08	4.4098D+07	3.2129D+07
SECOND COLLISION SLOWED DOWN NEUTRON										
9.0889D+09	9.3710D+09	9.6647D+09	6.2367D+10	6.4447D+10	5.47675D+13	6.6723D+10	7.7742D+11	8.4714D+11	9.2857D+11	4.7843D+12
4.9465D+12	5.3582D+12	1.0170D+13	1.0780D+13	1.2630D+13	0.0	2.0150D+12	2.2079D+12	3.3775D+11	1.0112D+11	1.6217D+10
5.2754D+09	2.4787D+09	1.3264D+09	7.4622D+08	5.3118D+08	0.0	1.9524D+08	1.0292D+08	4.9416D+07	2.0217D+07	1.4729D+07
FIRST COLLISION DEPOSIT NEUTRON ENERGY										
2.0263D-02	2.0878D-02	2.1525D-02	1.4112D-01	1.4582D-01	8.2868D+01	1.5094D-01	1.2400D-01	1.3507D+00	1.4805D+00	6.6115D+00
6.8471D+00	7.4315D+00	1.3929D+01	1.4779D+01	1.5406D+01	0.0	5.8064D+00	6.3458D+00	8.3706D-01	2.5061D-01	3.2262D-02
1.0495D-02	4.9311D-03	2.6388D-03	1.4846D-03	6.9862D-04	0.0	3.8841D-04	2.0474D-04	9.8306D-05	4.0219D-05	2.9303D-05
SECOND COLLISION DEPOSIT NEUTRON ENERGY										
4.7260D-03	4.8729D-03	5.0257D-03	3.2536D-02	3.3621D-02	2.9510D+01	3.4808D-02	4.1225D-01	4.4923D-01	4.9241D-01	2.5883D+00
2.6759D+00	2.8968D+00	5.4954D+00	5.8252D+00	6.0875D+00	0.0	1.0624D+00	1.1643D+00	1.7574D+01	5.2616D+02	8.7377D+03
2.8424D-03	1.3355D-03	7.1470D-04	4.0208D-04	1.8922D-04	0.0	1.0320D-04	5.3454D-05	2.6626D-05	1.0893D-05	7.9363D-06
FIRST COLLISION NEUTRON ENERGY LOSS										
1.5797D-03	1.2065D-02	3.4546D-02	7.1512D-02	1.2849D-01	8.28665D+01	2.0307D-01	5.1186D-01	1.0495D+00	2.1059D+00	5.5538D+00
8.0233D+00	7.7690D+00	9.0951D+00	1.0631D+01	1.3306D+01	0.0	1.3149D+01	1.1206D+01	1.3923D-02	2.4296D-06	3.8937D-08
3.1282D-09	2.4247D-10	1.7253D-11	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
SECOND COLLISION NEUTRON ENERGY LOSS										
3.8030D-03	3.9185D-03	4.0400D-03	2.6137D-02	2.7070D-02	2.9509D+01	2.7950D-02	3.4752D-01	3.7855D-01	4.1493D-01	2.6012D+00
2.6939D+00	2.9238D+00	5.4790D+00	5.8131D+00	6.0596D+00	0.0	1.2817D+00	1.4008D+00	1.7489D-02	5.2361D-03	2.1104D-05
6.8652D-06	3.2257D-06	1.7262D-06	9.7113D-07	1.6585D-08	0.0	9.2094D-09	4.8543D-09	2.3309D-09	9.5382D-10	6.9963D-11
ESCAPE NEUTRON SPECTRUM BY GROUP										
1.3262D+14	1.1274D+13	9.4313D+12	7.7419D+12	6.3119D+12	1.83762D+14	5.1787D+12	4.2294D+12	2.7765D+12	1.8345D+12	1.3636D+12
1.0022D+12	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
TOTAL NEUTRON SPECTRUM BY GROUP										
1.3262D+14	1.1274D+13	9.4313D+12	7.7419D+12	6.3119D+12	1.83762D+14	5.1787D+12	4.2294D+12	2.7765D+12	1.8345D+12	1.3636D+12
1.0022D+12	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
TOTAL RELEASED NEUTRON										
8.2292D+09	6.2818D+10	1.7983D+11	3.7233D+11	6.6889D+11	2.26024D+14	1.0575D+12	1.6420D+12	3.3693D+12	6.7635D+12	1.1507D+13
1.6619D+13	1.6088D+13	1.8910D+13	2.2101D+13	2.7650D+13	0.0	5.3127D+13	4.5216D+13	6.8282D+11	1.1920D+08	5.8720D+07
4.7181D+06	3.6575D+05	2.6028D+04	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
