





DEVELOPMENT OF CR, ND: GSGG LASER AS A PUMPING

**SOURCE OF TI: SAPPHIRE LASER** 

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Koji TAMURA and Takashi ARISAWA

日本原子力研究所 Japan Atomic Energy Research Institute

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# Development of Cr,Nd:GSGG Laser as a Pumping Source of Ti:Sapphire Laser

Koji TAMURA and Takashi ARISAWA+

Department of Materials Science Tokai Research Establishment Japan Atomic Energy Research Institute Tokai-mura, Naka-gun, Ibaraki-ken

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Since efficiency of Cr,Nd doped gadolinium scandium gallium garnet (GSGG) laser is in principle higher than that of Nd:YAG laser, it can be a highly efficient pumping source for Ti:sapphire laser. We have made GSGG laser, and measured its oscillation properties. It was two times more efficient than Nd:YAG laser at free running mode operation. At Q-switched mode operation, fundamental output of 50mJ and second harmonics output of 8mJ were obtained. The developed laser had appropriate spatial profile, temporal duration, long time stability for solid state laser pumping. Ti:sapphire laser oscillation was achieved by the second harmonics of GSGG laser.

Keywords: GSGG laser, Ti:sapphire Laser, DLAP, AVLIS, Chromium, Neodymium

<sup>&</sup>lt;sup>+</sup> Advanced Photon Research Center, Kansai Research Establishment

チタンサファイアレーザー励起用光源としての Cr,Nd:GSGG レーザーの開発

### 日本原子力研究所東海研究所物質科学研究部 田村 浩司·有沢 孝<sup>+</sup>

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クロム、ネオジムドープの GSGG レーザーはヤグレーザよりも原理的に効率が高く、チタンサファイアレーザーなどの高効率励起用光源となる可能性がある。そのため、GSGG レーザーを試作しその発振特性を調べた。発振効率はフリーランニングモードで Nd:YAG レーザーの 2 倍であった。Q スイッチモードにおいて基本波で 50mJ、第二高調波で 8mJ を得た。レーザーの空間分布、時間波形、長時間安定性は固体レーザー励起に適したものであった。GSGG レーザーの第二高調波励起によりチタンサファイアレーザーの発振が達成できた。

東海研究所:〒319-1195 茨城県那珂郡東海村白方白根 2-4

<sup>+</sup> 関西研究所光量子科学研究センター

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#### 1. Introduction

A compact tunable laser light source is desirable as an alternative to conventional dye laser for high resolution spectroscopy<sup>1)</sup>, remote sensing<sup>2)</sup>, and isotope separation<sup>3)</sup>. Titanium-doped sapphire laser (TSL) is promising because of its advantageous features such as the broad gain region extending over 400nm, easy handling of solid state brick, and stable operation without vibration by dye circulation. Especially, TSL has possibility to become one of the promising light source for atomic vapor laser isotope separation (AVLIS), because some of the requirements of AVLIS light source such as broad tunability, narrow line width, stable operation are achievable by TSL.

The light source for material processing like isotope separation requires high efficiency for the economy of the process. The main pumping source for TSL is frequency doubled Nd:YAG laser. But its efficiency is usually not more than a few percent. Therefore, it is desirable to increase the pumping efficiency. Sensitization of the laser medium is a possible approach to high efficiency. One of the promising materials for the sensitization is  $Cr^{3+}$ ,  $Nd^{3+}$  doped gadolinium scandium gallium garnet (GSGG). Broad absorption band of doped Cr ion well match the emission spectrum of flash lamp spectra and absorb the lamp energy efficiently. The excited energy is non-radiatively transferred from Cr ion to upper level of Nd in GSGG host crystal, although it is slow and inefficient in YAG host crystal due to relatively small lattice parameter<sup>4</sup>). Since the energy transfer rate from Cr ion to Nd is very efficient (0.86), it is almost as efficient as direct Nd pumping<sup>5</sup>). Therefore, it is expected that Cr,Nd:GSGG laser can be a more efficient pump source compared with Nd:YAG laser. Zharikov et al. reported that the differential efficiency of Cr,Nd:GSGG was about 2 times higher than that of Nd:YAG at free running mode operation<sup>6</sup>).

Rather large diameter crystal of GSGG up to 2.5-inch with high optical quality can be grown by Czochralski method<sup>7)</sup>, although it is difficult for YAG bole which generates a center core during the crystal growth. It is desirable feature to scale up the laser for massive material processing which requires large laser medium.

The coexistence of Cr and Nd in GSGG crystal can reduce solarization by radiation. It was reported that GSGG laser output was unaffected to  $^{60}$ Co  $\gamma$ -ray irradiation up to 10Mrad, although Nd:YAG output dropped by an order of magnitude after 1Mrad  $\gamma$ -ray exposure  $^{8)}$ . These properties make GSGG more desirable laser for atomic energy applications or space applications where strong radiation is expected.

As a pump laser of TSL smooth spatial profile to avoid damage to the crystal surface and several J/cm<sup>2</sup> of fluence to oscillate dispersive cavity are required. About 50mJ of fundamental output is desirable to oscillate TSL laser considering SHG efficiency, spatial profile, and losses during optical delivery.

In this study, in order to investigate the possibility of efficient tunable laser system for AVLIS pumped by GSGG laser, Nd, Cr:GSGG laser was developed. Free running and Q-switched mode operations were achieved. TSL oscillation pumped by the second harmonics of GSGG was demonstrated.

## 2. Apparatus

The goal of our development was to generate green light from GSGG laser that can oscillate TSL resonator. Figure 1 shows schematic diagram of the developed laser system. The laser system consisted of GSGG oscillator (GSGG crystal, mirror, lamp house, E/O Q-switch, polarizer, shutter), SHG crystal, TSL oscillator. The GSGG oscillator cavity consisted of a flat output coupler and a highly reflecting concave mirror, with a lamp house between mirrors. Nd, Cr:GSGG crystal (4mm diameter × 75mm length) was obtained from Tokin Co. Ltd. The concentration of Cr and Nd was

1% and 2%, respectively. The crystal was set in a silver coated lamp house with elliptical shape, which was cooled with purified water. The silver coating has better reflection especially in UV region than the conventional Au coating, which would improve pumping efficiency of flash lamp. A Xe flash lamp light was focused into GSGG crystal. The flash lamp was simmered and driven by a conventional discharge circuit. The half-height pulse width of the current to the flash lamp was  $340 \mu$  s. The cavity was supported with 3 invar rod (20 mm  $\phi \times 1.2$  mL) to prevent thermal and mechanical disturbance. An iris of 1.5 mm  $\phi$  was inserted in the resonator in order to limit the number of transverse modes in operation and to investigate its effect to output energy.

#### 3. Results

The designed laser resonator was used as a free running oscillator to optimize the laser system and to compare the oscillation efficiency. The oscillation was in  $1.06 \,\mu$ m wavelength range which corresponds to  ${}^4F_{3/2}$ - ${}^4I_{11/2}$  transition of Nd<sup>3+</sup> ion. Figure 2 shows the laser output energy as a function of electrical input energy to the flash lamp for a Nd, Cr:GSGG rod and a Nd:YAG rod tested in the same resonator. The result shows that in this experiment GSGG was about two times more efficient than YAG laser, although oscillation threshold energy of GSGG was slightly higher than that of YAG.

Figure 3 shows the output energy versus the lamp input at the free running mode operation with output coupler of 50 percent reflectivity and 1m cavity length. The repetition rate was 1Hz. By insertion of iris, the output power decreased considerably, and the oscillation threshold energy increased. On the other hand, the beam size 1.5 m away from the output coupler decreased from 3.6 mm  $\phi$  to 1.7 mm  $\phi$ , and smooth spatial mode close to TEM<sub>00</sub> mode was obtained.

Based on these investigations about the optical components such as mirror reflectivity, mirror curvature, the influence of iris, etc., the oscillator configuration was optimized. The optimized laser resonator consisted of a concave total reflector with 3 m radius of curvature and a plane mirror of 20 percent reflectivity separated by 65 cm. The reflectivity and the curvature of the selected mirrors maximized the output energy among several available mirrors. After these optimizations, the slope efficiency increased form 0.2 % to 1 %, and a smooth spatial mode close to  $TEM_{00}$  was obtained without iris.

Figure 4 shows the temporal profile of GSGG fundamental output at the free running mode operation detected with a PIN photo diode. The repetition rate was 10 Hz, and the electrical input to the flash lamp was 12.2 J/pulse. The GSGG laser pulse appeared as a series of short spikes. The duration of the pulse envelope was 120  $\mu$  sec. Figure 5 shows the spatial profile of fundamental output detected with CCD camera (Hitachi Co. KP-140).