TOTAL RELEASED NEUTRON ENERGY										
1.8587D-02	1.4194D-01	4.0642D-01	8.4149D-01	1.5119D+00	5.10201D+02	2.3897D+00	3.7068D+00	7.6015D+00	1.5255D+01	2.5992D+01
3.7485D+01	3.6294D+01	4.2662D+01	4.9868D+01	6.2406D+01	0.0	1.1993D+02	1.0218D+02	1.5465D+00	2.6989D-04	1.3297D-04
1.0683D-05	8.2604D-07	5.6919D-08	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
TOTAL SLOWED DOWN NEUTRON										
2.2484D+10	2.3167D+10	2.3885D+10	1.5655D+11	1.6177D+11	9.2011D+13	1.6743D+11	1.3742D+12	1.4969D+12	1.6407D+12	7.3437D+12
7.6075D+12	8.2567D+12	1.5503D+13	1.6449D+13	1.7146D+13	1.7146D+13	6.3955D+12	6.9896D+12	9.1781D+11	2.7479D+11	3.5373D+10
1.1507D+10	5.4067D+09	2.8934D+09	1.6278D+09	1.6600D+08	0.0	4.2587D+08	2.2449D+08	1.0779D+08	4.4098D+07	3.2129D+07

TOTAL DEPOSIT NEUTRON ENERGY									
2.4989D-02	2.5751D-02	2.6531D-02	1.7365D-01	1.7944D-01	1.8575D-01	1.6322D+00	1.7999D+00	1.9729D+00	9.1998D+00
9.5231D+00	1.0330D+01	1.9425D+01	2.0604D+01	2.1143D+01	6.8689D+00	7.5101D+00	1.0128D+00	3.0323D+01	4.0999D-02
1.3537D-02	6.2666D-03	3.3535D-03	1.8867D-03	6.1878D-04	4.9361D-04	2.6019D-04	1.2493D-04	5.1112D-05	3.7239D-05
TOTAL NEUTRON ENERGY LOSS									
5.3827D-03	1.5983D-02	3.8586D-02	9.7648D-02	1.5599D-01	2.3103D-01	8.5938D-01	1.4281D+00	2.5209D+00	8.1550D+00
1.0717D+01	1.0693D+01	1.4574D+01	1.6445D+01	1.9366D+01	1.4431D+01	1.2607D+01	3.1412D-02	5.2385D-03	2.1143D-05
6.8684D-06	3.2259D-06	1.7262D-06	9.7115D-07	1.6565D-08	9.2094D-09	4.8545D-09	2.3309D-09	9.5362D-10	6.9963D-11
ESCAPE NEUTRON PROBABILITY BY GROUP									
7.8529D-01	7.8530D-01	6.4996D-01	4.6212D-01	4.6669D-01	7.2313D-01	9.7749D-01	9.9927D-01	9.9997D-01	1.0000D+00
7.3828D-01	7.3976D-01	5.8879D-01	3.9896D-01	4.0316D-01	6.7513D-01	9.7180D-01	9.9910D-01	9.9997D-01	1.0000D+00
7.0360D-01	7.0649D-01	5.4684D-01	3.5963D-01	3.6336D-01	6.4170D-01	9.6738D-01	9.9897D-01	9.9996D-01	1.0000D+00
6.6339D-01	6.6829D-01	5.0140D-01	3.2023D-01	3.2388D-01	6.0475D-01	9.6200D-01	9.9882D-01	9.9996D-01	1.0000D+00
6.2212D-01	6.2952D-01	4.5823D-01	2.8536D-01	2.8876D-01	5.6819D-01	9.5625D-01	9.9865D-01	9.9995D-01	9.9999D-01
5.9164D-01	6.0114D-01	4.2829D-01	2.6256D-01	2.6592D-01	5.4291D-01	9.5170D-01	9.9852D-01	9.9994D-01	9.9999D-01
5.7193D-01	5.8292D-01	4.0984D-01	2.4891D-01	2.5221D-01	5.2672D-01	9.4870D-01	9.9843D-01	9.9994D-01	9.9999D-01
5.6394D-01	5.7562D-01	4.0269D-01	2.4352D-01	2.4666D-01	5.2083D-01	9.4764D-01	9.9840D-01	9.9994D-01	9.9999D-01
5.6059D-01	5.7268D-01	3.9994D-01	2.4116D-01	2.4402D-01	5.1939D-01	9.4751D-01	9.9840D-01	9.9994D-01	9.9999D-01
5.5423D-01	5.6710D-01	3.9481D-01	2.3677D-01	2.3909D-01	5.1657D-01	9.4724D-01	9.9839D-01	9.9994D-01	9.9999D-01
5.4673D-01	5.6053D-01	3.8879D-01	2.3171D-01	2.3383D-01	5.1352D-01	9.4693D-01	9.9839D-01	9.9994D-01	9.9999D-01