A Q-switched mode oscillation was performed, which was achieved with a Pockel cell and a polarizer in the cavity. Figure 6 shows the temporal profile of generated laser pulse measured with photomultiplier at the repetition rate of 10 Hz. Post lasing caused by the residual distortion of E/O Q-switch observed 0.7-several  $\mu$  sec after the main oscillation was reduced sufficiently by adjusting optical alignment. The output beam from the resonator was vertically polarized. Figure 7 shows spatial profile of the fundamental output. The observed transverse mode profile was close to  $\text{TEM}_{00}$  mode, and it was appropriate for TSL pumping without intense spot that would cause crystal damage. The beam size was 2.5 mm  $\phi$  and the beam divergence was 2.4 mrad.

Figure 8 shows the fundamental output of Q-switched laser versus flash lamp input energy at a repetition rate of 10Hz. In this case, the laser output was proportional to the flash lamp input power. The maximum fundamental output of 50 mJ was obtained. The conversion efficiency from electric power to laser output was 0.4 percent, and the slope efficiency was 1 percent.

The thermal properties of GSGG host crystal is poor compared with that of YAG crystal. In order to examine the thermal effect to GSGG output, long time operation for over 60 minutes was measured

at the repetition rate of 1Hz and 10Hz. Figure 9 shows the laser output energy as a function of operation time. At a repetition rate of 1Hz, the output power did not change considerably in time. On the other hand, at a repetition rate of 10Hz, the output power decreased by 6 percent within the initial 1 minute after the start of oscillation, and stabilized. Therefore, within these repetition rates, stable operation for daily operational time was possible without serious thermal effect to the output power.

The second harmonics of the fundamental output was generated with a nonlinear crystal, deuterated L-Arginine Phosphate Monohydrate (DLAP) which was obtained from Creveland crystal Co. The size of the crystal was 10 mm×10 mm×15 mm length. It was reported that LAP has better SHG conversion efficiency than KDP and the damage threshold is two times higher than that of KDP<sup>9</sup>, which were desirable for efficient laser operation. Figure 10 shows the temporal profile of the generated second harmonics. The half-height temporal duration was 12 nsec. The profile was smooth and appropriate for TSL pumping. The divergence was 1.8 mrad. Figure 11 shows SHG output fluence versus fundamental laser fluence. The maximum SHG energy of 8 mJ was obtained at the fundamental input of 50 mJ, which corresponds to conversion efficiency of 16 percent. In this experiment, the SHG conversion efficiency was low, which was caused mainly by several nm of broad spectral line width of GSGG oscillator and rather low fluence of the fundamental output. It is possible to improve the conversion efficiency by narrowing the line width with a dispersive element, or intensify the fundamental fluence.

In order to investigate the acceptable temperature range of the laser when it was operated at a high repetition rate or at high power, the temperature dependence of the laser output was measured. Figure 12 shows the SHG output as a function of DLAP house temperature. In this experimental setup, the acceptable temperature range was about 5degree. Figure 13 shows the relative SHG output versus tuning angle of DLAP. In this case, the acceptable tuning angle was about  $\pm 1$ mrad.

The generated second harmonics was then introduced to TSL oscillator. The TSL oscillator consisted of a plane total reflector, a plane output coupler of 95 percent reflection, and a TSL crystal (Union Carbide Co., 5 mm×5 mm×12 mmL). The Ti concentration was 0.15 wt%, and FOM was 90. The pump laser was focused with a lens (f=50cm) before the crystal surface, and longitudinally pumped the crystal. Figure 14 shows the temporal profile of the generated TSL pulse which was 50 nsec delayed from the pump beam. The temporal duration of TSL was 12 nsec. Figure 15 shows the laser output as a function of incident pump energy. The threshold fluence of oscillation was 2.0 J/cm². The slope efficiency was 6.7 %. Figure 16 shows the vertical and horizontal beam profile of generated laser. The beam divergence was 1.8 mrad. The beam profile of the obtained TSL was smooth, and was appropriate for various applications.

We have developed Nd,Cr:GSGG oscillator that had sufficient output power and spatial profile to oscillate TSL. Although at this stage, the conversion efficiency of the total system is low, and further optimizations of the laser system are required, it is demonstrated that the GSGG laser is potentially more efficient pump source for TSL than YAG laser.

Since the thermal properties of GSGG crystal is not good, the influence of the thermal lensing or thermally induced birefringence might be larger compared with that of YAG. However, these influence can be reduced by using a slab type resonator with zigzag optical path<sup>10)</sup> or a flow cooled multiple disk type resonator<sup>11)</sup>, which make high average power GSGG laser system possible.

#### 4. Conclusions

Cr,Nd doped GSGG laser has been developed. It was about two times more efficient than Nd:YAG laser at the free running mode operation. At the Q-switched mode operation, fundamental output of 50 mJ, and SHG output of 8 mJ were obtained. Appropriate spatial profile, temporal duration, and long time stability for TSL pumping were achieved. Oscillation of TSL was demonstrated by the second harmonics of GSGG laser.

# Acknowledgment

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#### References

- 1) Vassen W., Zimmermann C., Kallenbach R. and Hansch T.W.: Optics Comm., 75, 435 (1990).
- 2) Brockman P., Barnes C.H., Hess R.V. and Browell E.V.: Opt. Lett., 11, 712 (1986).
- 3) Tamura K., Oba M. and Arisawa T.: Appl. Opt., 32, 987 (1993).
- 4) Avanesov A.G., Denker B.I., Osiko V.V., Ostroumov V.G., Sakun V.P., Smirnov V.A. and Shcherbakov I.A.: Sov.J.Quantum Electron., 12, 421 (1982).
- 5) Pruss D., Huber G. and Beimowski A.: Appl. Phys., <u>B28</u>, 355 (1982).
- 6) Zharikov E.V., Il'ichev N.N., Laptev V.V., Malyutin A.A., Ostroumov V.G., Pashinin P.P., Pimenov A.S., Smirnov V.A. and Shcherbakov I.A.: Sov.J.Quantum Electron., 13, 82 (1983).
- 7) Bruni F.J.: UCRL-15735 (1985).
- 8) Zharikov E.V., Kuratev I.I., Laptev V.V., Naselskii S.P., Ryabov A.I., Toropkin G.N., Shestakov A.V. and Shcherbakov I.A.: Bull.Acad.Sci.USSR, Phys.Ser., 48, 103 (1984).
- 9) Eimerl D., Velsko S., Davis L., Wang F., Loiacono G. and Kennedy G.: IEEE J.Quantum Electron., 25, 179 (1989).
- 10) Yoshida K., Yoshida H. and Kato Y.: IEEE J.Quantum Electron., 24, 1188 (1988).
- 11) Emmett J.L., Krupke W.F. and Sooy W.R.: UCRL-53571 (1984).

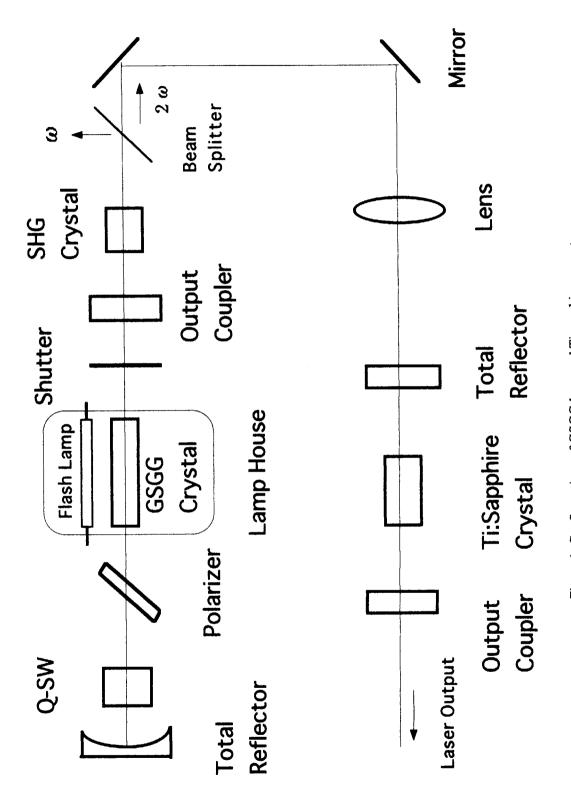


Figure 1. Configurations of GSGG laser and Ti:sapphire resonator.

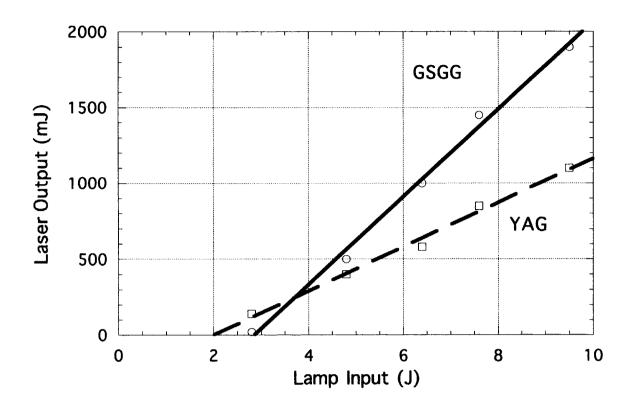


Figure 2. The output energy of GSGG and YAG resonator as a function of lamp input energy.

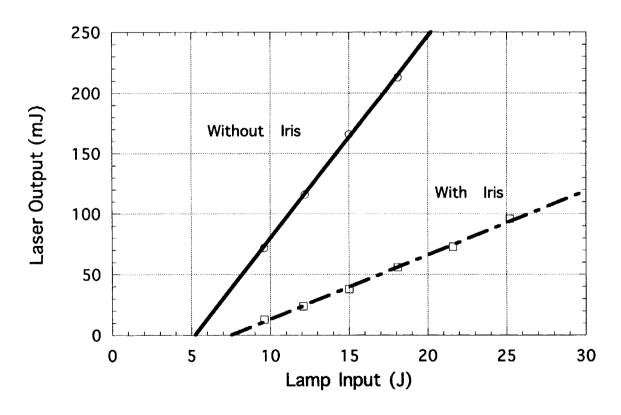


Figure 3. The output energy of GSGG laser as a function of lamp input energy with and without iris.

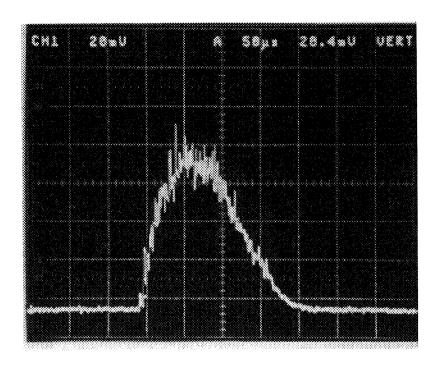


Figure 4. Temporal profile of GSGG fundamental output at a free running mode operation.

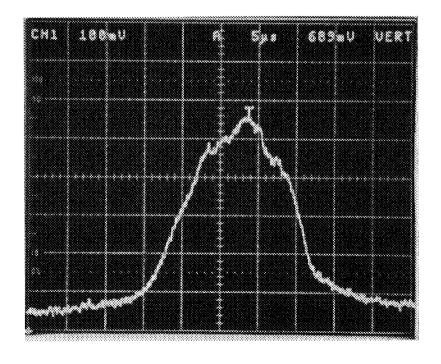


Figure 5. Spatial profile of GSGG fundamental output.

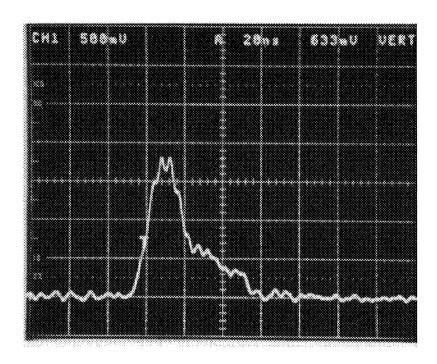


Figure 6. Temporal profile of GSGG fundamental output.

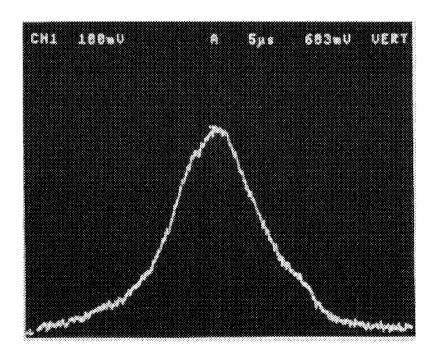


Figure 7. Spatial profile of Q-switched GSGG fundamental output.

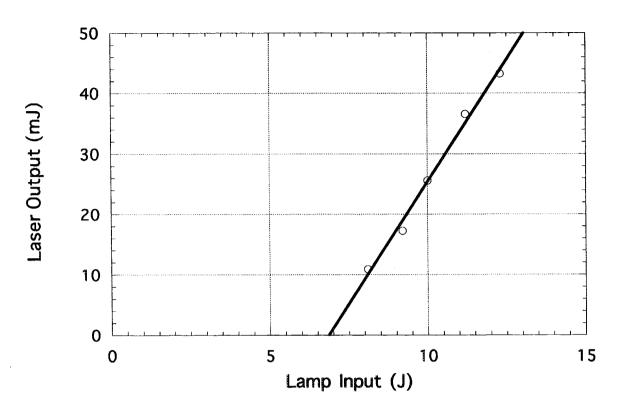


Figure 8. Fundamental output of Q-switched GSGG laser as a function of lamp input energy.

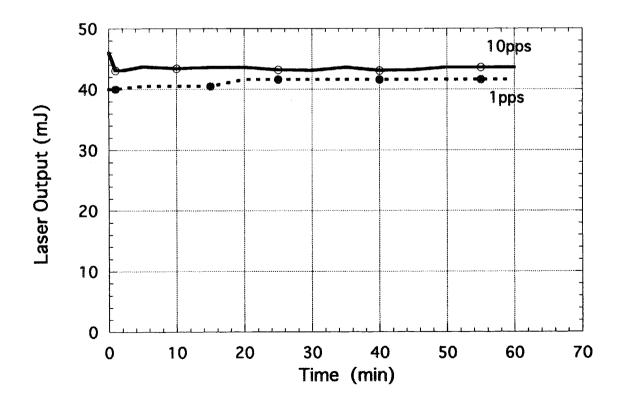


Figure 9. Laser output energy as a function of operation time at repetition rates of 1Hz and 10Hz.

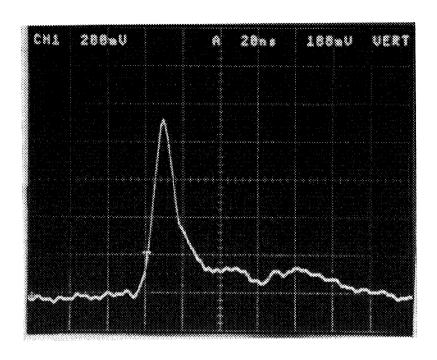


Figure 10. Temporal profile of second harmonics of GSGG laser.

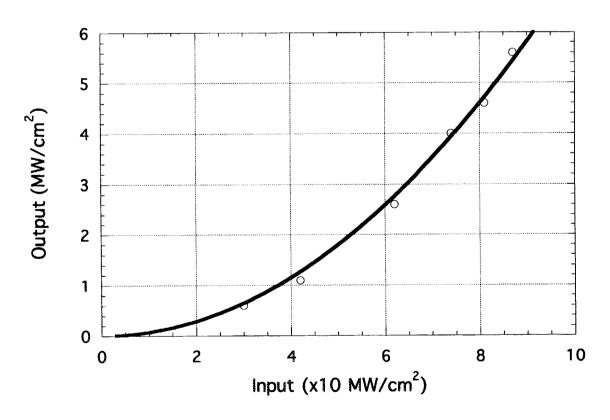


Figure 11. Second harmonics fluence as a function of fundamental fluence.

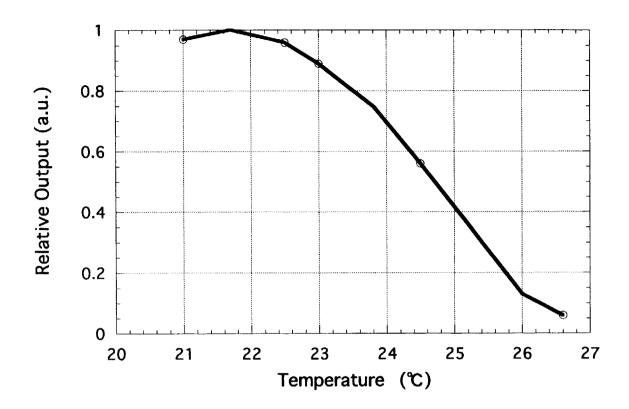


Figure 12. Relative SHG output as a function of the crystal temperature.

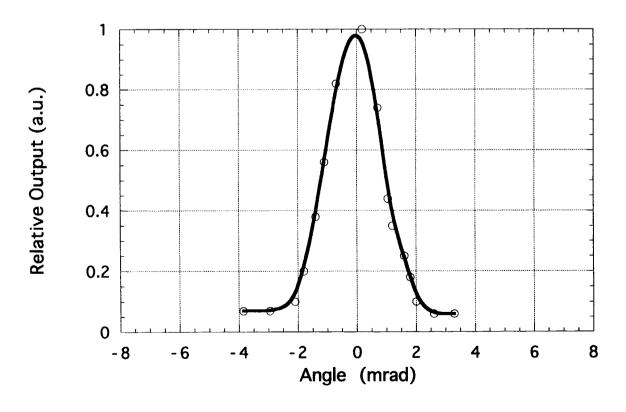


Figure 13. Relative SHG output as a function of the tuning angle.

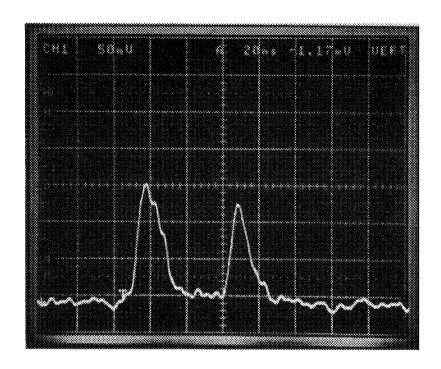


Figure 14. Temporal profile of GSGG laser and generated TSL.

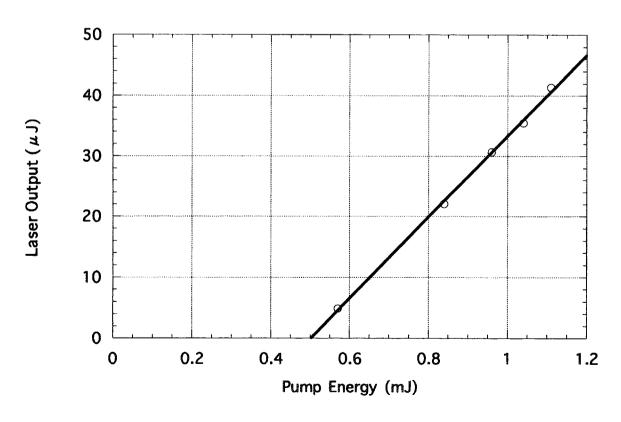


Figure 15. TSL output energy as a function of GSGG input energy.

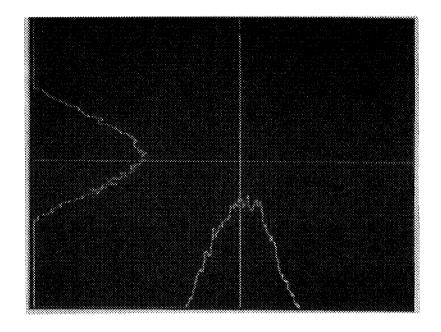


Figure 16. Vertical and horizontal beam profile of TSL.

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# 国際単位系 (SI) と換算表

表1 SI基本単位および補助単位

量		名 称	記 号_
長	5	メートル	m
質	量	キログラム	kg
時	間	秒	s
電	流	アンペア	Α
熱力学	温度	ケルビン	K
物質	量	モル	mol
光	度	カンデラ	cd
平面	角	ラジアン	rad
立体	角	ステラジアン	sr

表3 固有の名称をもつ SI 組立単位

遺	名 称	記号	他のSI単位 による表現
周 波 数	ヘルッ	Hz	<b>s</b> -1
カ	ニュートン	N	m·kg/s²
圧 力 , 応 力	パスカル	Pa	N/m²
エネルギー,仕事,熱量	ジュール	J	N∙m
工率, 放射束	ワット	W	J/s
電気量,電荷	クーロン	C	A·s
電位,電圧,起電力	ボルト	V	W/A
静電容量	ファラド	F	C/V
電 気 抵 抗	オーム	Ω	V/A
コンダクタンス	ジーメンス	S	A/V
磁 束	ウェーバ	Wb	V-s
磁束密度	テスラ	Т	Wb/m <sup>2</sup>
インダクタンス	ヘンリー	Н	Wb/A
セルシウス温度	セルシウス度	℃	
光 束	ルーメン	lm	cd·sr
照 度	ルクス	lx	lm/m²
放 射 能	ベクレル	Bq	s-1
吸収線量量	グレイ	Gy	J/kg
線量当量	シーベルト	Sv	J/kg

表2 SIと併用される単位

名 称	記号
分, 時, 日 度, 分, 秒	min, h, d
リットル	l, L
トン	t
電子ボルト	eV
原子質量単位	u

 $1 \text{ eV} = 1.60218 \times 10^{-19} \text{ J}$  $1 u = 1.66054 \times 10^{-27} kg$ 

表4 SIと共に暫定的に 維持される単位

	名 称		記	号
オンタ	グストロ	- L	Ä	
バ		ン	k	)
ノギ	-	ル	b	ar
ガ		ル	G	al
丰	ューリ	-	C	'i
レン	ノトケ	・ン	F	₹
ラ		F	га	ad
レ		4	re	em

 $1 \text{ Å} = 0.1 \text{ nm} = 10^{-10} \text{ m}$ 1 b=100 fm<sup>2</sup>= $10^{-28}$  m<sup>2</sup> 1 bar=0.1 MPa=10<sup>5</sup> Pa  $1 \text{ Gal} = 1 \text{ cm/s}^2 = 10^{-2} \text{ m/s}^2$  $1 \text{ Ci} = 3.7 \times 10^{10} \text{ Bg}$  $1 R = 2.58 \times 10^{-4} C/kg$  $1 \text{ rad} = 1 \text{ cGy} = 10^{-2} \text{Gy}$  $1 \text{ rem} = 1 \text{ cSv} = 10^{-2} \text{ Sv}$ 

表 5 SI接頭語

倍数	接頭語	記号
1018	エクサ	E
1015	ペタ	P
1012	テ ラ	T
109		G
10°	メーガー	M
10³	+ 0	k
10²	ヘクト	h
10 ¹	デ カ	da
10-1	デ シ	d
10-2	センチ	С
$10^{-3}$	ミ リ	m
10-6	マイクロ	μ
10-9	ナノ	n
$10^{-12}$	ピコ	р
10-15	フェムト	f
10-18	アト	а

(注)

- 1. 表1-5は「国際単位系」第5版,国際 度量衡局 1985年刊行による。ただし、1 eV および 1 u の値は CODATA の 1986 年推奨 値によった。
- 2. 表4には海里、ノット、アール、ヘクタ ールも含まれているが日常の単位なのでこ こでは省略した。
- 3. barは、JISでは流体の圧力を表わす場 合に限り表2のカテゴリーに分類されてい
- 4. EC閣僚理事会指令では bar, barn およ び「血圧の単位」mmHg を表2のカテゴリ - に入れている。

力	N(=10 <sup>5</sup> dyn)	kgf	lbf
	1	0.101972	0.224809
	9.80665	1	2.20462
	4.44822	0.453592	1

mt + D (NT / 2) +0.D(18 -> -t)/ //

粘度	$1 \text{ Pa·s}(\text{N·s/m}^{\prime}) = 10 \text{ P}(\pi 7 \text{ X}) (\text{g/(cm·s)})$
動粘度	$1 \text{ m}^2/\text{s} = 10^4 \text{St}(2 \text{ h} - 2 \text{ m})(\text{cm}^2/\text{s})$

圧	MPa(=10 bar)	kgf/cm²	atm	mmHg(Torr)	lbf/in²(psi)
	1	10.1972	9.86923	7.50062 × 10 <sup>3</sup>	145.038
カ	0.0980665	1	0.967841	735.559	14.2233
	0.101325	1.03323	1	760	14.6959
	1.33322 × 10 <sup>-4</sup>	$1.35951 \times 10^{-3}$	$1.31579 \times 10^{-3}$	1	1.93368 × 10 <sup>-2</sup>
	6.89476 × 10 <sup>-3</sup>	$7.03070 \times 10^{-2}$	$6.80460 \times 10^{-2}$	51.7149	1

I	J(=10'erg)	kgf•m	kW•h	cal(計量法)	Btu	ft • lbf	eV	1 cal = 4.18605 J (計量法)
ネル	1	0.101972	2.77778 × 10 <sup>-1</sup>	0.238889	9.47813 × 10 <sup>-4</sup>	0.737562	6.24150 × 10 18	= 4.184 J (熱化学)
ギー	9.80665	1	$2.72407 \times 10^{-6}$	2.34270	9.29487 × 10 <sup>-3</sup>	7.23301	6.12082 × 10 <sup>19</sup>	= 4.1855 J (15 °C)
・ 仕 事	$3.6 \times 10^{6}$	3.67098 × 10 <sup>5</sup>	1	8.59999 × 10 5	3412.13	2.65522 × 10 <sup>6</sup>	2.24694 × 10 <sup>25</sup>	= 4.1868 J (国際蒸気表)
•	4.18605	0.426858	1.16279 × 10 <sup>-6</sup>	1	$3.96759 \times 10^{-3}$	3.08747	2.61272×1019	仕事率 1 PS (仏馬力)
熱量	1055.06	107.586	2.93072 × 10 <sup>-4</sup>	252.042	1	778.172	$6.58515 \times 10^{21}$	$= 75 \text{ kg} \text{f} \cdot \text{m/s}$
	1.35582	0.138255	$3.76616 \times 10^{-7}$	0.323890	1.28506 × 10 <sup>-3</sup>	1	8.46233 × 10 <sup>18</sup>	= 735.499 W
	1.60218 × 10 <sup>-19</sup>	1.63377 × 10 <sup>-20</sup>	$4.45050 \times 10^{-26}$	$3.82743 \times 10^{-20}$	1.51857 × 10 <sup>-22</sup>	1.18171 × 10 <sup>-19</sup>	1	

放	Bq	Ci
射	1	2.70270 × 10 <sup>-11</sup>
能	3.7 × 10 <sup>10</sup>	1

吸	Gy	rad
吸収線量	1	100
觝	0.01	1

照	C/kg	R
射線量	1	3876
重	2.58 × 10 <sup>-4</sup>	1

線	Sv	rem
線量当量	1	100
重	0.01	1